

COSMIC-RAY DRIVEN DYNAMO IN GALACTIC DISKS

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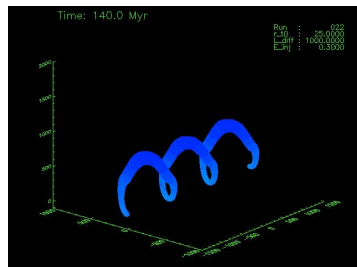
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(Hanasz & Lesch 2000, ApJ, 543, 235)

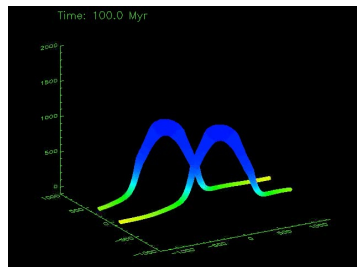
⇒ helical magnetic loops form on initially azimuthal magnetic field due to buoyancy of cosmic rays and the Coriolis force



⇒ small scale loops reconnect to form larger loops

⇒ generation of the large-scale radial m.f.

⇒ differential rotation: generation of the azimuthal m.f.



$$\frac{\partial \mathbf{V}}{\partial t} + (\mathbf{V} \cdot \nabla) \mathbf{V} = -\frac{1}{\rho} \nabla(p + p_{CR}) + \mathbf{g} - 2\boldsymbol{\Omega} \times \mathbf{V} + \boldsymbol{\Omega}^2 \mathbf{r} + \frac{1}{\rho} \nabla \left(\frac{B^2}{8\pi} \right) + \frac{\mathbf{B} \cdot \nabla \mathbf{B}}{4\pi\rho}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{V}) = 0$$

$$\frac{\partial e}{\partial t} + \nabla \cdot (e \mathbf{V}) = -p \nabla \cdot \mathbf{V}$$

$$p = (\gamma - 1)e$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{V} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B}$$

Diffusion - advection approximation

(eg. Schlickeiser & Lerche 1985, A&A, 151, 151)

$$\frac{\partial e_{\text{cr}}}{\partial t} + \nabla(e_{\text{cr}} \mathbf{V}) = -p_{\text{cr}} \nabla \mathbf{V} + \nabla(\hat{K} \nabla e_{\text{cr}}) \quad (1)$$

+ CR sources (SN remnants)

$$p_{\text{cr}} = (\gamma_{\text{cr}} - 1)e_{\text{cr}} \quad (2)$$

Anisotropic diffusion of CRs

(Giaccalone & Jokipii 1998 , Jokipii 1999, Ryu et al. 2003)

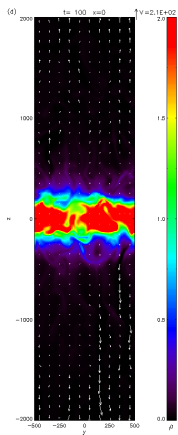
$$K_{ij} = K_{\perp} \delta_{ij} + (K_{\parallel} - K_{\perp}) n_i n_j, \quad n_i = B_i / B, \quad (3)$$

$$K_{\parallel} = 3 \cdot 10^{28} \text{cm}^2 \text{s}^{-1}, \quad K_{\perp} = (1 - 10)\% (K_{\parallel})$$

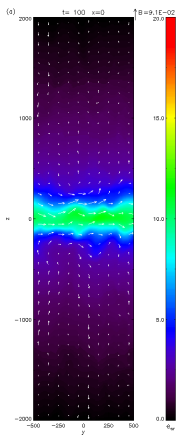
Original idea: Parker (1992)

Numerical model: Hanasz, Kowal, Otmianowska-Mazur & Lesch, 2004)

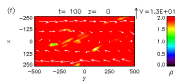
- **the cosmic ray component:** diffusion-advection transport equation (Hanasz and Lesch 2003 - numerical algorithm).
- **localized sources of cosmic rays:** supernova remnants, exploding randomly in the disk volume, SN shocks & thermal effects neglected
- **resistivity of the ISM** (see Hanasz, Otmianowska-Mazur and Lesch 2002, and Hanasz and Lesch 2003, Kowal, Hanasz & Otmianowska-Mazur 2003) $\Rightarrow$ magnetic reconnection.
- **sheared rotation** (+ Coriolis and tidal forces in local simulations)
- **realistic vertical disk gravity** following the model of ISM in the Milky Way by Ferriere (1998)



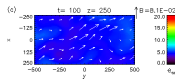
$$\rho(y, z), \vec{V}(y, z)$$



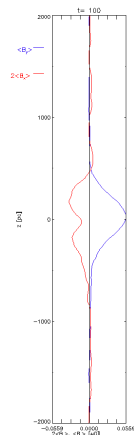
$$e_{cr}(y, z), \vec{B}(y, z)$$



$$\rho(x, y), \vec{V}(x, y)$$



$$e_{cr}(x, y), \vec{B}(x, y)$$



$$\langle B_r \rangle(z), \langle B_\phi \rangle(z)$$

Slices through the computational box at $R_G = 5\text{kpc}$

& horizontally averaged magnetic field components.

Hanasz, et al 2004 ApJ, 543, 235; 2006 AN 327, 469; 2009 A&A 498, 335)

Similar results for SN-thermally driven dynamo by Gressel et al 2008

- Weak influence of SN rate on magnetic field amplification in the range of 10% - 100% of the realistic value ($f_{SN} = 130/\text{Myr} / \text{kpc}^2$ at $R \simeq 5\text{kpc}$).
- The efficiency of m.f. amplification increases with rising $K_{\parallel}$ (up to $3 \times 10^{28} \text{cm}^2 \text{s}^{-1}$) and decreases with rising $K_{\perp} \Rightarrow$ **anisotropic diffusion** $K_{\perp}/K_{\parallel} \leq 10\%$ necessary for efficient action of CR-driven dynamo.
- Final saturated magnetic fields reach energetic equipartition with gas kinetic energy
- CR energy density remains an order of magnitude larger than gas and magnetic energy densities in local CR-driven dynamo models.

- Dynamo coefficients in Parker unstable disks with cosmic rays and shear (Kowal et al. 2006, A&A 445, 915)
- Computation of α and η tensors in the framework of Kleeorin et al. (2003) and Blackman & Field (2003) nonlinear $\alpha - \omega$ dynamo theories (Otmianowska-Mazur et al. 2007, ApJ, 668, 110)
- Synthetic radio-maps of a global galactic disk based on local CR-driven dynamo models exhibit X-type structures (Otmianowska-Mazur et al 2009, ApJ, 693,1)
- Demonstration that CR-driven dynamo can work in physical conditions (rotation, shearing rate) of irregular galaxies (Siejkowski et al 2010, A&A, 510, 97)

- CR-driven dynamo works in a wide range of parameter space!
- but ... CR excess is worrying.

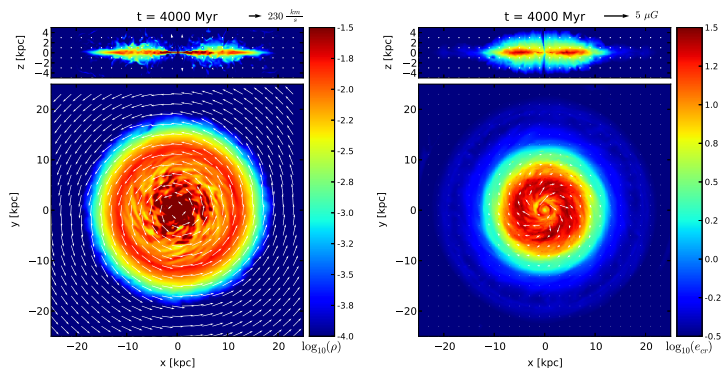
Improvements necessary:

- Global galactic disk simulations → escape of CR along horizontal magnetic field in galactic plane, to avoid the CR excess.

GALACTIC DISK MODEL

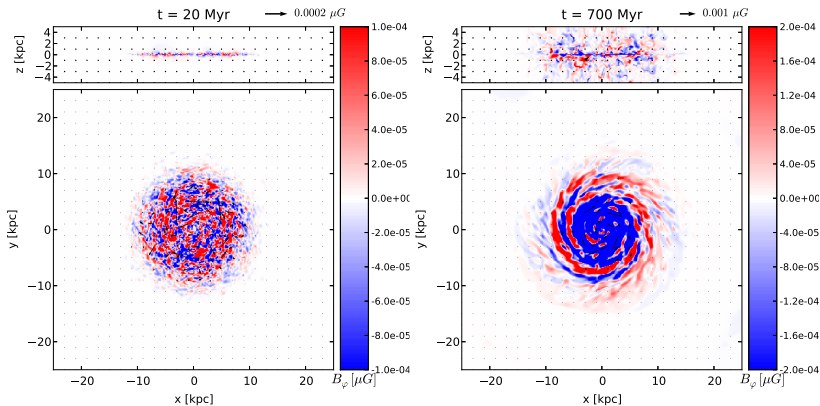
- Galactic gravitational potential: halo+bulge+disk (Allen & Santillan 1991)
- Flat rotation curve for $R_G \geq 3\text{kpc}$
- Interstellar gas: Global model of ISM for the Milky Way (Ferriere 1998), molecular ring at $R_G = 4.5\text{kpc}$
- **No magnetic field and no CRs at $t = 0$**
- SN rate $\propto$ star formation rate $\propto$ to gas column density: maximum of SN activity at $R_G = 4.5\text{kpc}$, Gaussian distribution of SNe in z-coordinate ($H = 200\text{pc}$)
- 10% of of SN energy output is converted to CR energy.
- **weak ($10^{-4}\mu\text{G}$) dipolar, small scale ($r \sim 50\text{pc}$) randomly oriented magnetic field is supplied locally with every SN explosion for $t \leq 1\text{Gyr}$**
- resistive dissipation of small-scale magnetic fields.

Multifluid, parallel (MPI) magnetohydrodynamical code PIERNIK (Hanasz et al. 2008): <http://piernik.astro.uni.torun.pl>, GALERA, TASK Gdańsk, resol. 500x500x200 up to 1000x1000x200 grid cells, $\simeq 100\text{k}-250\text{k}$ CPU h per experiment, 400-1600 CPU cores.



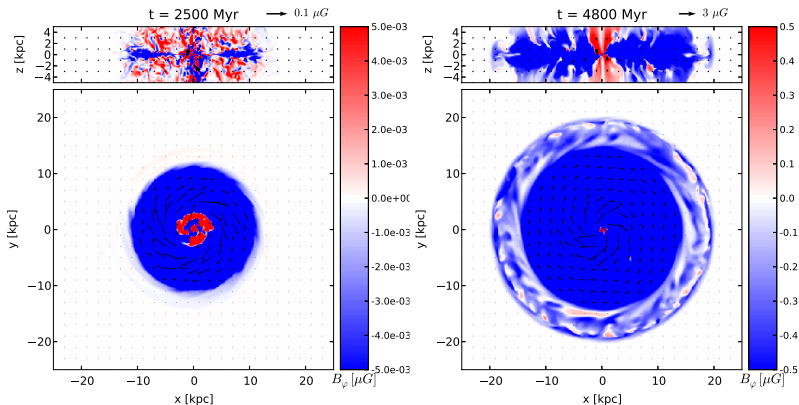
Gas density + vectors of gas velocity (left), a and cosmic ray energy density + vectors of magnetic field at $t = 4\text{Gyr}$.

Hanasz, Woltański, Kowalik 2009, ApJ Letter 706L, 155



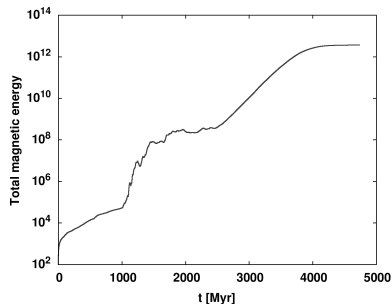
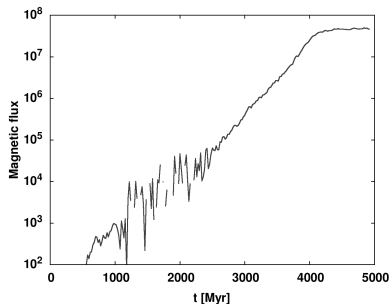
Colours: – azimuthal (toroidal) magnetic field component **blue**: $B_\phi < 0$,
red: $B_\phi > 0$

Exploding magnetized stars spread weak irregular magnetic fields in the interstellar medium



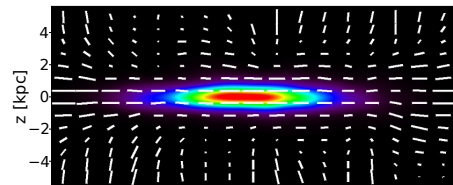
Colours: – azimuthal (toroidal) magnetic field component **blue**: $B_\phi < 0$,
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Cosmic rays resulting from Supernova Explosions, and disk rotation cause amplification and ordering of magnetic field in the interstellar medium

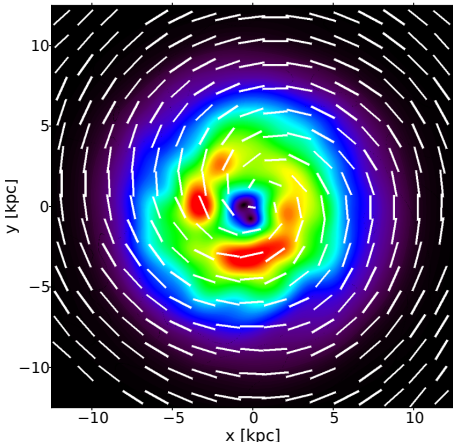


Amplification timescale of the large-scale magnetic field component:

$$T_{\langle B \rangle} = 270 \text{ Myr} \simeq T_{\text{rot}}$$



X-type structure in edge-on view



Synthetic radio-maps of the simulated galaxy reproduce the main features observed in real galaxies

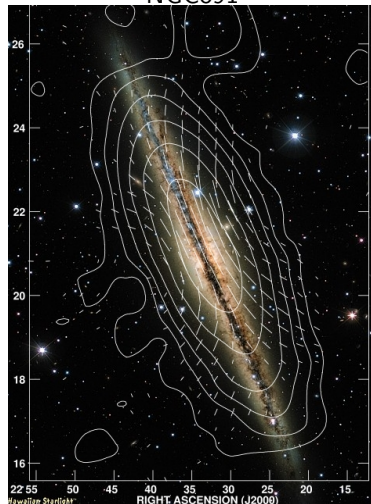
Spiral structure of magnetic field in the disk plain

M51

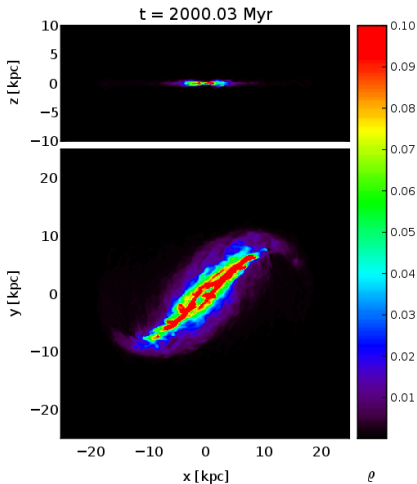
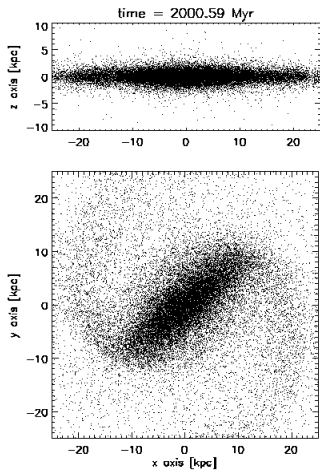


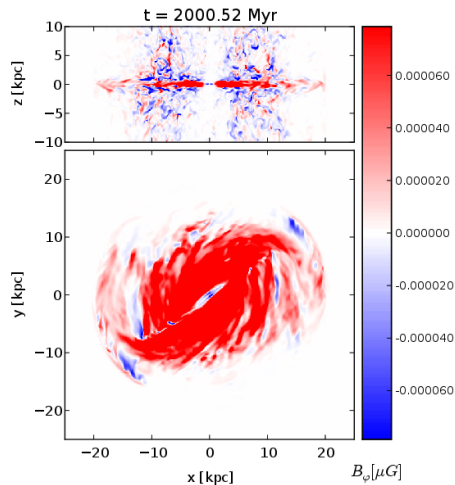
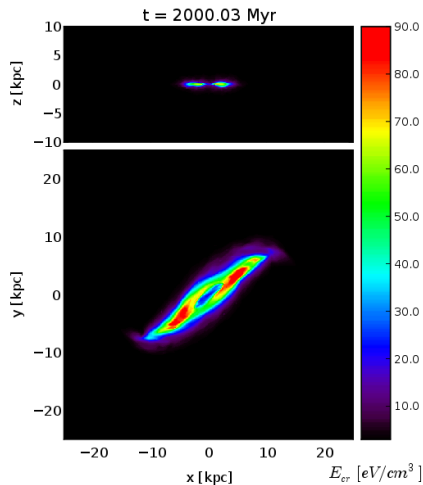
A. Fletcher et al. 2008

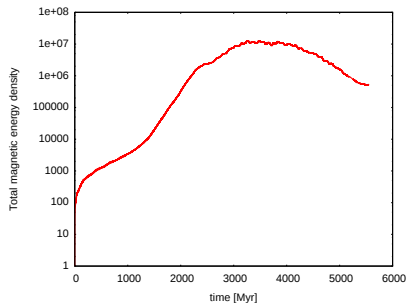
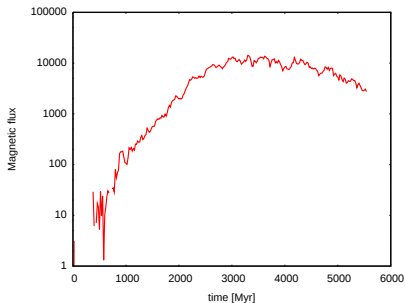
NGC891



M. Krause et al. 2008



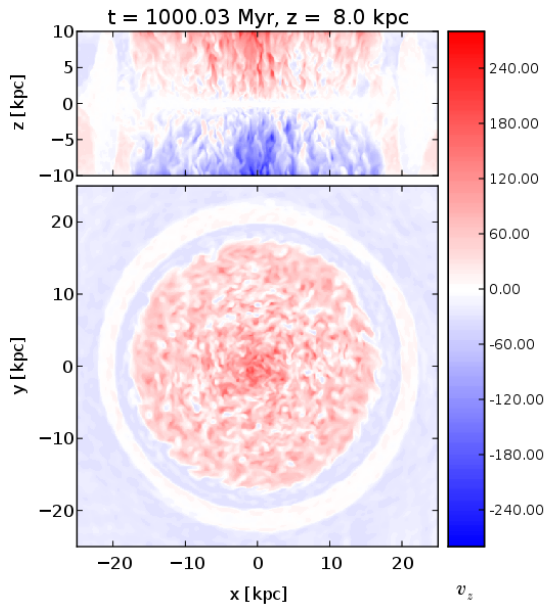




Magnetic field amplification stops too early

– due to the loss of angular momentum in the bar ?

WE NEED A BETTER N-BODY DISK



- COSMIC RAY DYNAMICS LEADS TO A VERY EFFICIENT MAGNETIC FIELD AMPLIFICATION IN GALACTIC DISKS
- Amplification timescale $\sim t_{rot} \Rightarrow$ growth of the large-scale magnetic field by several orders of magnitude, fast enough to expect $\sim 1\mu\text{G}$ magnetic field in galaxies at $z \sim 1 \div 2$
- Dipolar small-scale magnetic fields supplied by exploding stars build up a large scale magnetic field $\Rightarrow$ no need for seed fields of cosmological origin.
- Efficient regularization of the random magnetic field component.
- Growth of magnetic field, driven by SNe, far outside the star forming ring.
- Synthetic radio-maps resemble magnetic field structures in real galaxies.
- Simulations of the live disk need improvements