



Collisionless Accretion Disks: Role of Reconnection in Anisotropic Plasmas

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Collisionless Accretion Disk & MRI

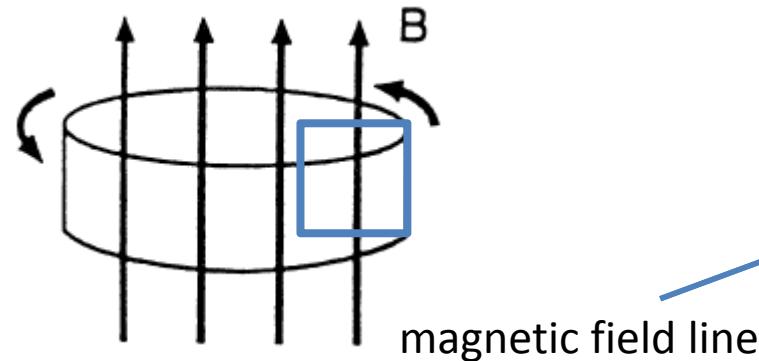
- ✓ Accretion Disk around Massive Black Holes (e.g., Sgr A*)
- ✓ Radiatively Inefficient Accretion Flow (RIAF)
- ✓ Nonthermal electrons, power law index $p \sim 1.3 < 2$

(e.g., Narayan+ 1998; Quataert+ 2002; Yuan+ 2003;
Aharonian+ 2008; Kusunose & Takahara 2012,...)

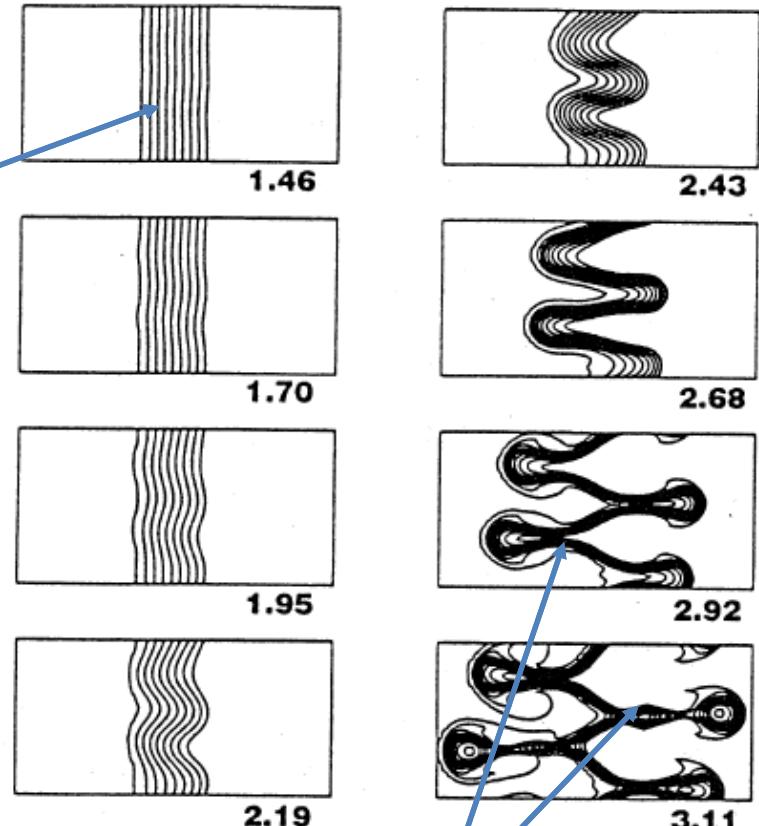
- ✓ MRI by using 2d & 3d particle-in-cell simulations
- ✓ Nonthermal particle acceleration with a hard spectrum
- ✓ Enhancement of Shakura-Sunyaev's “ α -parameter”, i.e.,
 $\alpha_{\text{collisionless}}/\alpha_{\text{MHD}} = O(10-100)$

Collisionless magnetic reconnection plays a dynamically important role on both particle acceleration and α -parameter

Role of Reconnection in MRI



Time evolution of MRI
Balbus and Hawley 1998

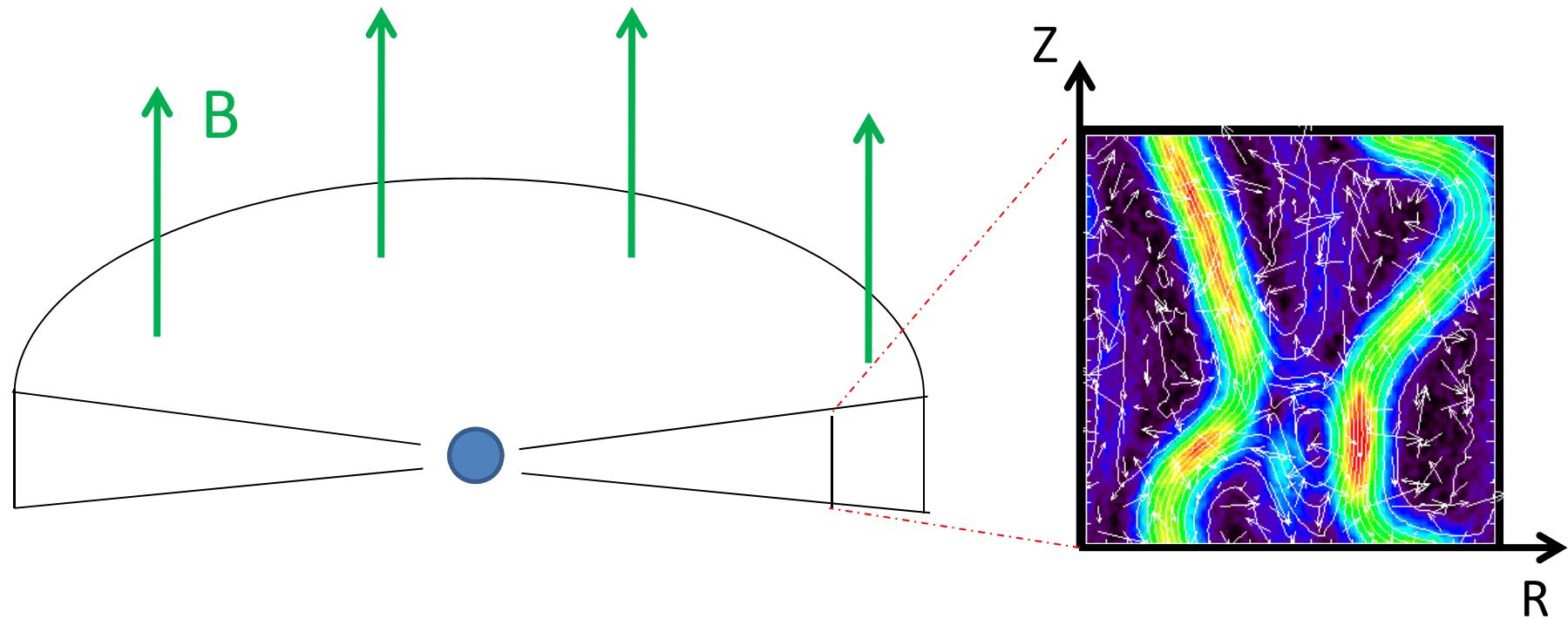


- ✓ Saturation of B-field in MRI is determined by balance between “B-field amplification by MRI” and “B-field dissipation by reconnection”
- ✓ Initial weak B-field ($\beta \gg 1$) $\rightarrow \beta \sim 10$

(e.g., Hawley+ 1995; Sano+ 2004...)

magnetic reconnection

Collisionless MRI in 2d PIC simulation

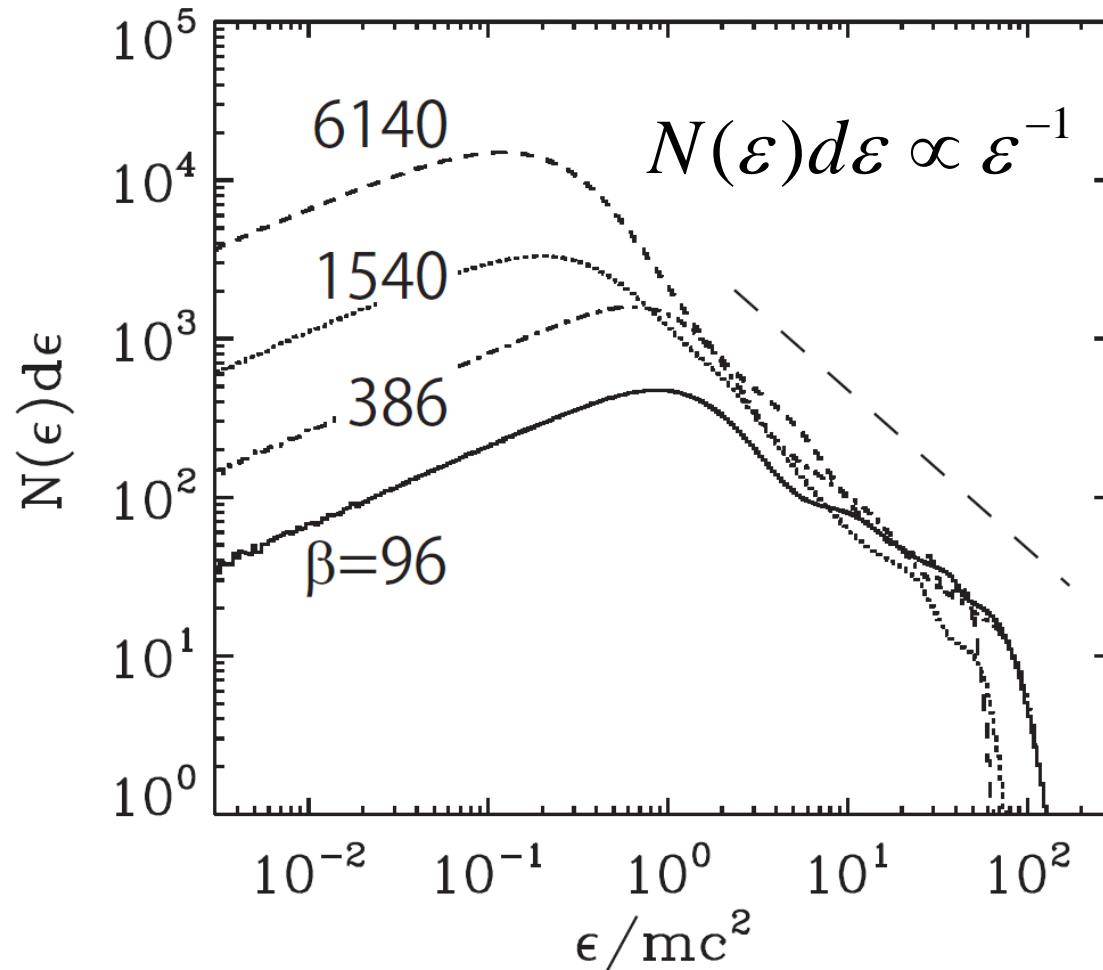


Kepler rotation Ω , $\Omega / \Omega_c = 0.1$,
 $\beta = 1540$, 200x200 grids 8000 particles/cell,
open shearing box boundary, electron-positron plasma

cf. Requelm + ApJ 2012; Shirakawa & MH PoP 2014

MH ApJ 2013

Energy Spectra during MRI-Reconnection

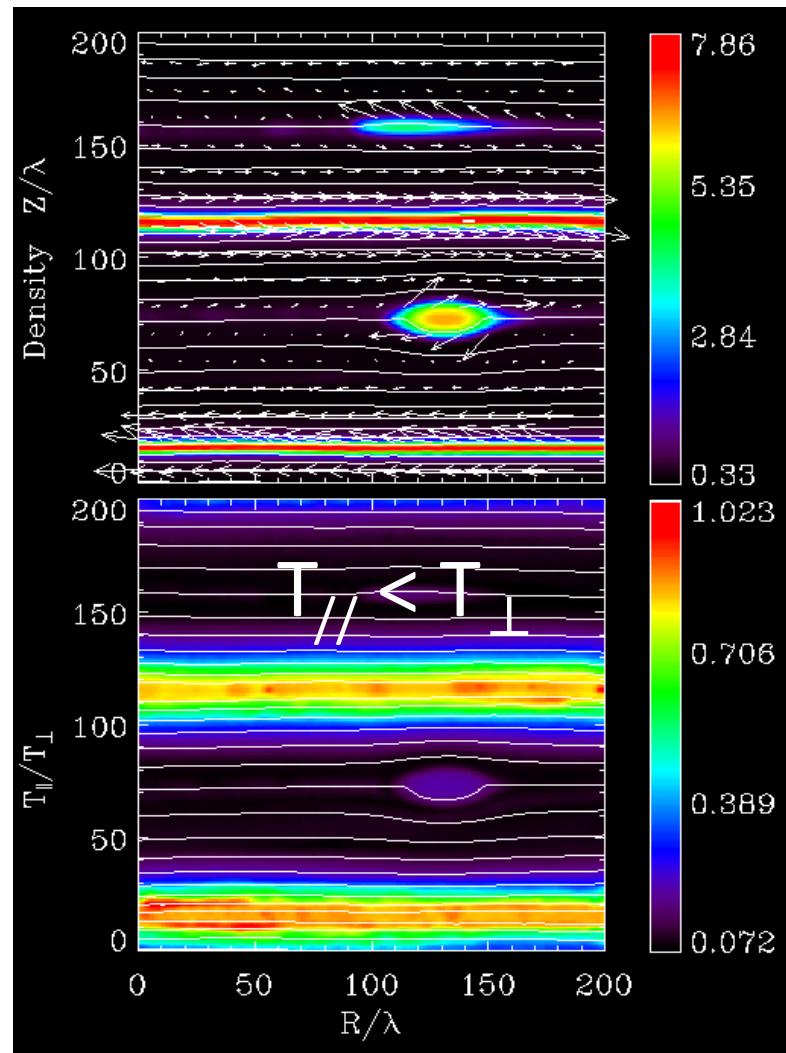


Dependence of initial plasma beta is weak

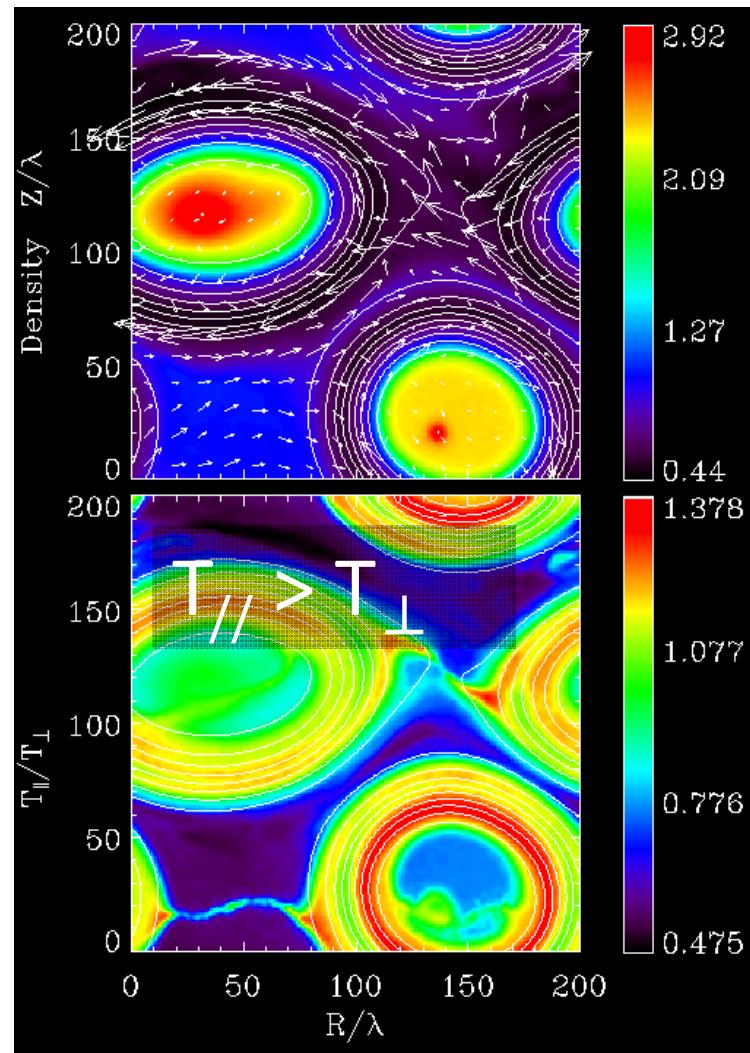
cf. Zenitani & MH ApJ 2001, reconnection talks in this meeting

Onset of Reconnection

Before onset of Reconnection



After onset of Reconnection



Production of Pressure Anisotropy during MRI & Reconnection

CGL (Chew-Goldberger-Low) or
Double adiabatic theory

$$\frac{D}{Dt} \left(\frac{p_{\perp}}{\rho B} \right) = 0, \quad \frac{D}{Dt} \left(\frac{p_{\parallel} B^2}{\rho^3} \right) = 0$$

MRI: B large $\Rightarrow p_{\perp} > p_{\parallel}$

Istropization by mirror inst. & ion-cyclotron inst.

(Quataert+ ApJ 2002; Sharma+ ApJ 2006)

Reconnection : B weak $\Rightarrow p_{\parallel} > p_{\perp}$

Istropization by Alfvén waves generated by ion beam inst. etc.

(e.g. MH+ 1998; Higashimori & MH 2015)

Role of Reconnection in Collisionless MRI

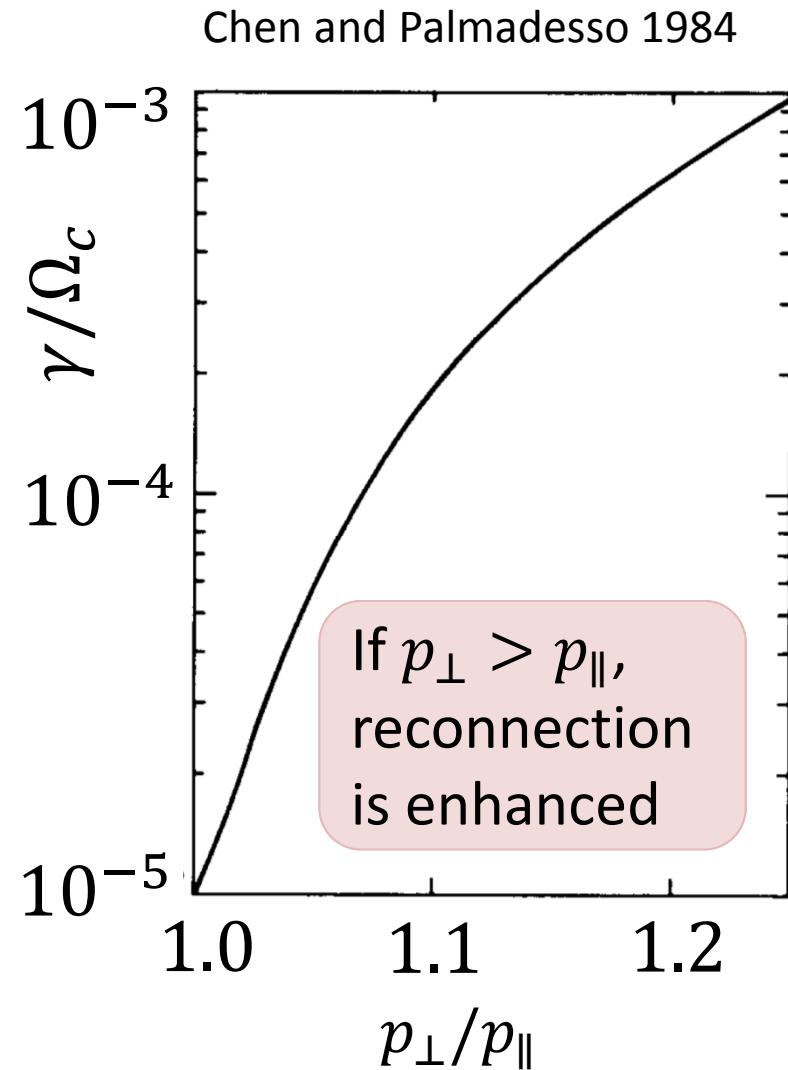
- ✓ Dynamic evolution of reconnection is strongly controlled by pressure anisotropy

(e.g., Laval & Pellat 1968; Chen & Palmadesso 1984; MH 1987)

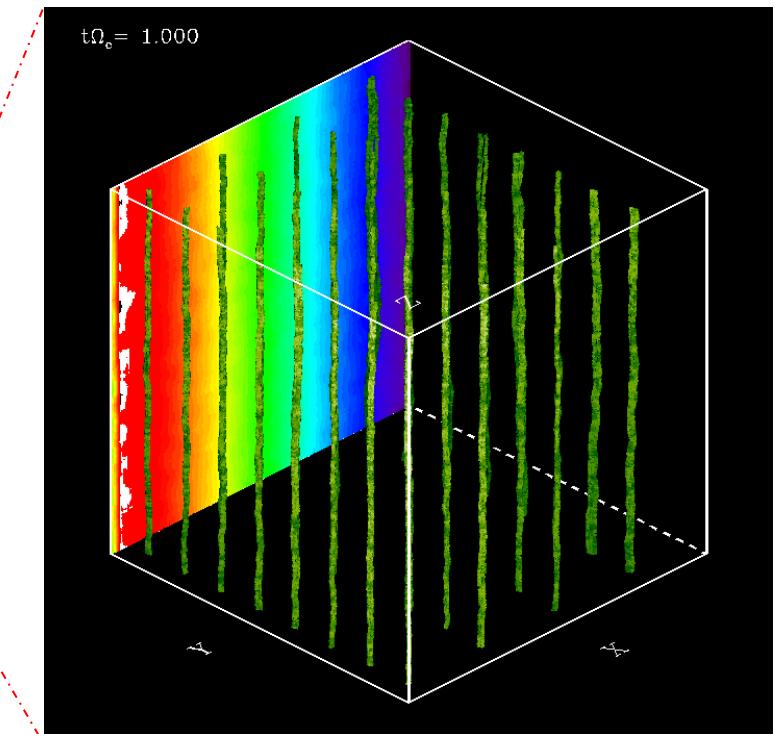
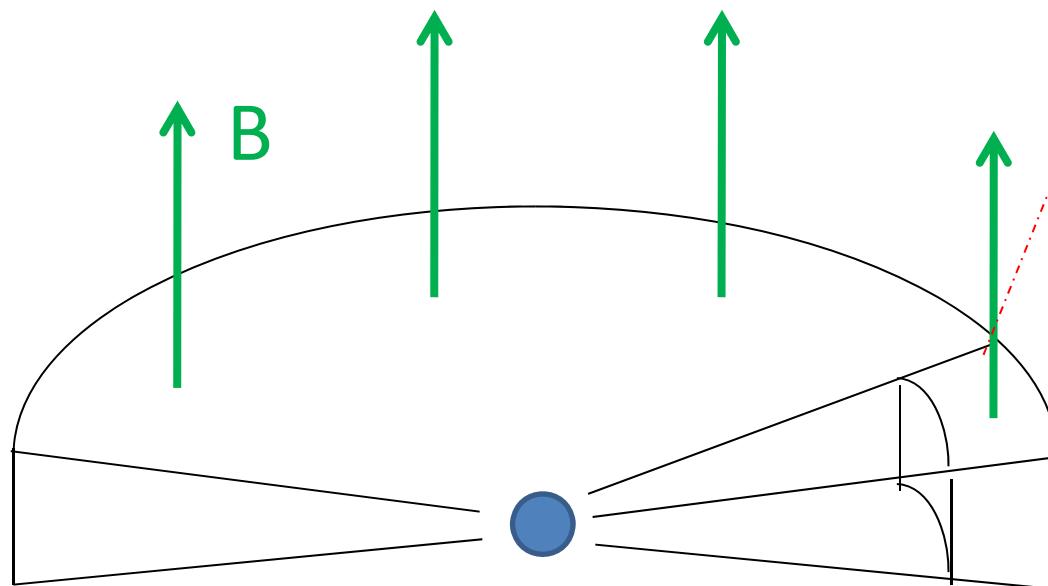
Tearing theory

$$\frac{\gamma}{kv_{th}} \approx \left(\frac{p_{\perp}}{p_{\parallel}} - 1 \right) + \left(\frac{r_g}{\delta} \right)^{3/2} \left(\frac{1 - k^2 \delta^2}{k \delta} \right)$$

If $p_{\parallel} > p_{\perp}$,
reconnection is suppressed



Angular Momentum Transport in 3d PIC simulation

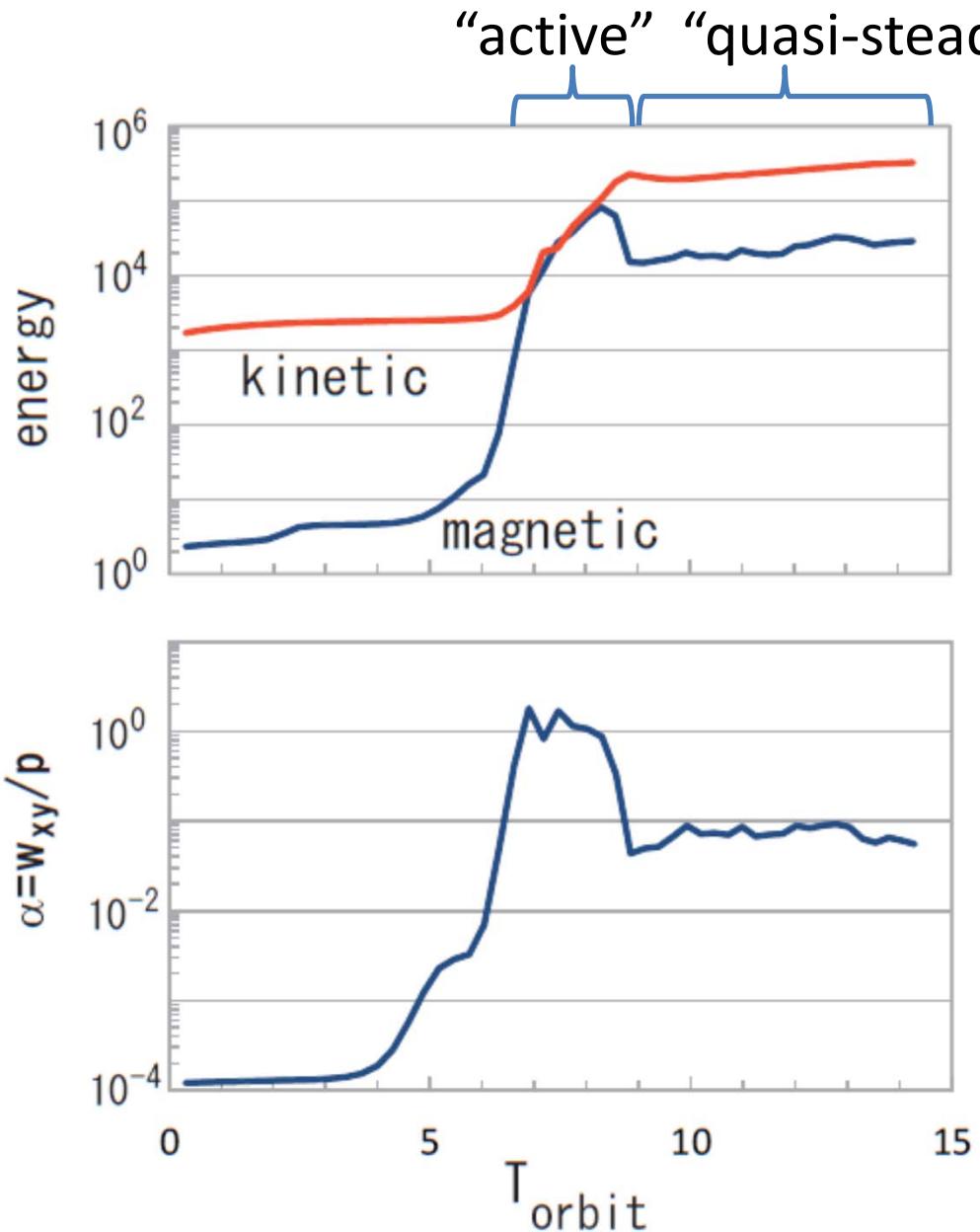


$\beta=1540$, Kepler rotation Ω
300x300x300 grids 40 particles/cell,
open shearing box boundary,
electron-positron plasma

green: magnetic field lines
color contour: angular velocity

MH PRL 2015

Energy and Stress Tensor Evolutions



Initial plasma $\beta = 1540$,
active phase $\beta \sim O(1)$
quasi-steady-state $\beta \sim O(10)$

stress tensors

$$w_{xy} = \rho v_x \delta v_y - \frac{B_x B_y}{4\pi} + \frac{(p_{\parallel} - p_{\perp})}{B^2} B_x B_y$$

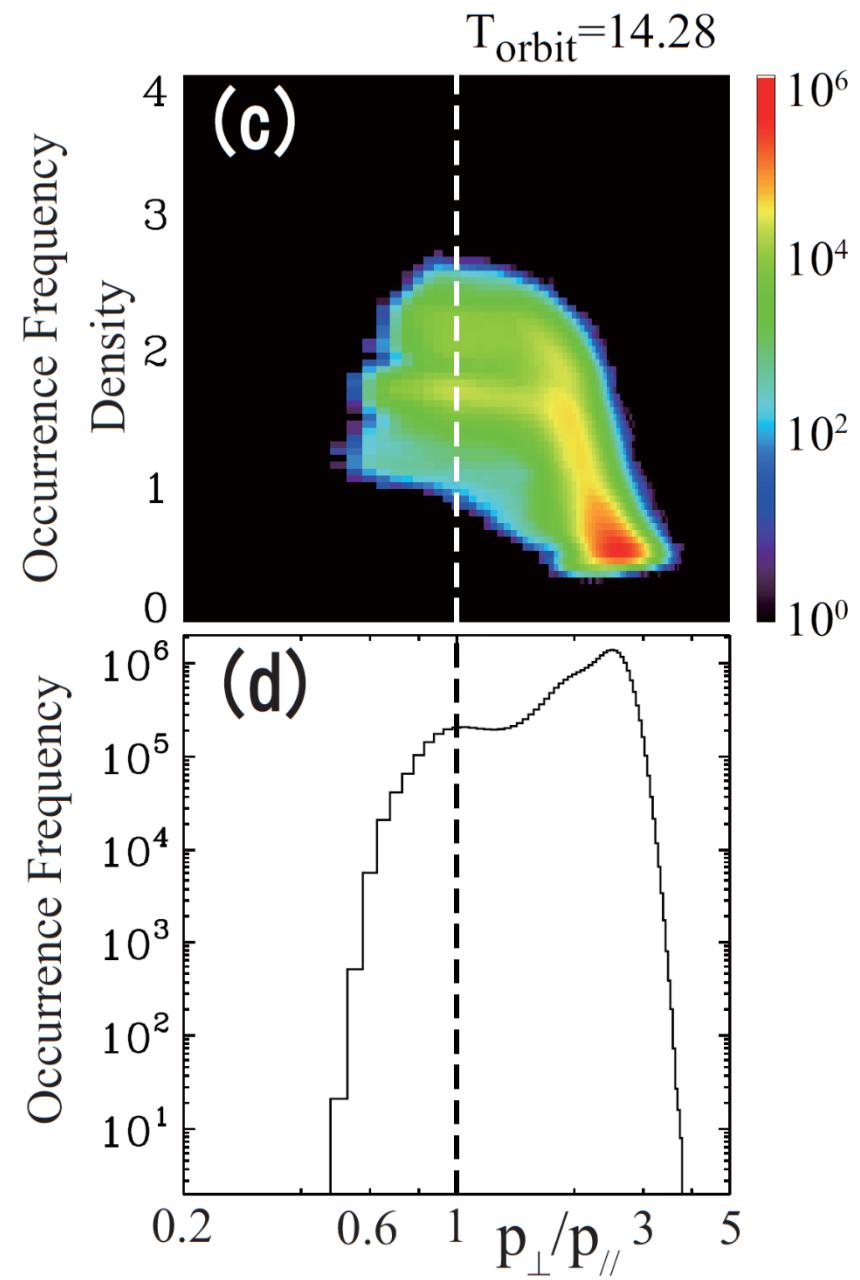
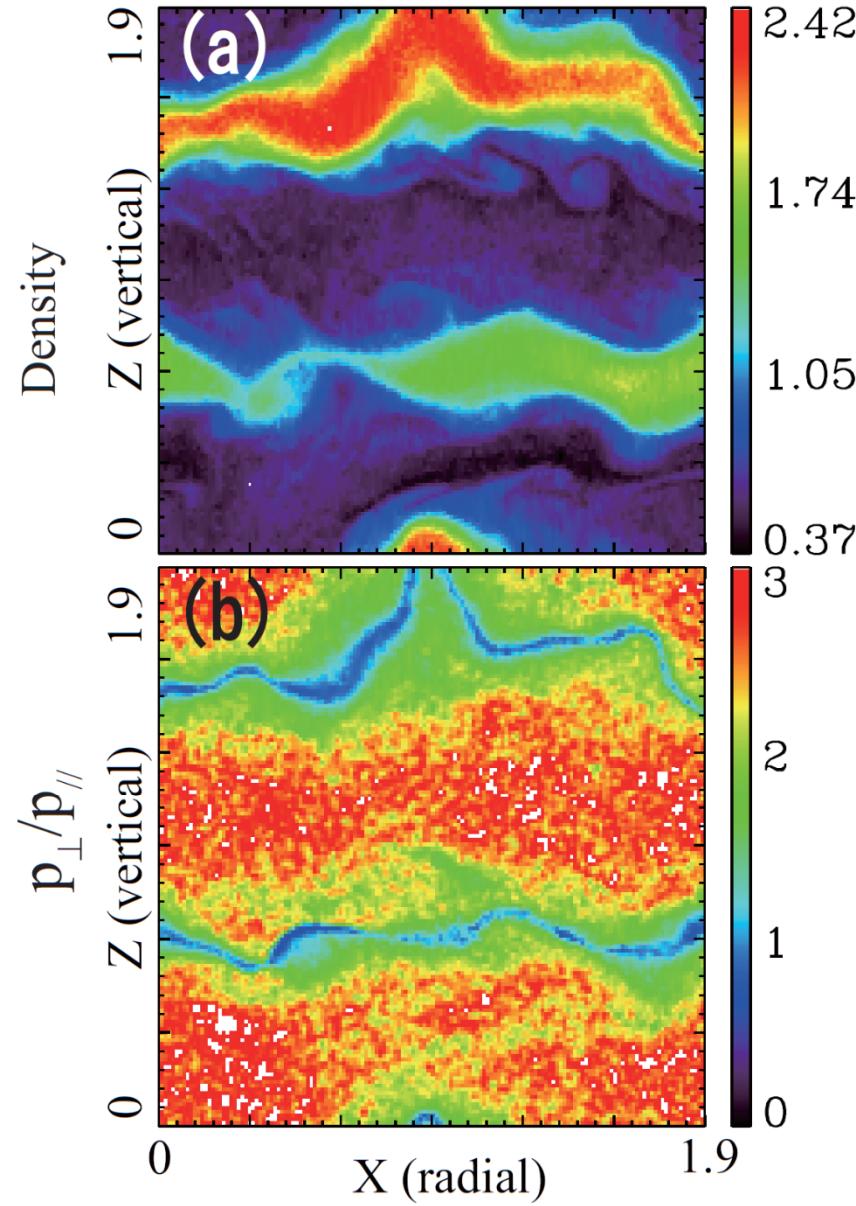
$$\begin{aligned} \alpha &= \frac{w_{xy}}{p} \approx - \frac{B_x B_y}{4\pi p} \\ &= - \frac{2B_x B_y}{B^2} \frac{B^2}{8\pi p} \approx \frac{B^2}{8\pi p} = \frac{1}{\beta} \end{aligned}$$

$$\alpha(\text{kinetic}) \sim O(0.1)$$

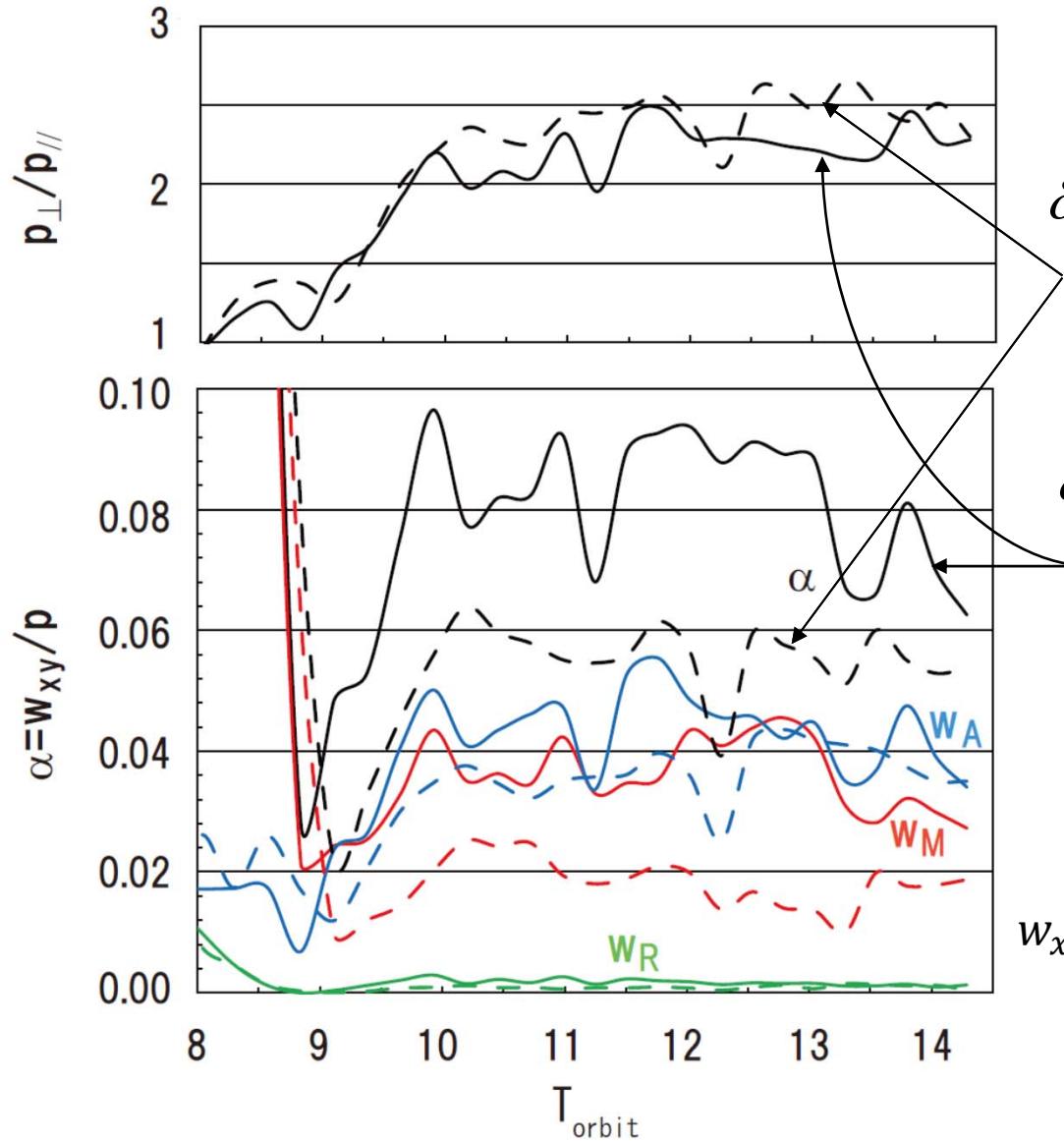
$$\alpha(\text{kinetic})/\alpha(\text{MHD}) > 10 - 100$$

(e.g., Hawley+ 1995; Sano+ 2004...)

Reconnection is suppressed by $p_{\parallel} > p_{\perp}$



$$\frac{d\vec{p}}{dt} = e \left(\vec{E} + \frac{\vec{v}}{c} \times (\vec{B} + \delta\vec{B}) \right) - m\gamma \left(2\Omega_0 \times \vec{v} - 2q\Omega_0^2 x \hat{x} \right)$$



$\delta\vec{B}$: random field in channel flow

With pitch-angle
scattering model

$\delta\vec{B} = 0$

Without pitch-angle
scattering model

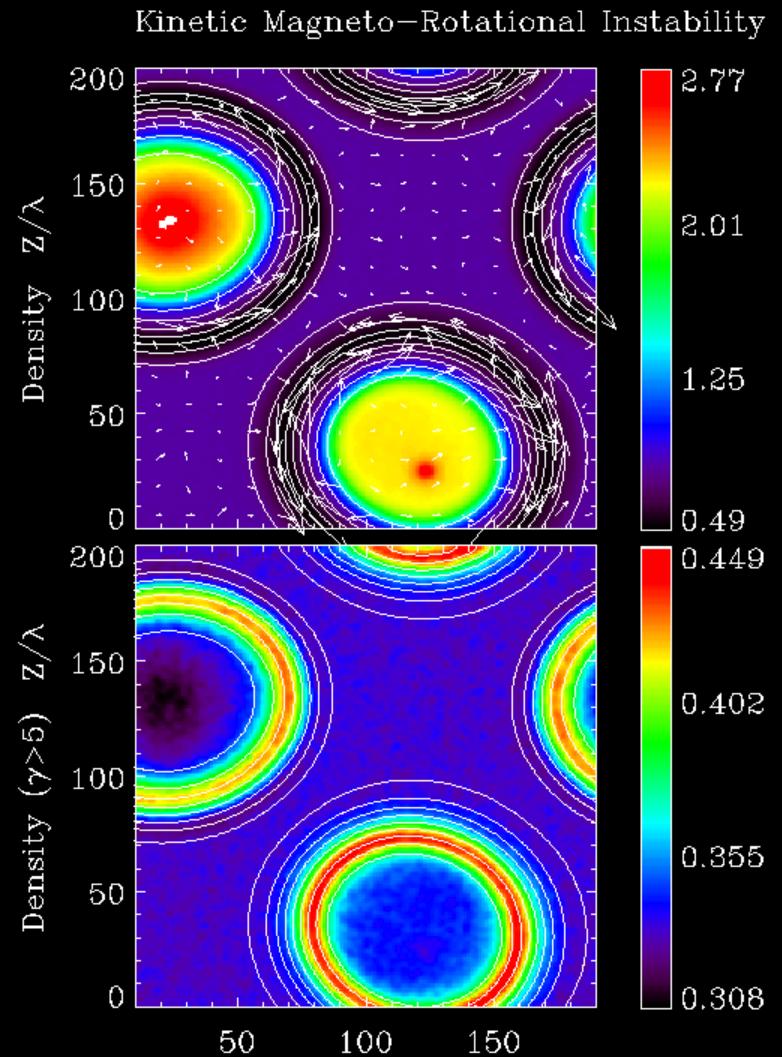
$$w_{xy} = \rho v_x \delta v_y - \frac{B_x B_y}{4\pi} + \frac{(p_{\parallel} - p_{\perp})}{B^2} B_x B_y$$

w_R

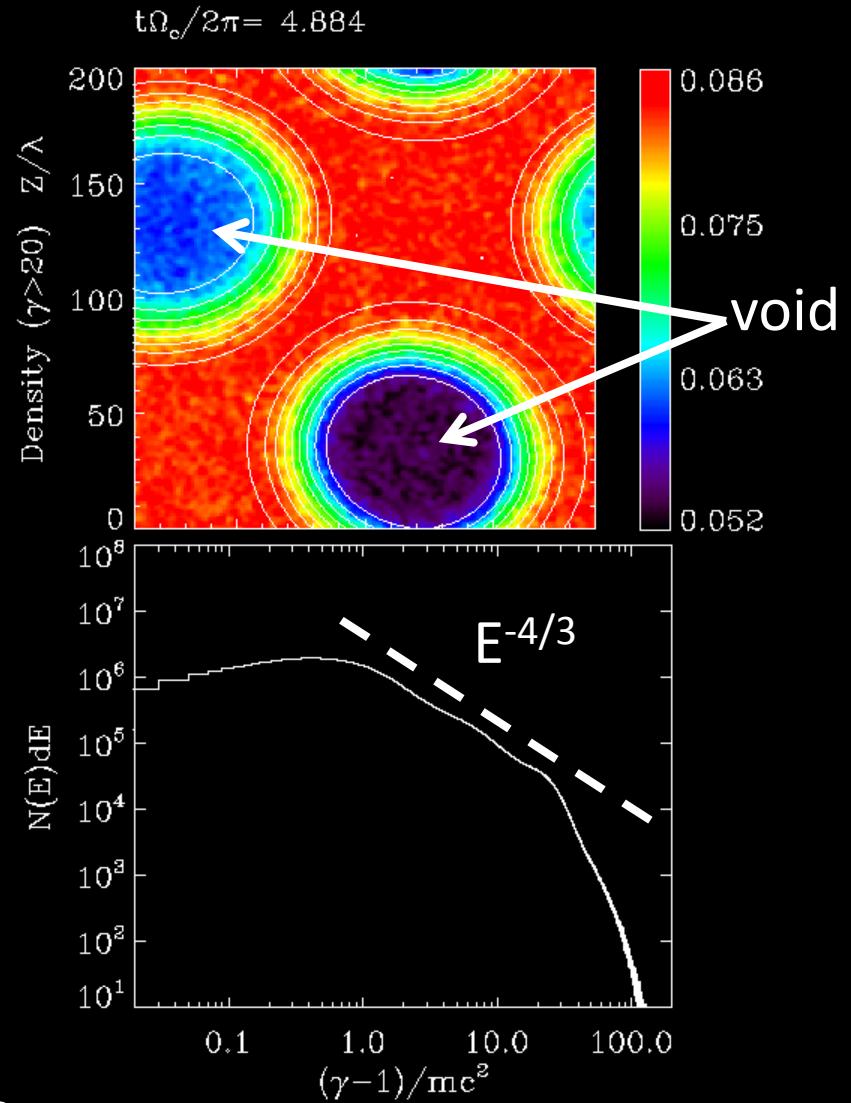
w_M

w_A

thermal plasma is confined
inside magnetic islands



high energetic particles ($\gamma>20$) are
located outside magnetic islands

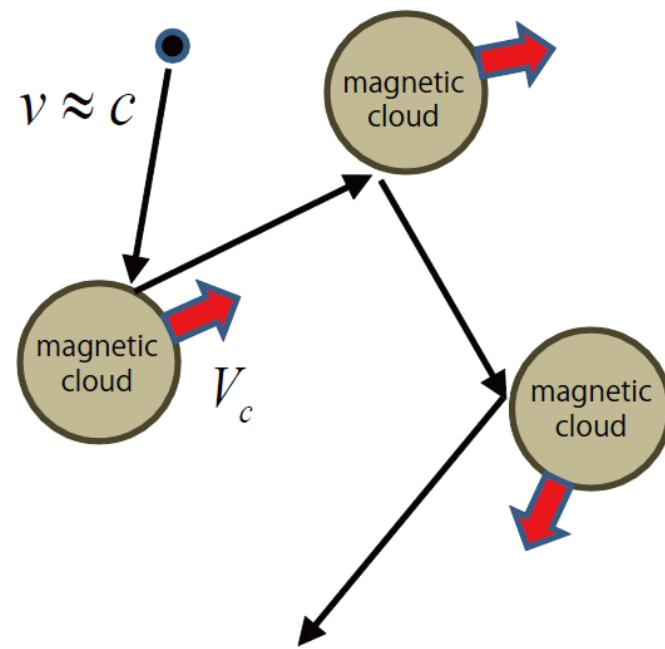


middle energetic particles ($\gamma>5$) are
located at outer edge of islands

Fermi-Reconnection Acceleration in Many Magnetic Islands

2nd order Acceleration

cosmic ray
(energetic particle)

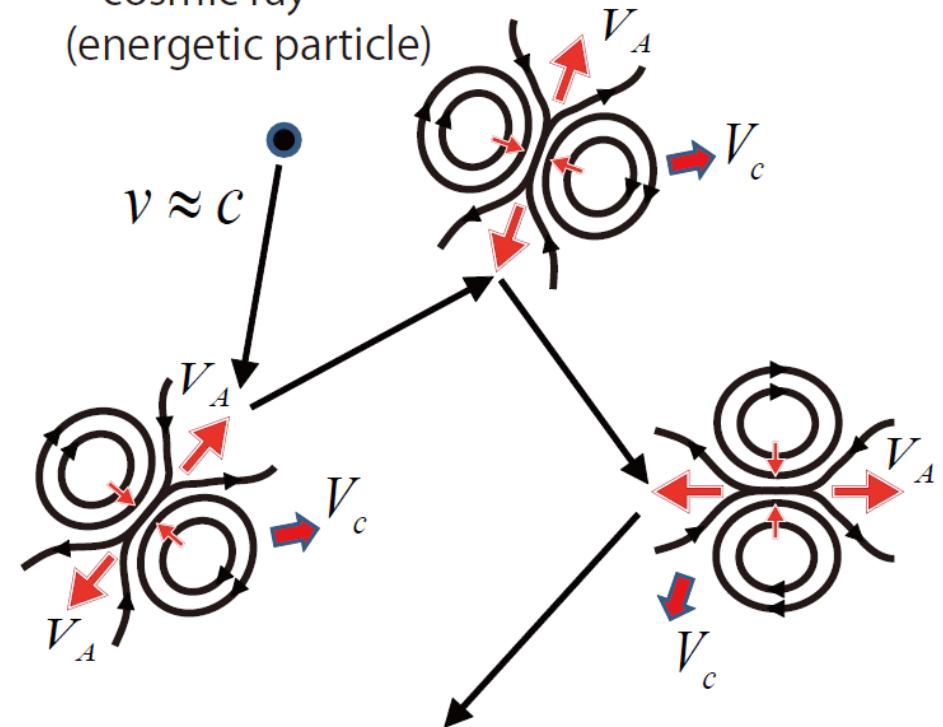


$$\frac{\Delta\epsilon}{\epsilon} \approx \left(\frac{V_c}{c} \right)^2$$

Fermi, Phys. Rev. (1949)

1st order Acceleration

cosmic ray
(energetic particle)



$$\frac{\Delta\epsilon}{\epsilon} \approx \left(\frac{V_A}{c} \right)$$

MH PRL (2012); MH & Lyubarsky SSR (2013)

Summary

"Particle Acceleration" and "Angular Momentum Transport" during MRI in Collisionless Accretion Disk

- $T_{\perp} > T_{\parallel\parallel}$ during MRI evolution
- Enhanced reconnection rate due to $T_{\perp} > T_{\parallel\parallel}$
- Strong particle acceleration by reconnection

- $T_{\perp} < T_{\parallel\parallel}$ in current sheet by reconnection
- Suppressed onset of reconnection due to $T_{\perp} < T_{\parallel\parallel}$
- Large B field and enhanced "a parameter"