

The most powerful persistent engine of Nature

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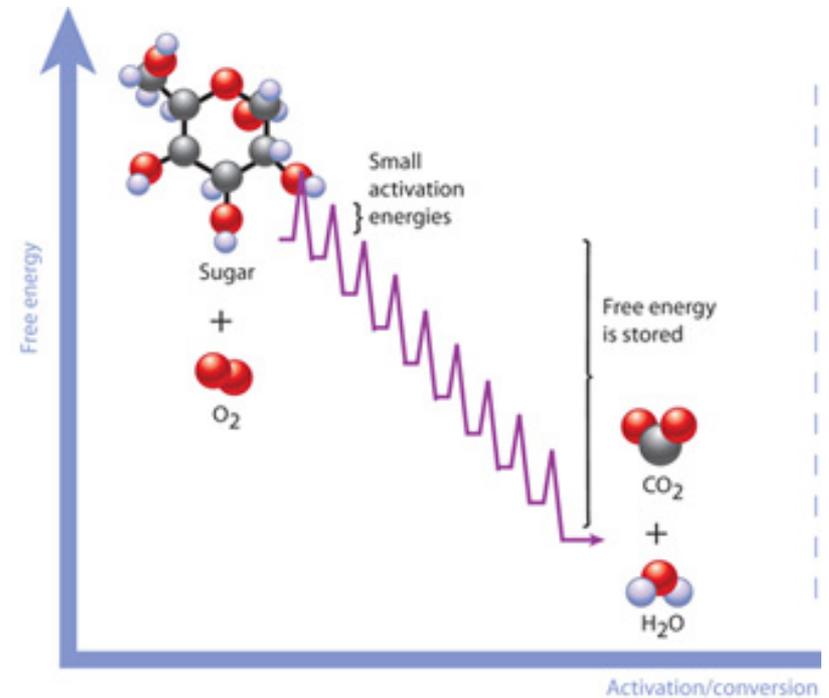
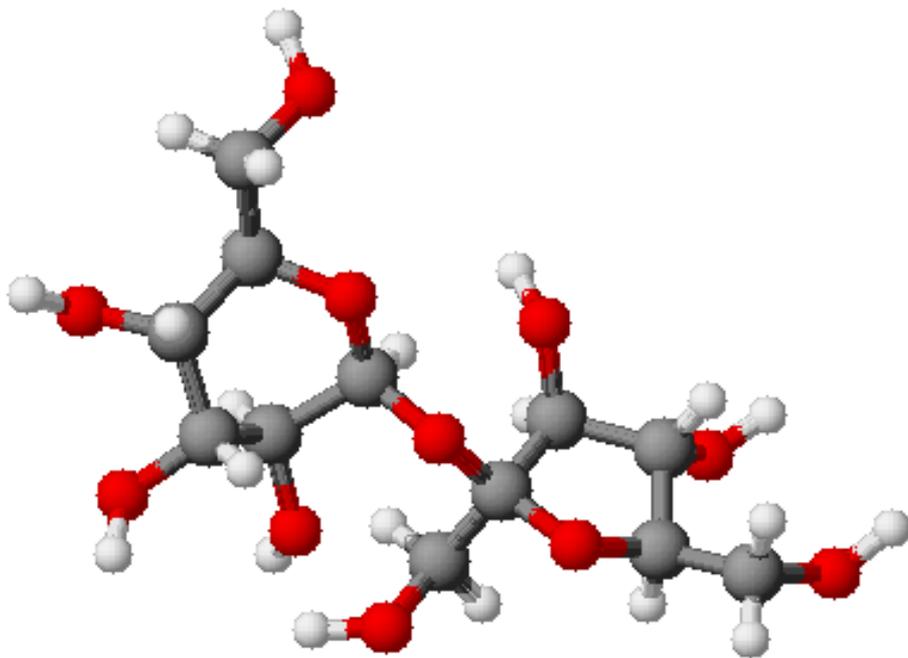
With:

Fabrizio Tavecchio, Laura Maraschi,
Annalisa Celotti, Tullia Sbarrato

Efficienc

y

$$\eta = \frac{E}{mc^2}$$



Sugar saccharose



1g \square 4 kcal = 16.2 kJ = 1e23 eV = 4.7 eV per molecule

$$\eta = \frac{E}{mc^2} = \frac{1.6 \times 10^{11} \text{ erg}}{9 \times 10^{20} \text{ erg}} = 1.8 \times 10^{-10}$$

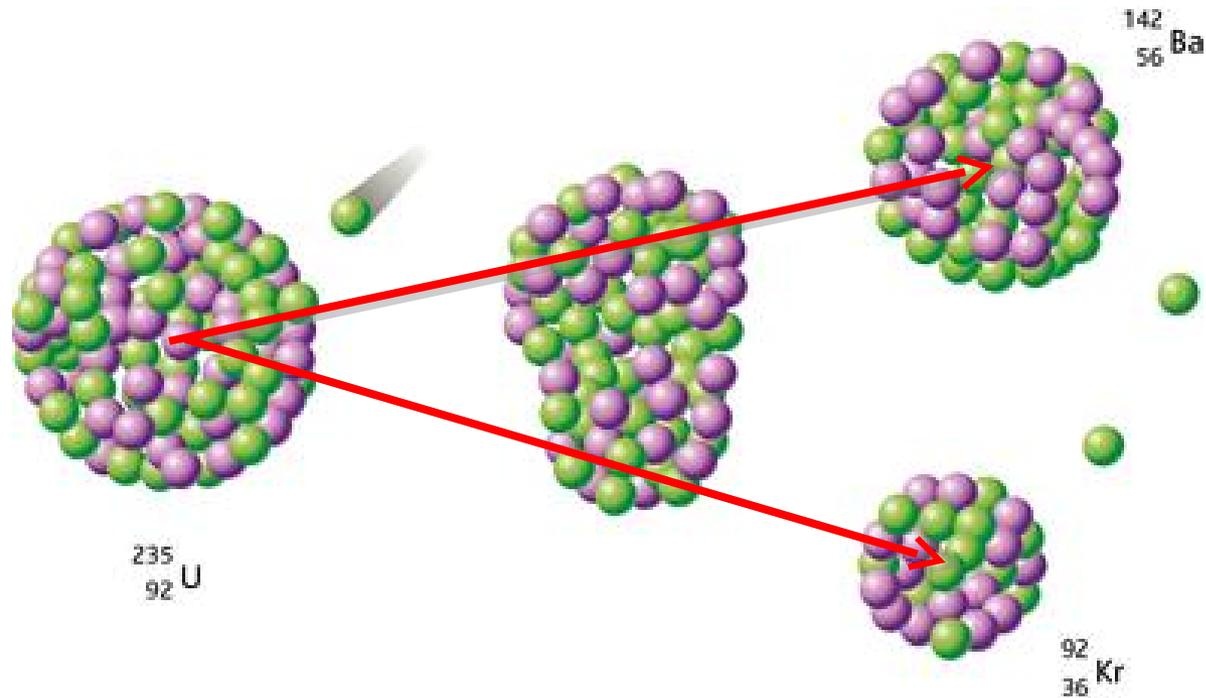


$$\eta = \frac{mg}{h mc^2} = \frac{980 \times 10^4 (h/100)}{9 \times 10^{20} \text{ erg}} \sim 10^{-14}$$



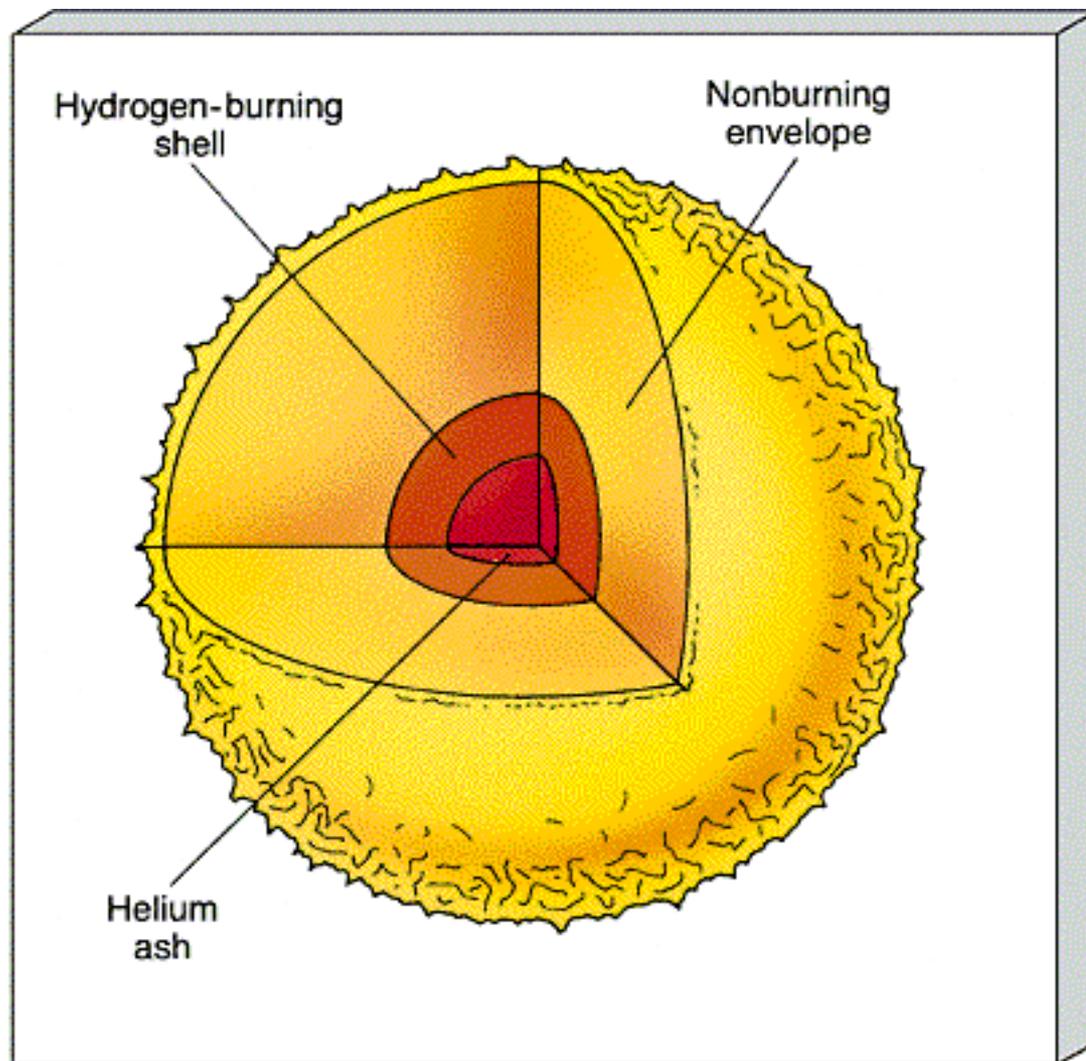
$$\eta = \frac{mv^2}{2mc^2} = \frac{(v/c)^2}{2} = 4 \times 10^{-15} (v/100 \text{ km/h})^2$$

Nuclear fission

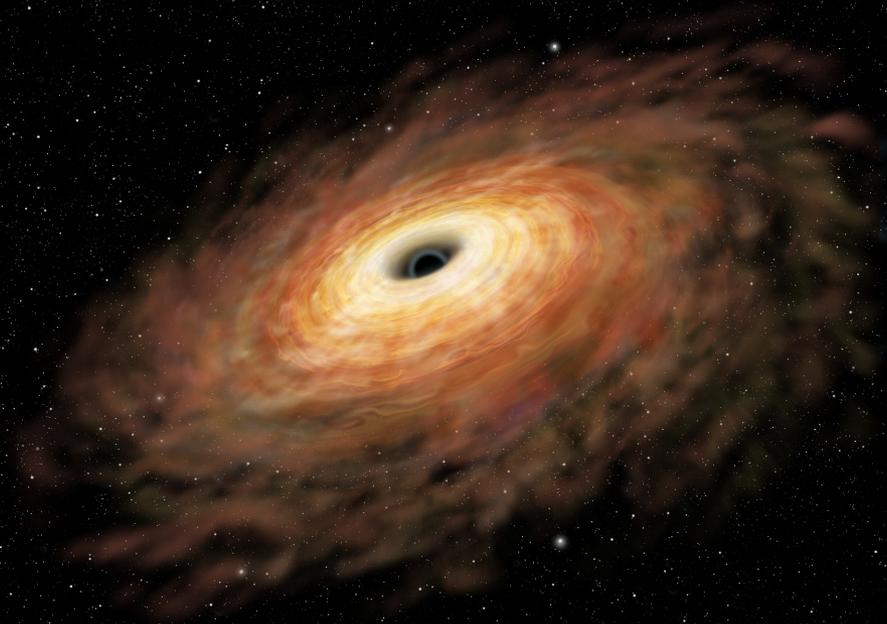


$$\eta = \frac{E}{mc^2} = \frac{0.2 \times 10^9 \text{ eV}}{235 \times 9.4 \times 10^8 \text{ eV}} \sim 9 \times 10^{-4}$$

Fusion



$$\eta = 0.008 \times 0.1 \sim 8 \times 10^{-4}$$

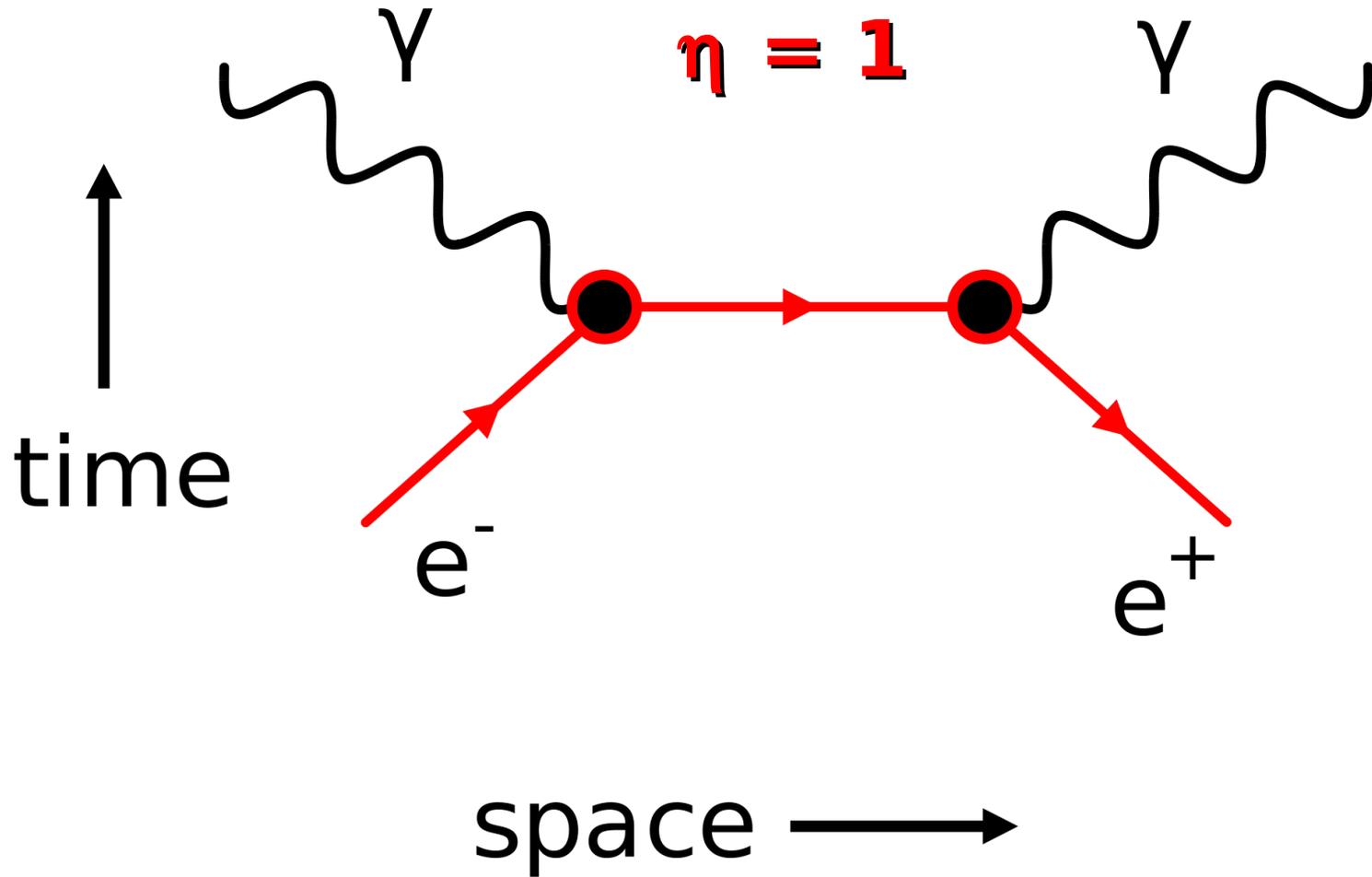


$$\eta = \frac{1}{2} \frac{GM}{R} \frac{m}{mc^2} = \frac{R_g}{2R} \quad \text{(Newton)}$$

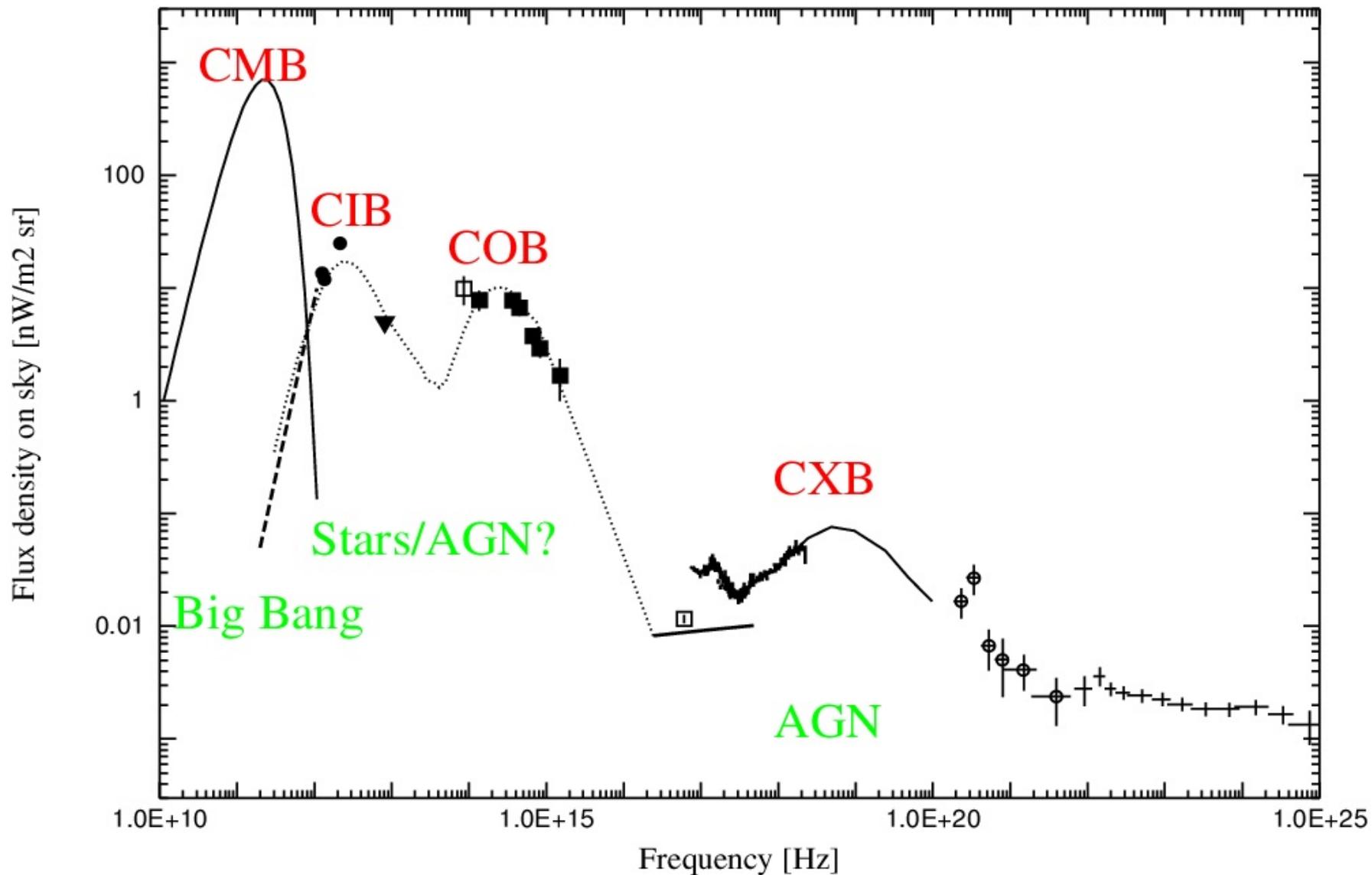
$R_{\min} = R_g$ for max spin

$\eta = 0.1$ up to 0.3 for accreting
Kerr (Thorne 1974)

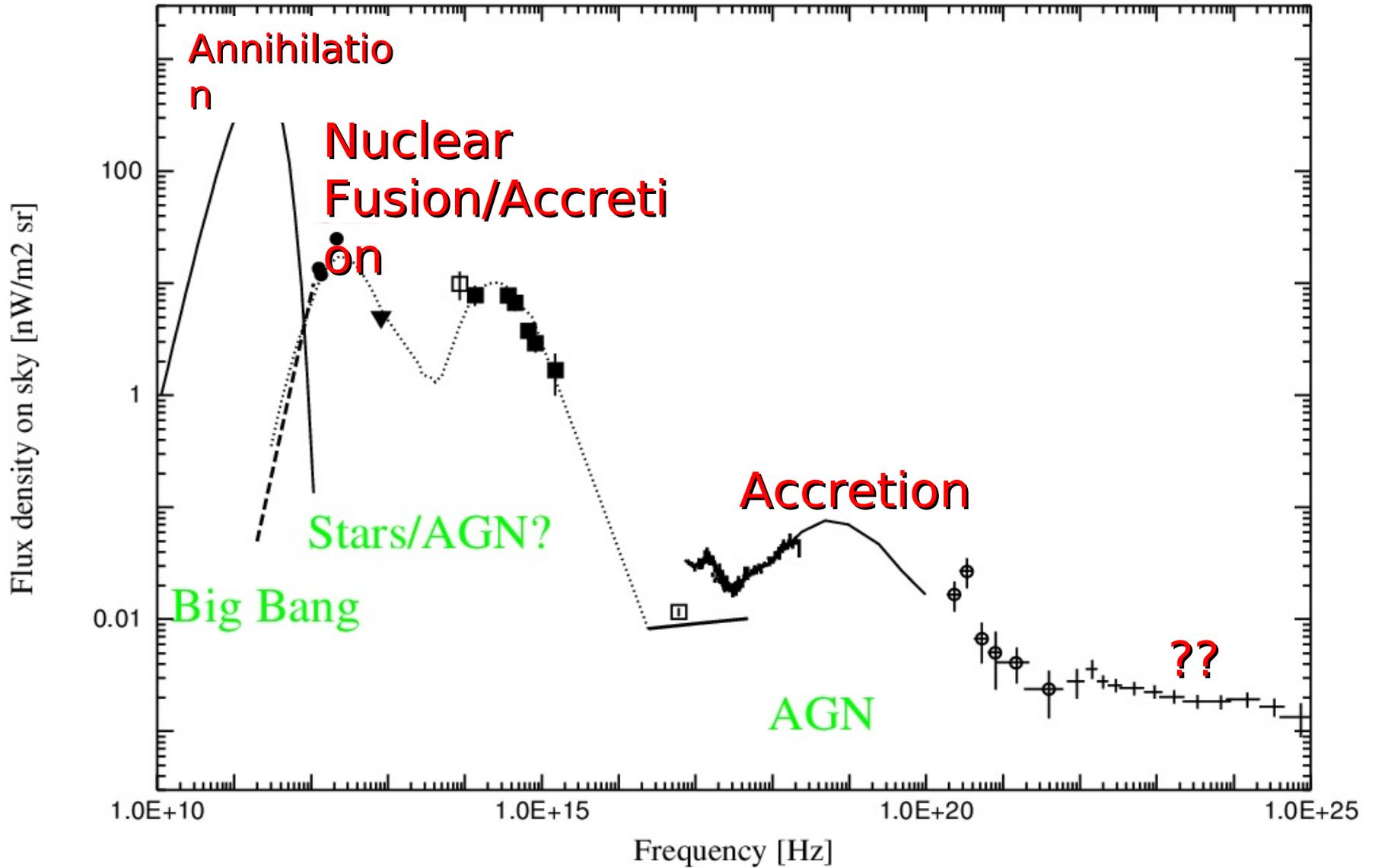
Annihilation



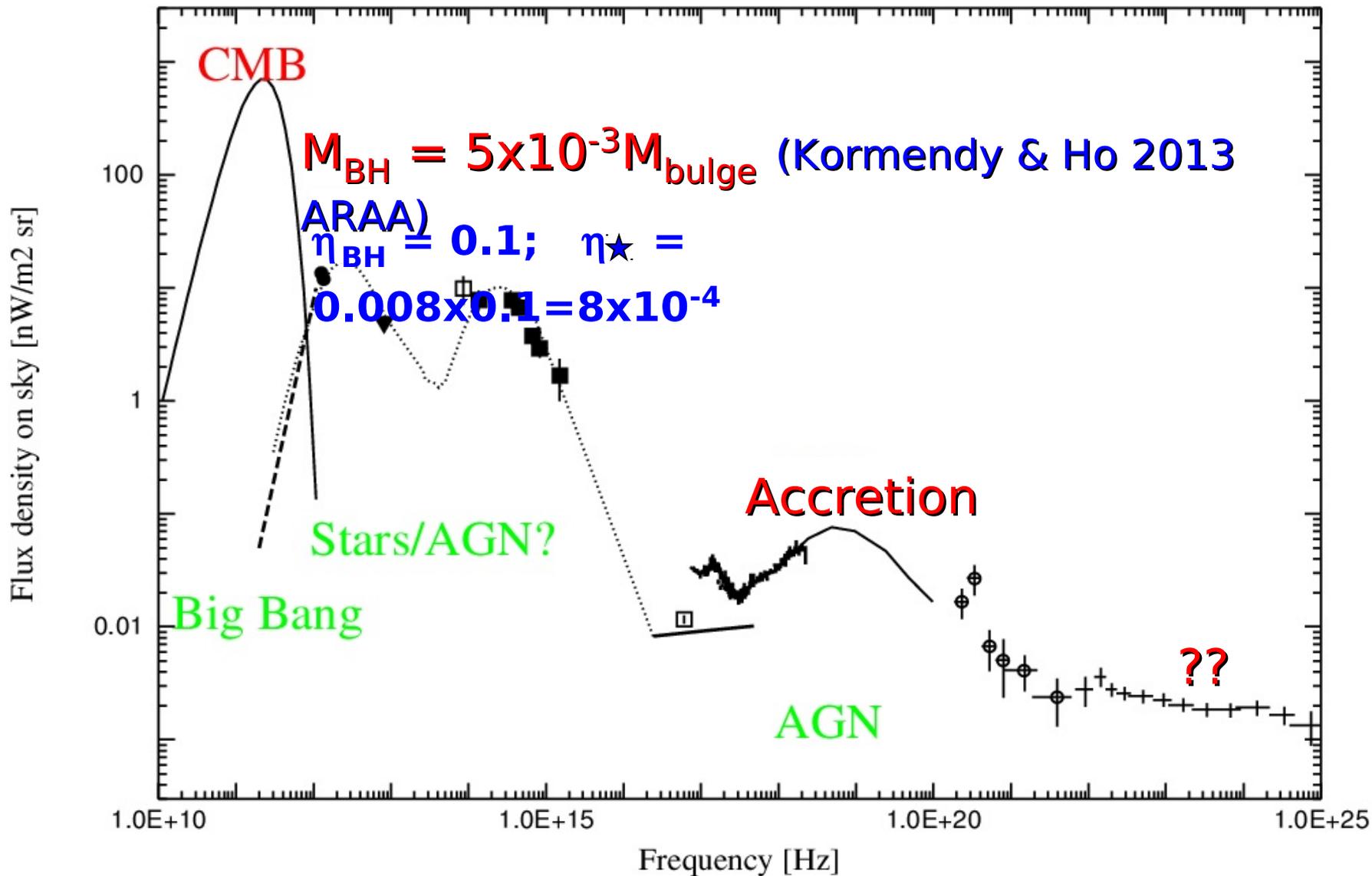
The Cosmic Energy Density Spectrum



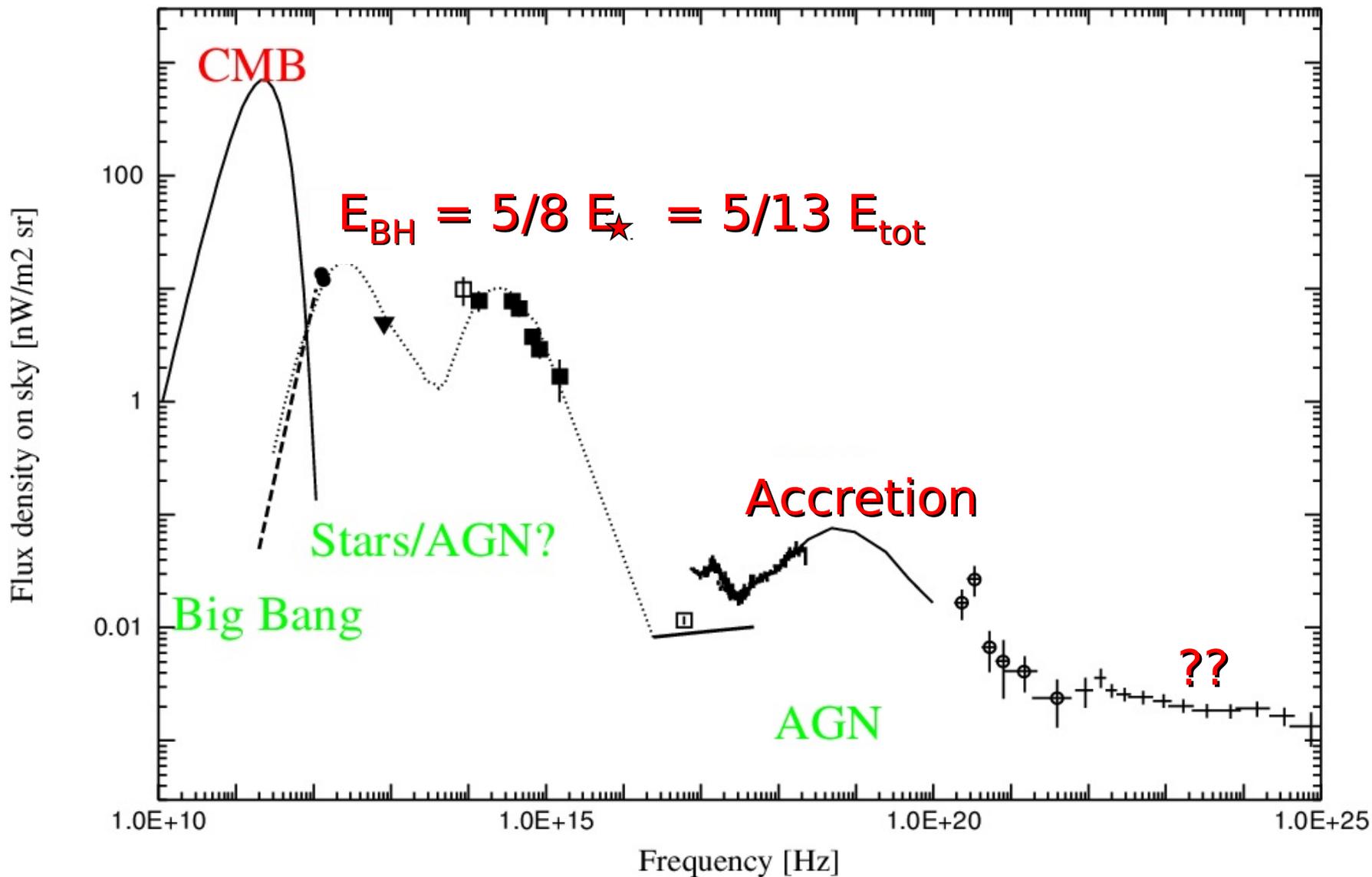
The Cosmic Energy Density Spectrum



The Cosmic Energy Density Spectrum

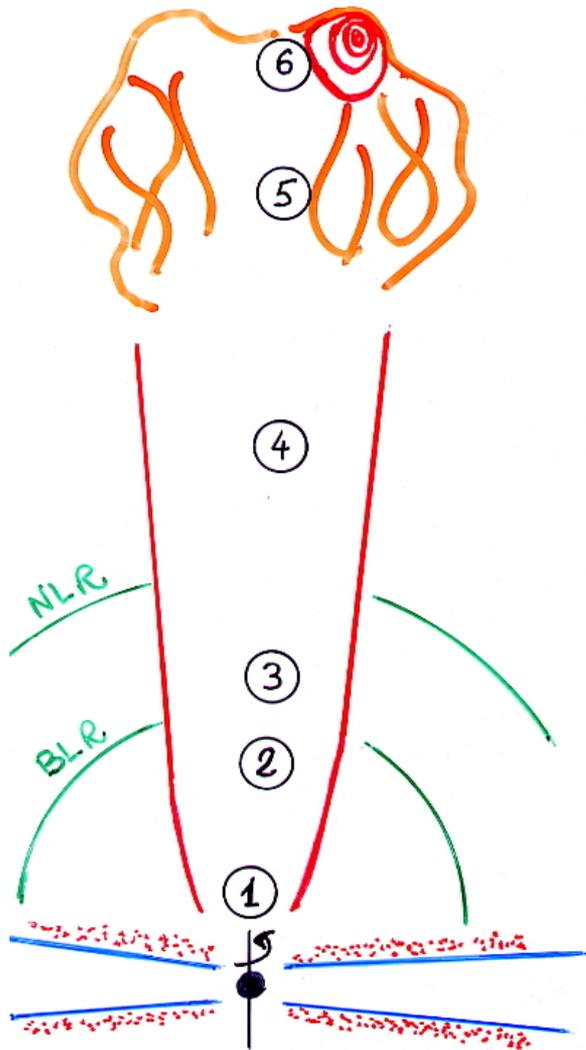


The Cosmic Energy Density Spectrum



Jets?

Power of jets in blazars



Hot spot

10^6 pc

Radio lobe

X-ray cavities

CHANDRA

10^5 pc

VLBI

10 pc

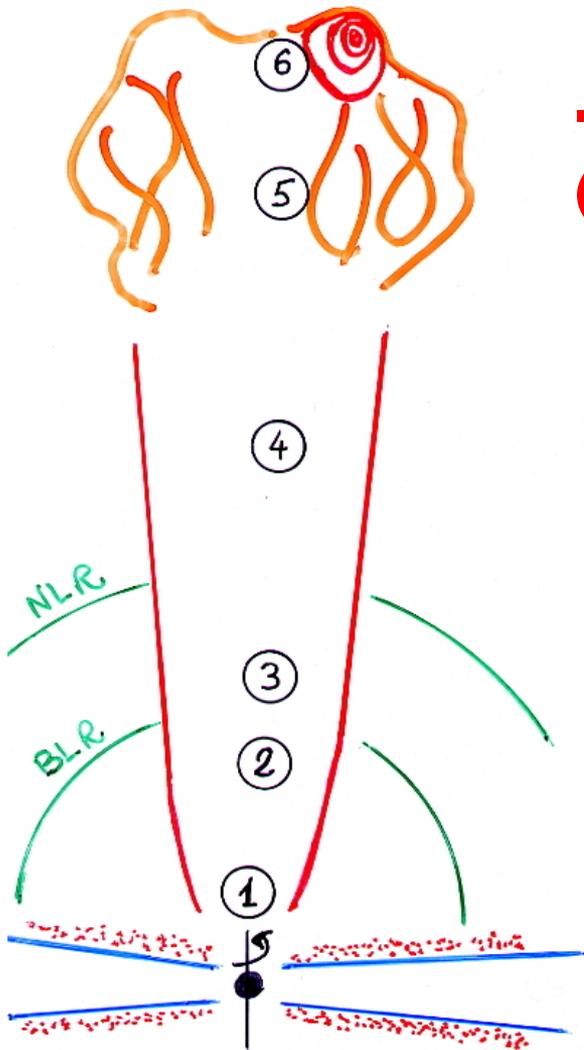
γ -ray zone

10^{-1} pc

Formation

10^{-4} pc

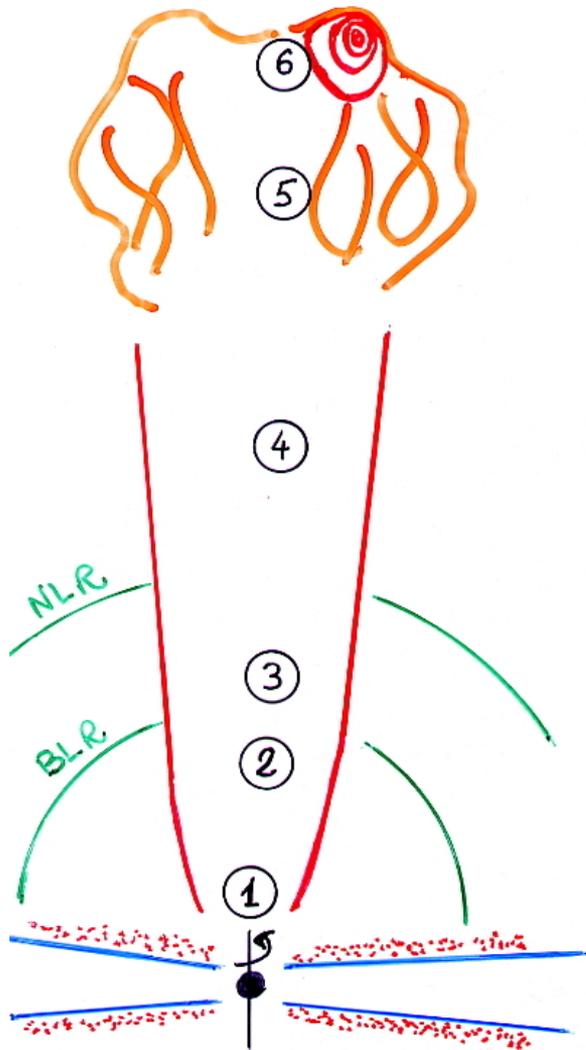
Power of jets in blazars



Hot spot	
Radio lobe	10^6 pc
X-ray cavities	
CHANDRA	10^5 pc
VLBI	10 pc
γ -ray zone	10^{-1} pc
Formation	10^{-4} pc

Minimum energy

t_{lobe}

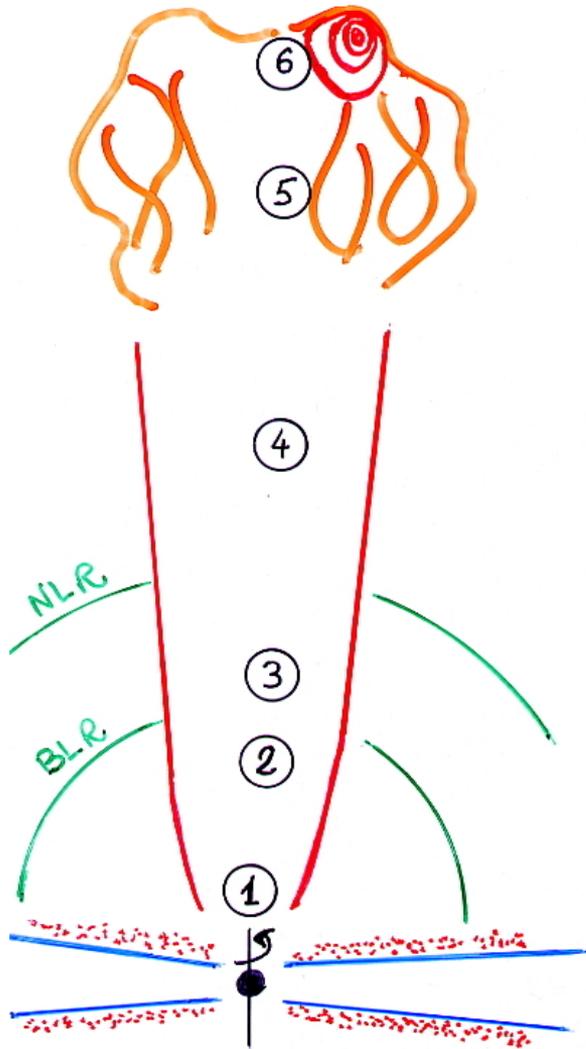


- Hot spot
- Radio lobe 10^6 pc
- X-ray cavities
- CHANDRA 10^5 pc
- VLBI 10 pc
- γ -ray zone 10^{-1} pc
- Formation 10^{-4} pc

Power of jets in blazars

PV work / t_{sound}

Power of jets in blazars



Hot spot

10^6 pc

Radio lobe

X-ray cavities

CHANDRA

10^5 pc

VLBI

10 pc

γ -ray zone

10^{-1} pc

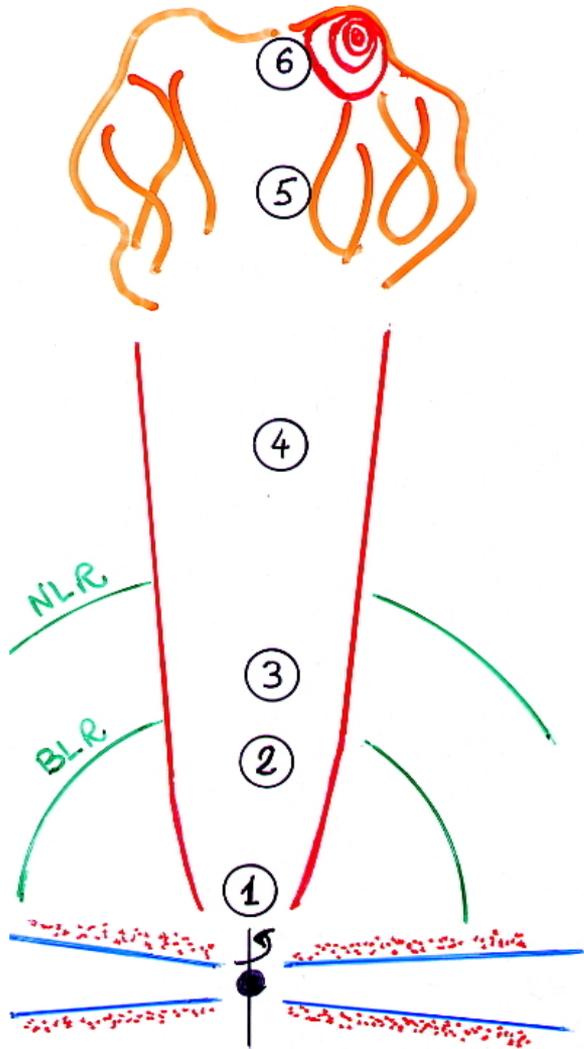
Formation

10^{-4} pc

Measure the size, apply SSC model, find δ and # of e^-

**1992: launch of EGRET on the
Compton Gamma-ray
Observatory**

Power of jets in blazars



Hot spot

10^6 pc

Radio lobe

X-ray cavities

CHANDRA

10^5 pc

VLBI

10 pc

γ -ray zone

10^{-1} pc

Formation

10^{-4} pc

jet power and accretion

luminosity

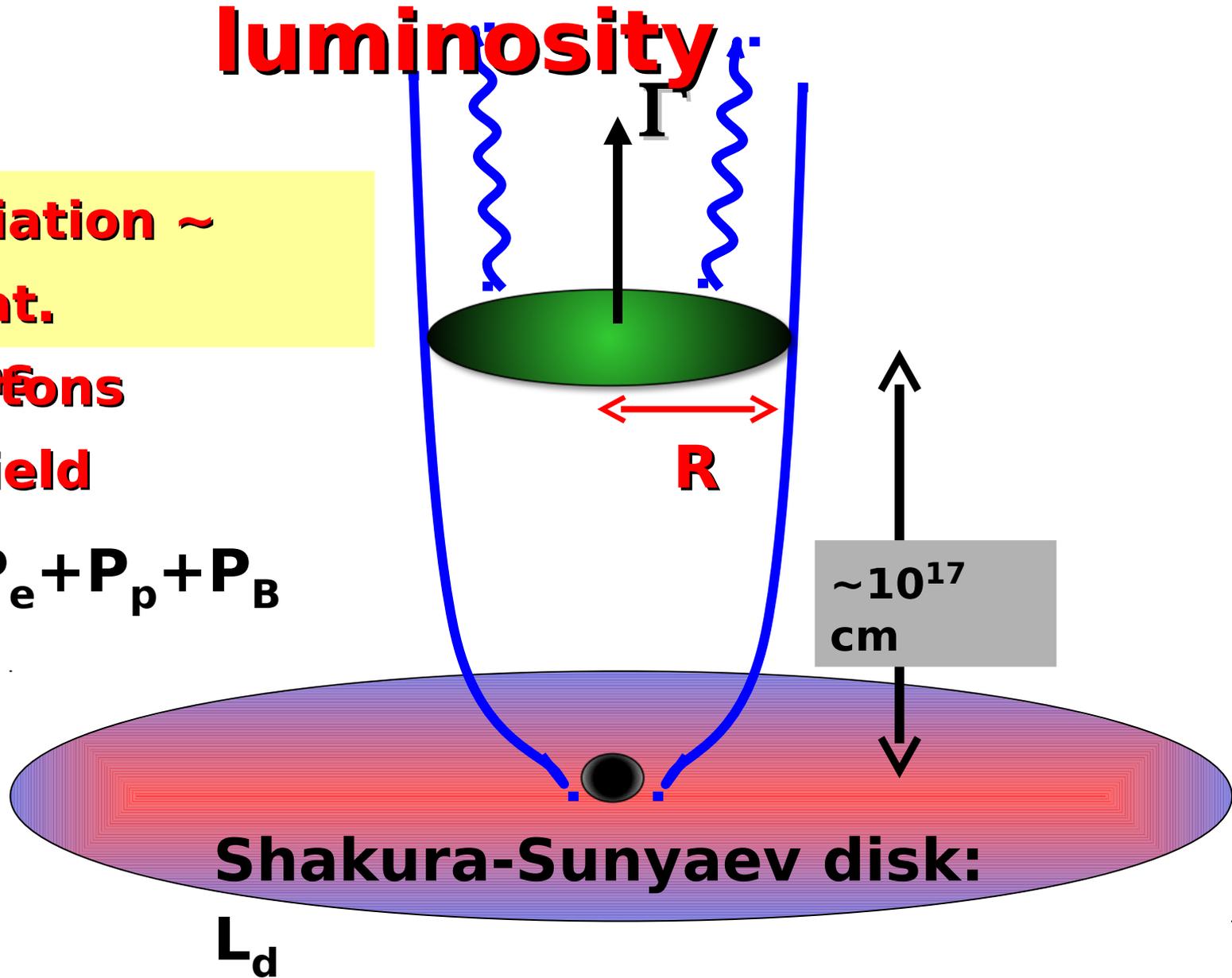
$P_r = \text{radiation} \sim$

$P_{\text{ebs}} = \Gamma^2 \text{relat.}$

$P_e = \text{electrons}$
 $P_p = \text{protons}$

$P_B = \text{B-field}$

$$P_{\text{jet}} = P_e + P_p + P_B$$



**If you want to compare disk
luminosity and jet power,
the best sample is:**

Blazars detected

by Fermi \square $L_{\text{obs, jet}}$

with broad emission lines \square

L_{disk}

Shaw+ 2012: FSRQs (~220 sources)

Shaw+ 2013: BL Lacs (26 with BLR)

The jet cannot have less power than what required to produce the observed luminosity:

$$P_{\text{jet}} > \frac{L_{\text{obs}}}{\Gamma^2}$$

If P_{jet} is twice as much, Γ halves.

We can take P_{rad} as the minimum P_{jet} .

This limit is model-independent.

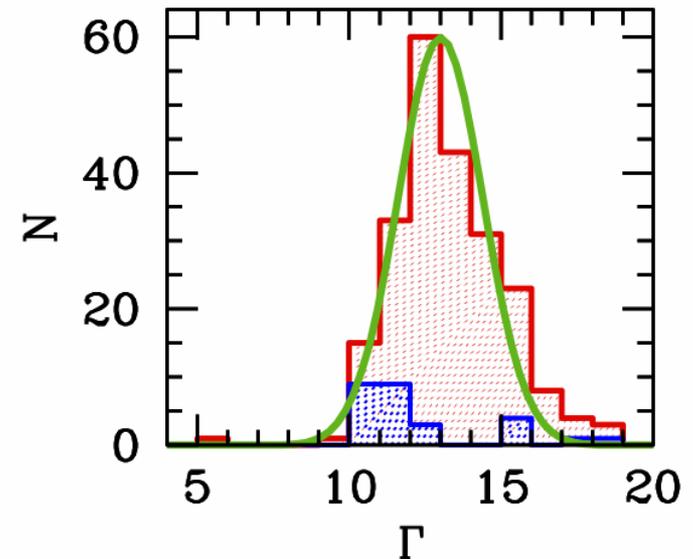
The jet cannot have less power than what required to produce the observed luminosity:

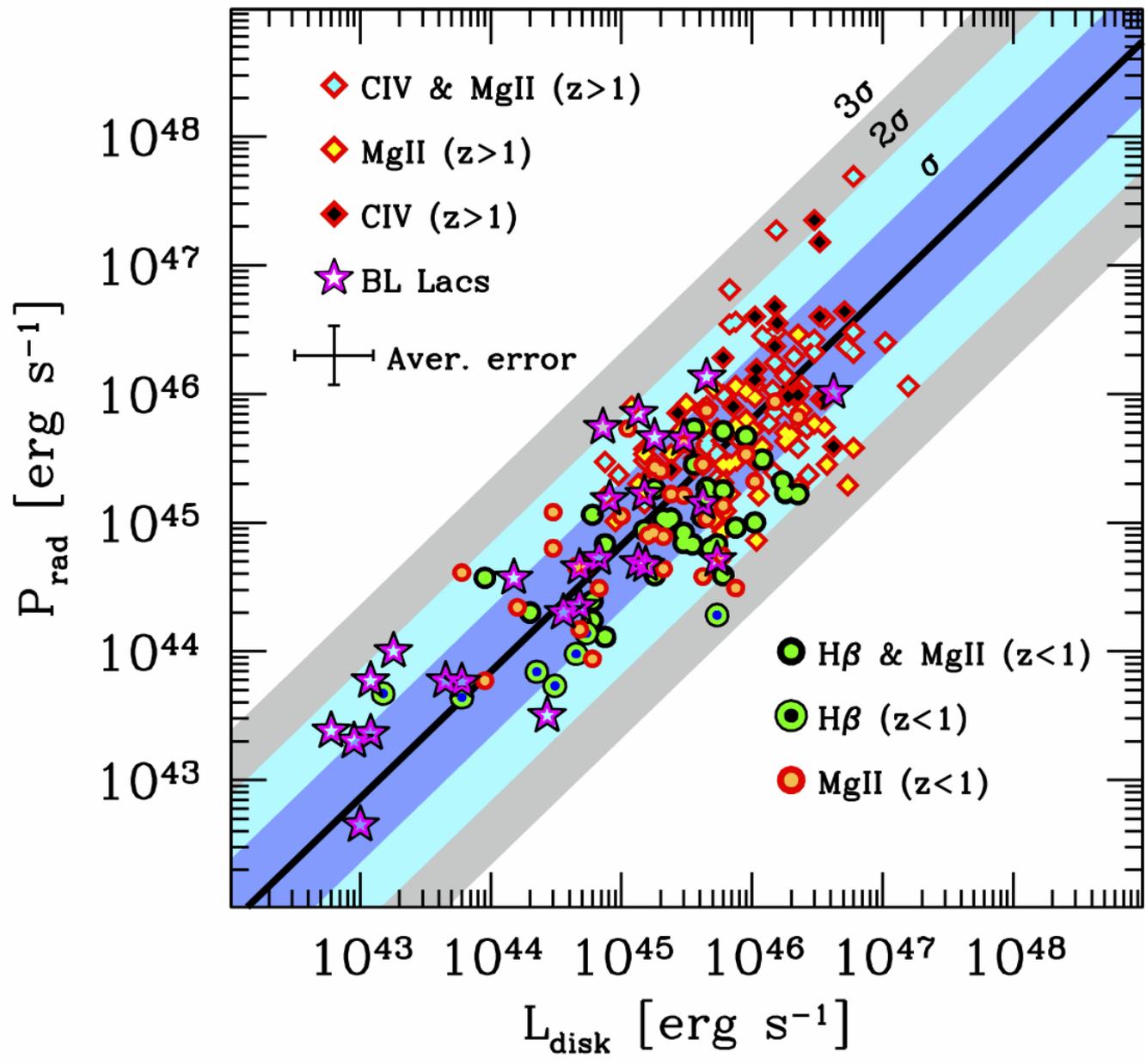
$$P_{\text{jet}} > \frac{L_{\text{obs}}}{\Gamma^2}$$

If P_{jet} is twice as much, Γ is

We can take P_{rad} as the m

This limit is model-indeper





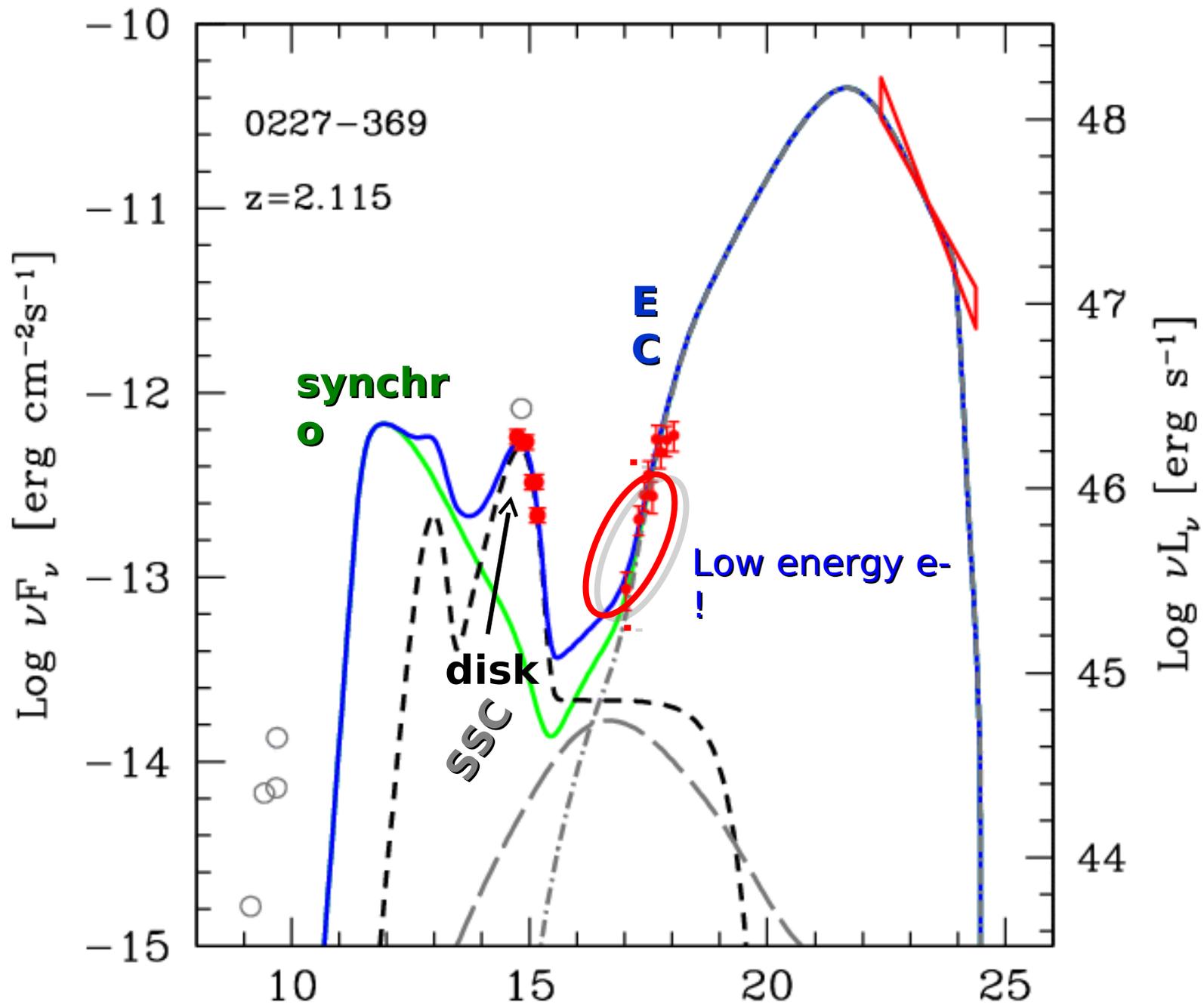
GG+ Nature, 2014

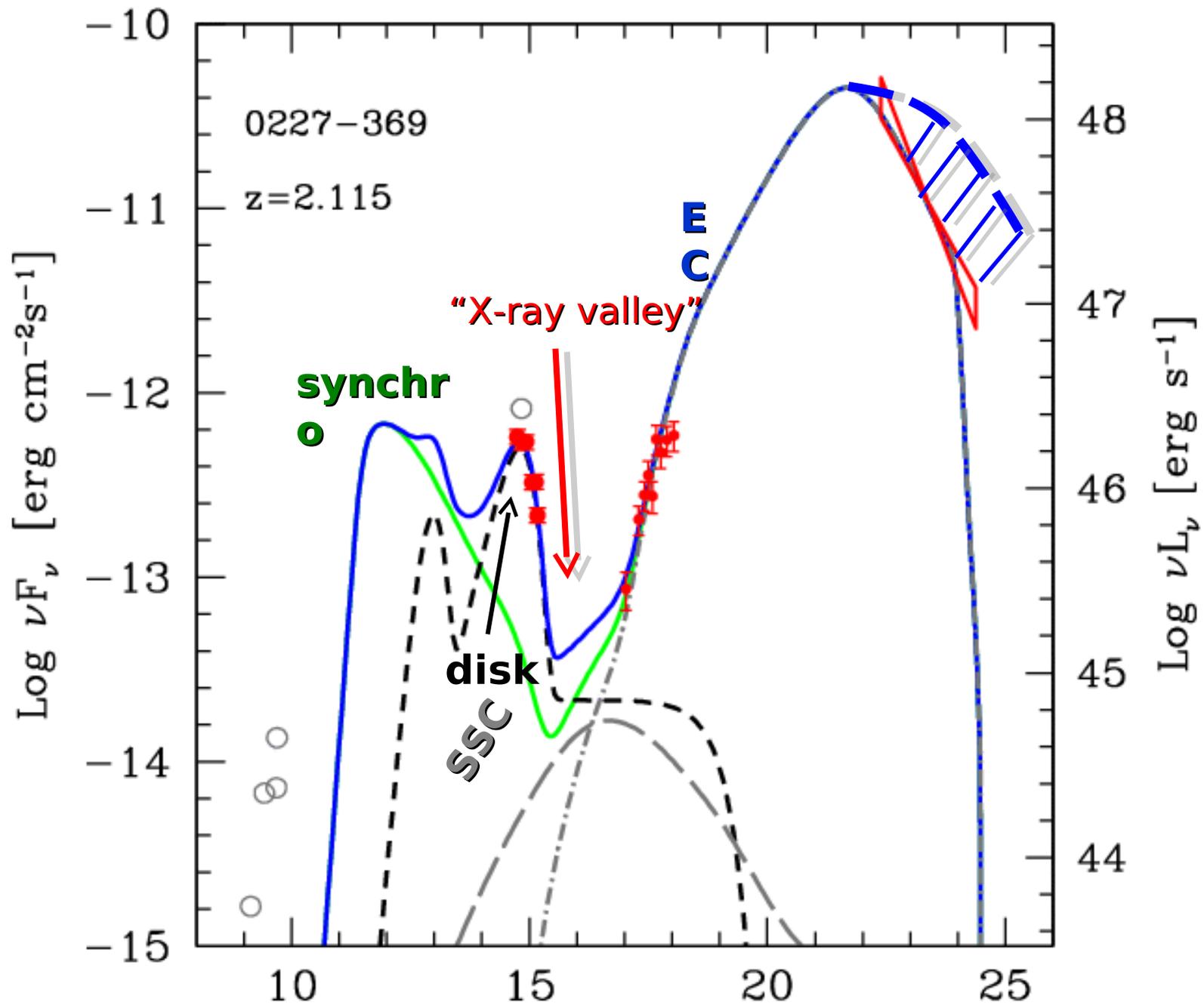
From P_{rad} to P_{jet}

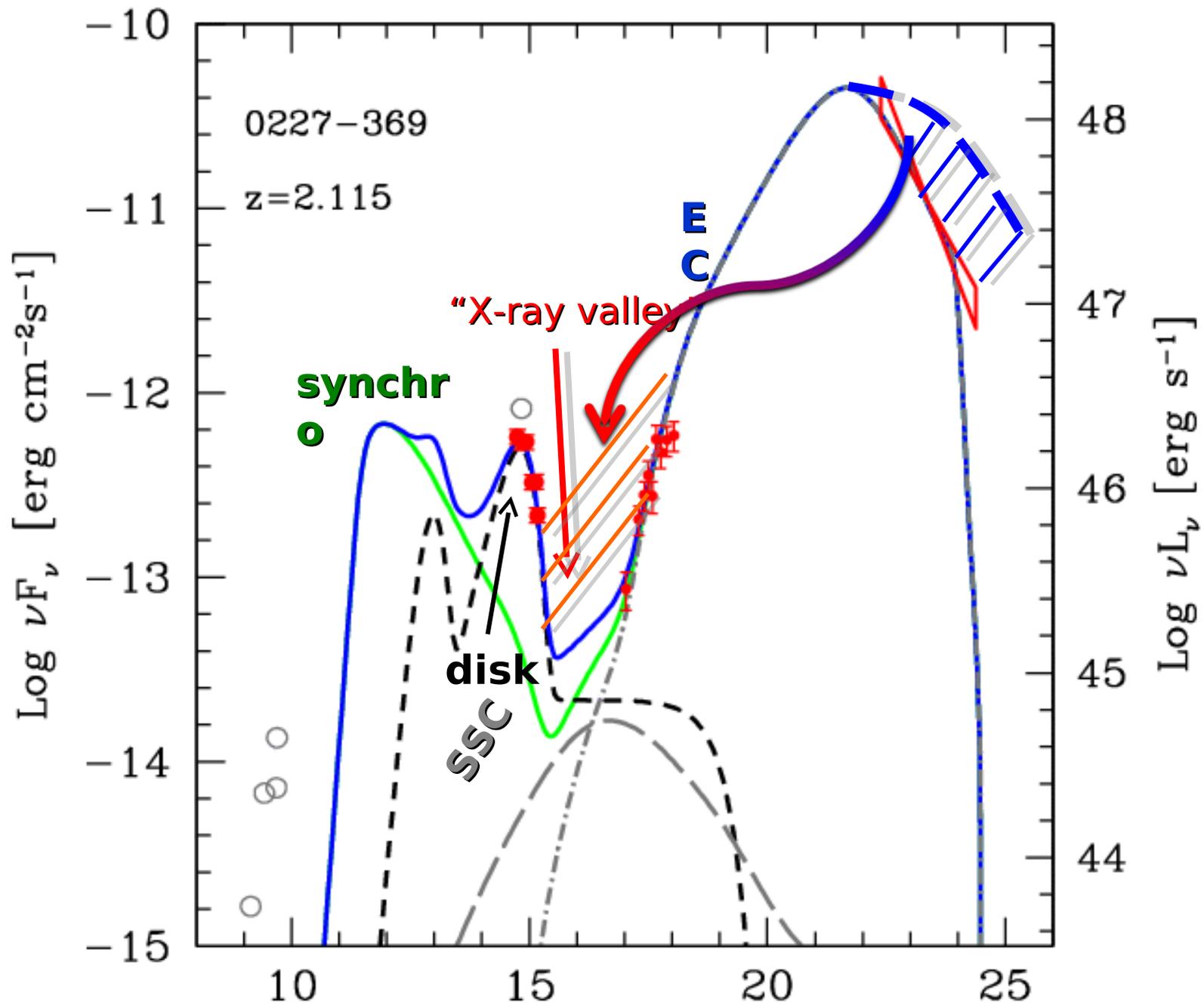
- γ_{min} \square total number of e- (and of protons, if no pairs...)
- Presence of electron positron pairs
- Radiative model: SSC vs EC vs baryons

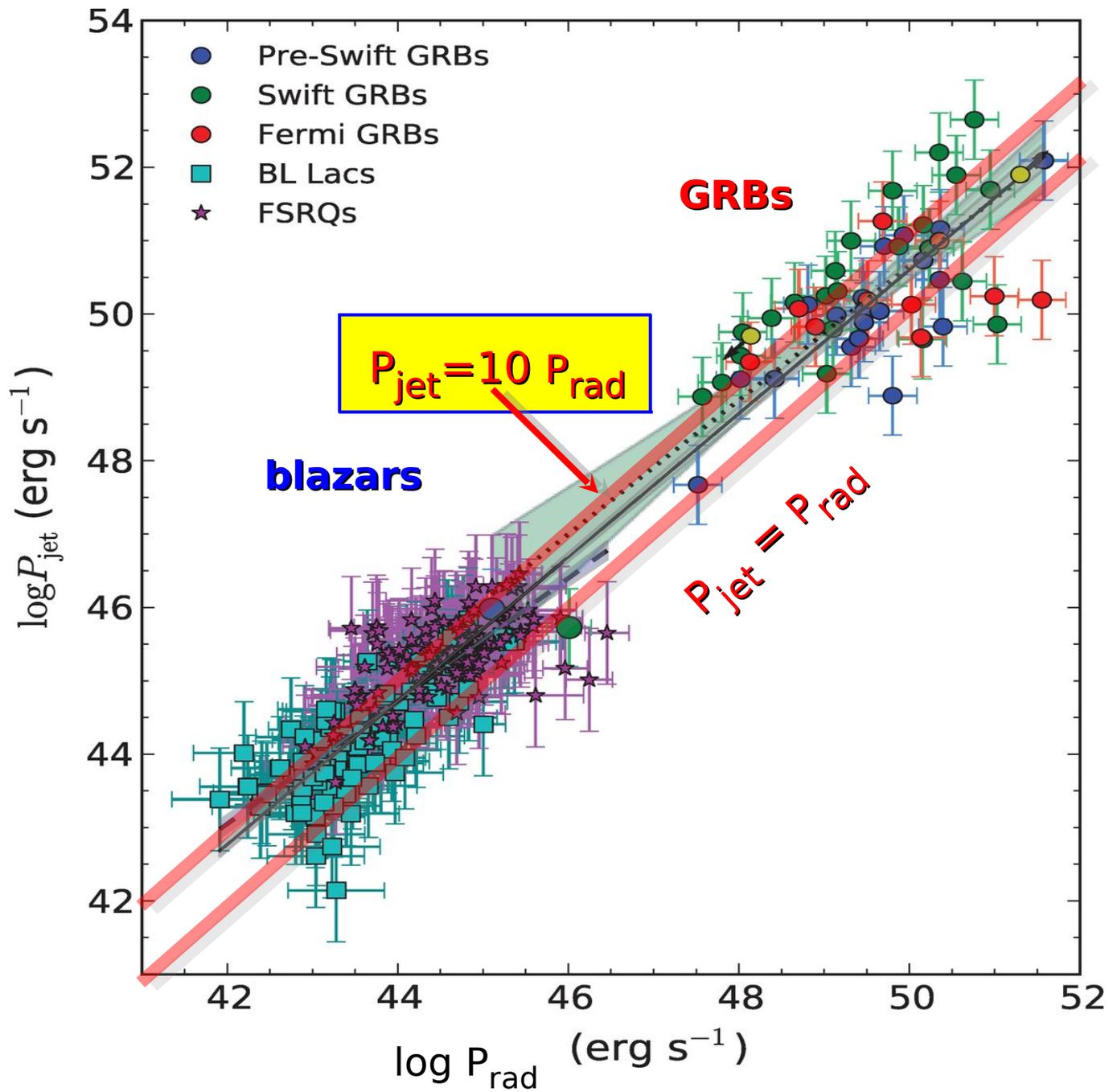
From P_{rad} to P_{jet}

- γ_{min} \square total number of e- (and of protons, if no pairs...)
- Presence of electron positron pairs
- Radiative model: SSC vs EC vs baryons
 - doesn't fit
 - even more power

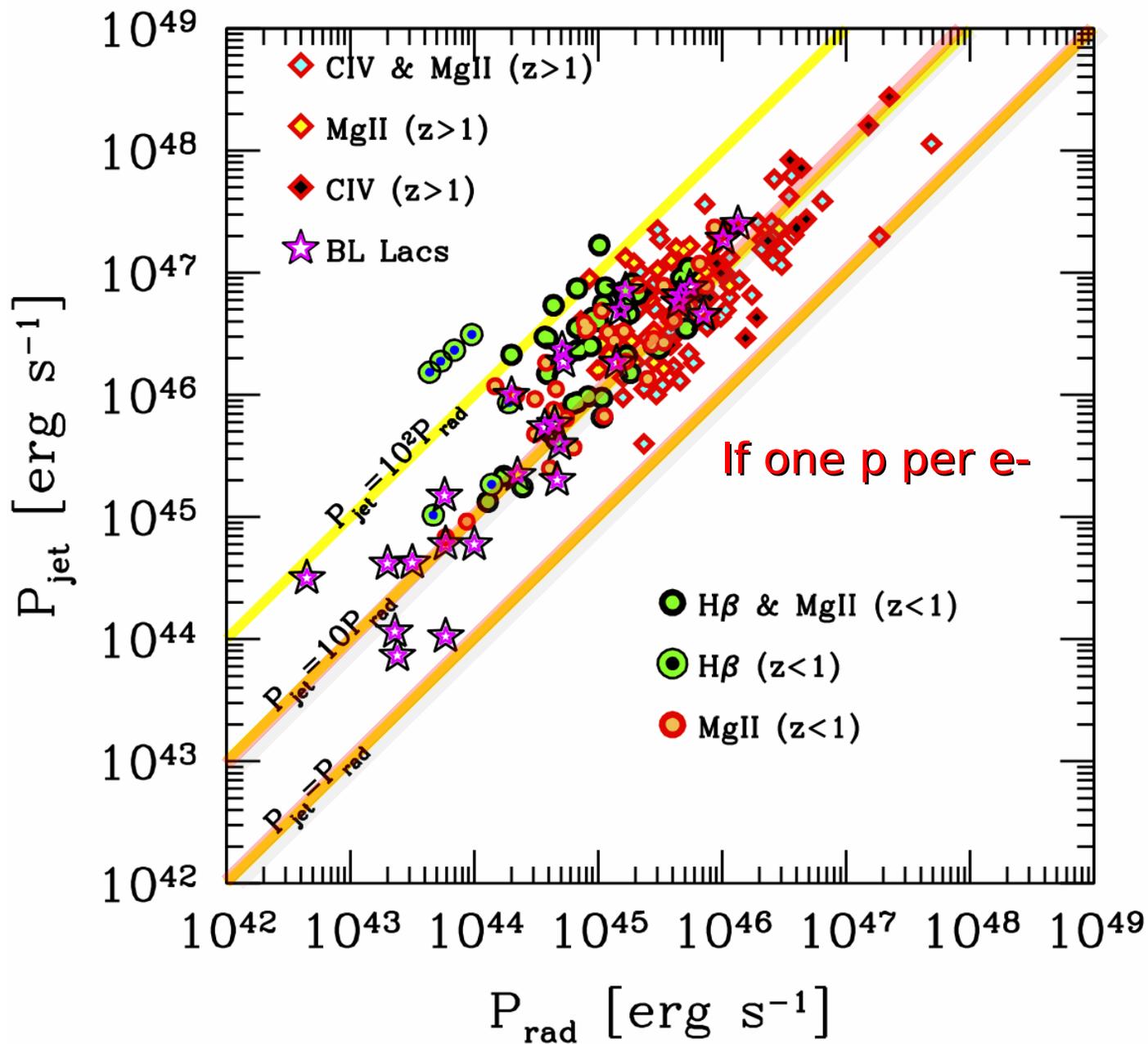






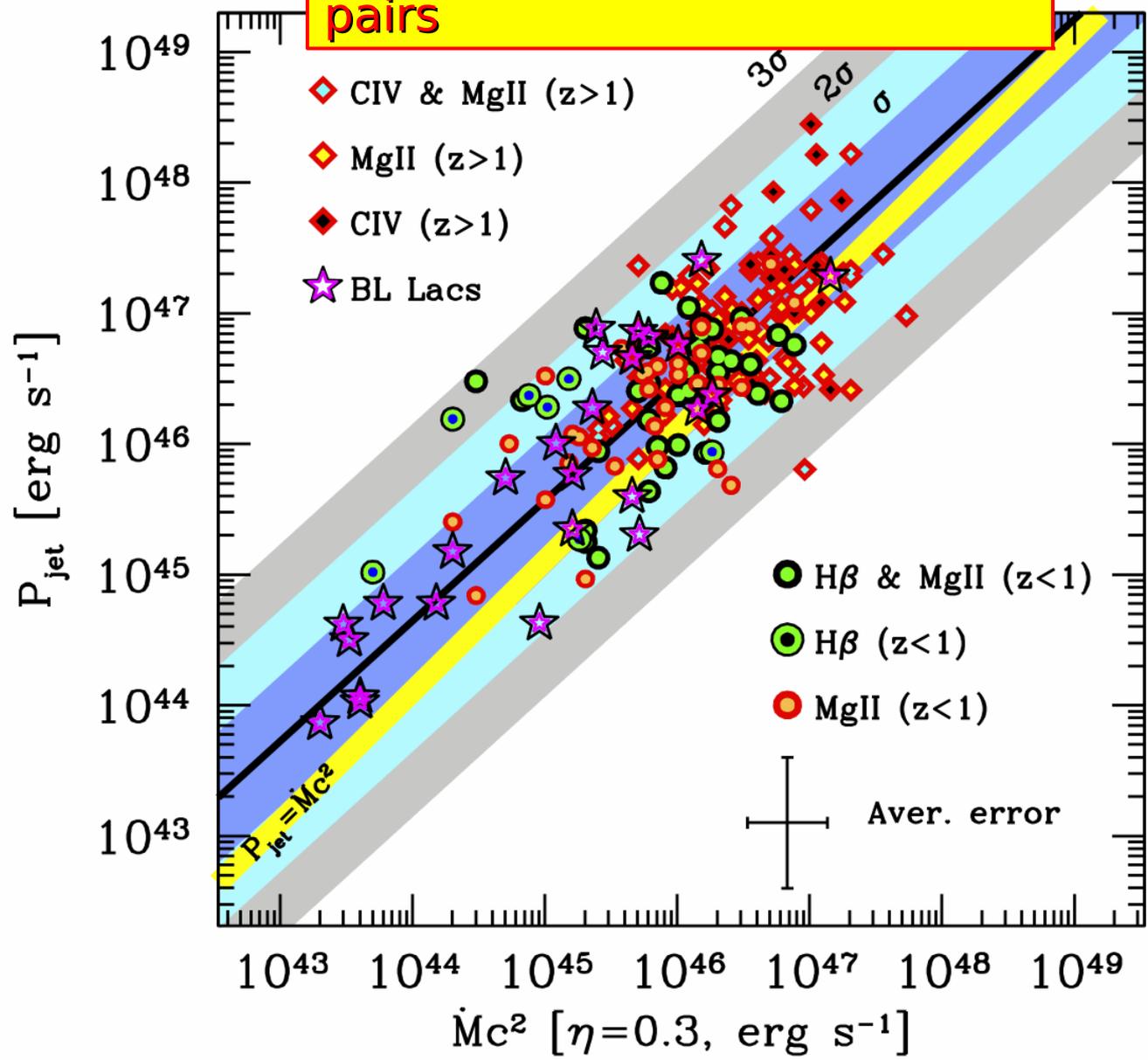


**Nemmen + Science
2012**

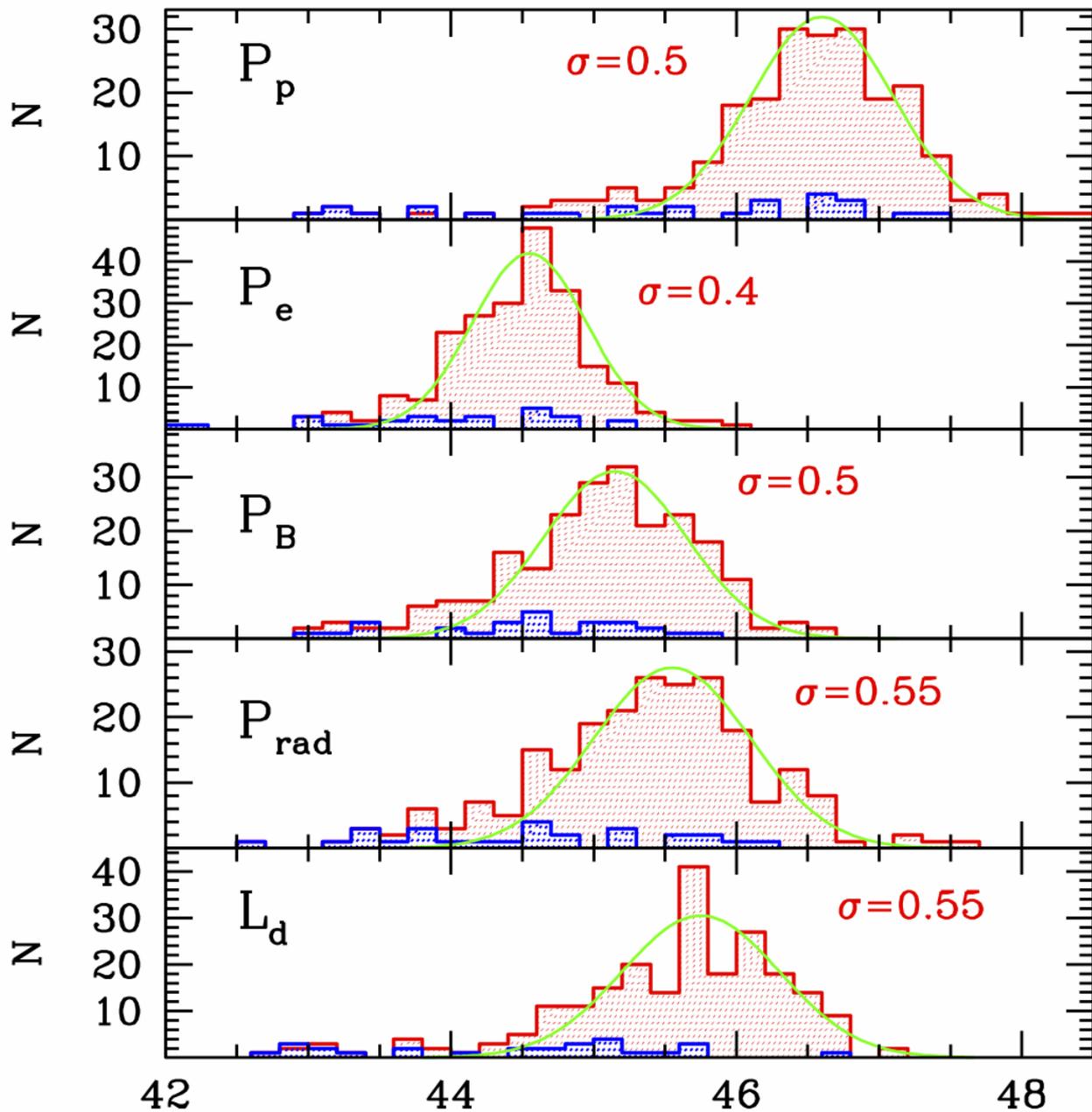


GG+ Nature 2014

1 proton per electron \Rightarrow no pairs



GG+ Nature, 2014



If one p per e-

Relat.
electrons

Magnetic Field

Radiation

Disk

Log Power [erg s⁻¹]

Apparent paradox:

Jet power proportional to accretion

And yet it is greater....

Blandford Znajek

$$P_{\text{BZ}} \sim a^2 B^2 M^2$$

Blandford Znajek

$$P_{\text{BZ}} \sim a^2 B^2 M^2$$



$$P_{\text{BZ}} \sim a^2 B^2 R_g^2 c$$

Poynting

Flux

Blandford Znajek

$$P_{\text{BZ}} \sim a^2 B^2 M^2$$



$$P_{\text{BZ}} \sim a^2 B^2 R_g^2 c \quad \text{Poynting}$$



Flux

$$P_{\text{BZ}} \sim a^2 \rho c^2 R_g^2 c \quad \rho c^2 \sim B^2 / 8\pi$$

$$P_{\text{BZ}} \sim a^2 \dot{M} c^2$$

Rotation > Accretion

B-field amplified by accretion can tap the spin energy of the hole

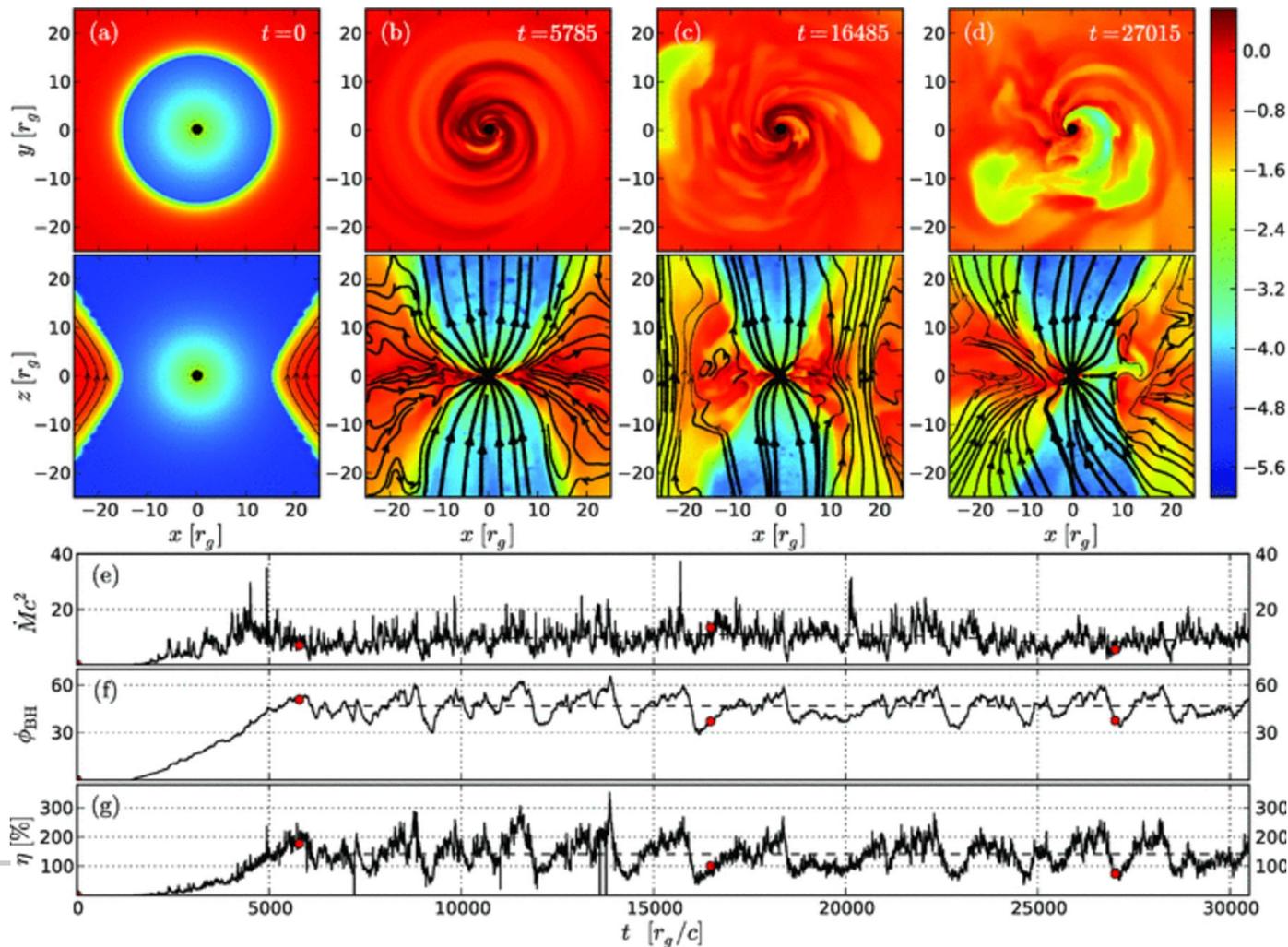
This process is very efficient

The B-field “does not work”, the jet power comes from the BH spin

Yet it is the catalyst for the process. No B no jet.

But B is linked to accretion, that's why P_{jet} is propto L_d

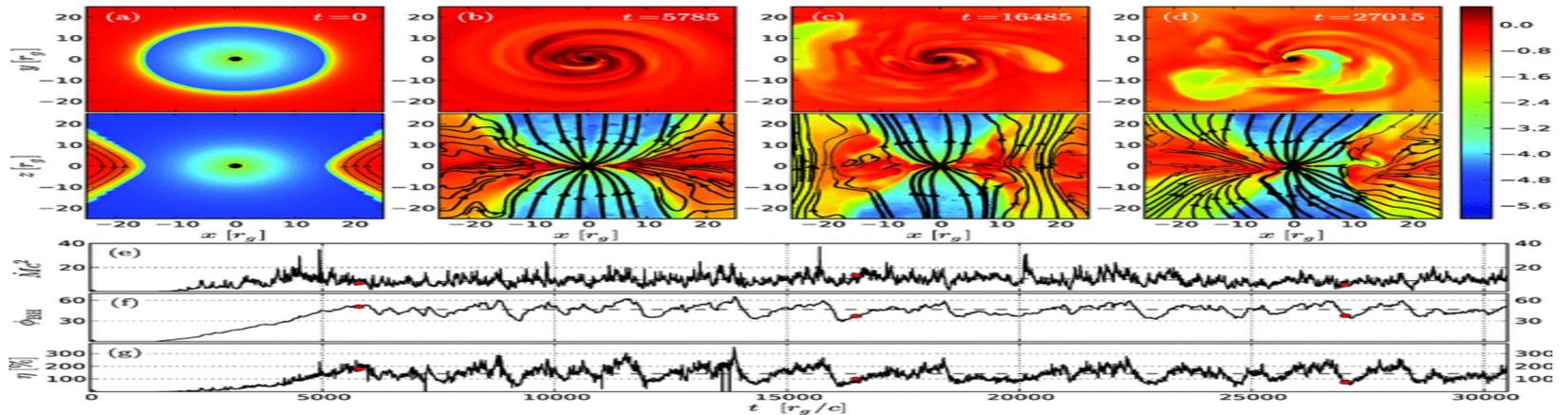
Shows results from the fiducial GRMHD simulation A0.99fc for a BH with spin parameter $a = 0.99$; see Supporting Information for the movie.



$$\frac{P_{\text{jet}}}{Mc^2}$$

Tchekhovskoy A et al. MNRAS 2011;418:L79-L83

Shows results from the fiducial GRMHD simulation A0.99fc for a BH with spin parameter $a = 0.99$; see Supporting Information for the movie.



On average: 1 g in $\approx 1.5 c^2$ erg out

Tchekhovskoy A et al. MNRAS 2011;418:L79-L83

Conclusion

- $P_{\text{jet}} \sim \dot{M}c^2$, larger than L_d

The jet uses to energy stored in the rotation of the hole, that was provided by accretion. But it takes it out faster than when it was put in.