Recent achievements on KMHD and applications in the ICM

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Collaboration: G. Kowal, E. de Gouveia dal Pino, R. Santos-Lima (IAG-USP) A. Lazarian (UW-Madison) Turbulence understanding have greatly improved in the past years;

Aims: understanding of energy transfer (scales) and diffusion of matter in the ISM

"direct" comparison of observables :

velocity (Esquivel et al. + Pogosian et al. + Falceta-Goncalves et al.)

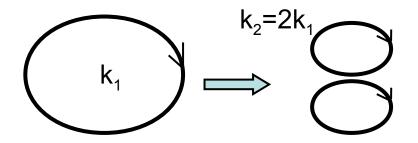
density (Kowal et al. + Burkhart et al.)

magnetic fields (Falceta-Goncalves et al. + Falgarone et al. + Heitch et al.)

Introduction

Turbulence: Kolmogorov's theory

incompressible fluids: Homogeneity + Isotropy + Scale invariance + Locality



Energy "flows" from large to small scales: $\dot{\varepsilon} \approx \delta v_l^2 / \tau_l$

being
$$\tau_l \approx l/\delta y_l$$
 therefore $\delta v_l \approx (\dot{\varepsilon}l)^{1/3}$ and, since
 $\delta v_l^2 \approx \int_{k=1/l}^{\infty} E(k') dk'_{we get} \qquad E(k) \approx \dot{\varepsilon}^{2/3} k^{-5/3}$

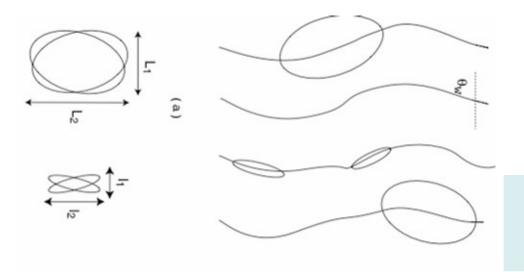
Magnetized Turbulence: Goldreich-Sridhar theory

strong wave-wave interactions

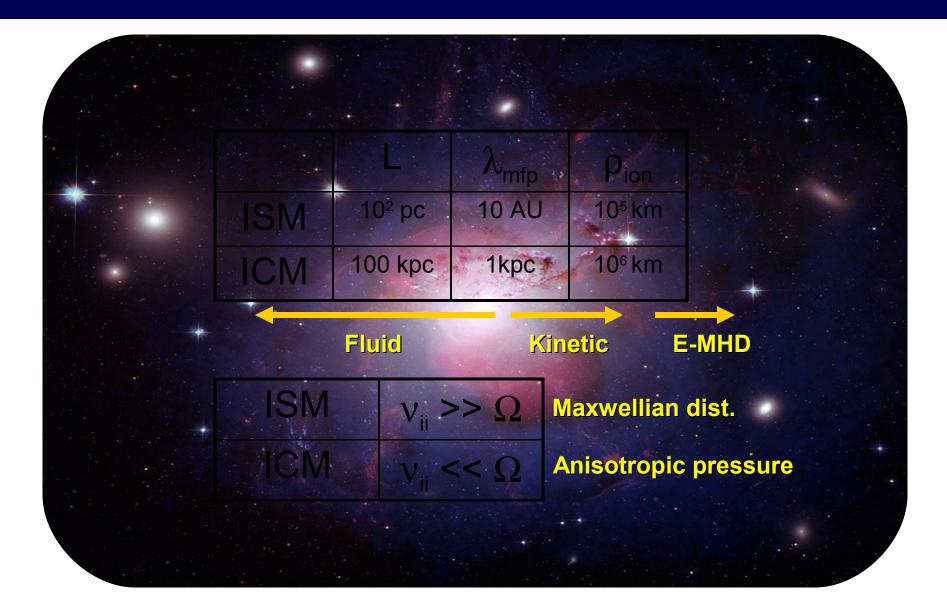
Anisotropy!

$$\tau_{\rm A} \sim \tau_{\rm s} \quad \Leftrightarrow \quad l_{\parallel}/l_{\perp} \sim v_{\rm A}/\delta u_l \qquad \qquad l_{\parallel} \sim v_{\rm A} \epsilon^{-1/3} l_{\perp}^{2/3} \sim k_{\parallel 0}^{-1} \left(l_{\perp}/l_{*}\right)^{2/3}$$

Kolmogorov spectrum for perpendicular direction



Anisotropy is: SCALE-DEPENDENT and LOCAL



Modeling collisionless plasma

... under fluid approximation

In the collisionless regime

moments of Vlasov equation

$$\begin{split} & \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho v) &= 0, \\ & \frac{\partial \rho v}{\partial t} + \nabla \cdot \left[\rho v v + \left(\mathsf{P} + \frac{B^2}{8\pi} \right) I - \frac{1}{4\pi} B B \right] = f, \\ & \frac{\partial B}{\partial t} - \nabla \times (v \times B) &= 0, \end{split}$$

$$\mathsf{P} = p_{\perp}\hat{I} + (p_{\parallel} - p_{\perp})\hat{b}\hat{b}$$

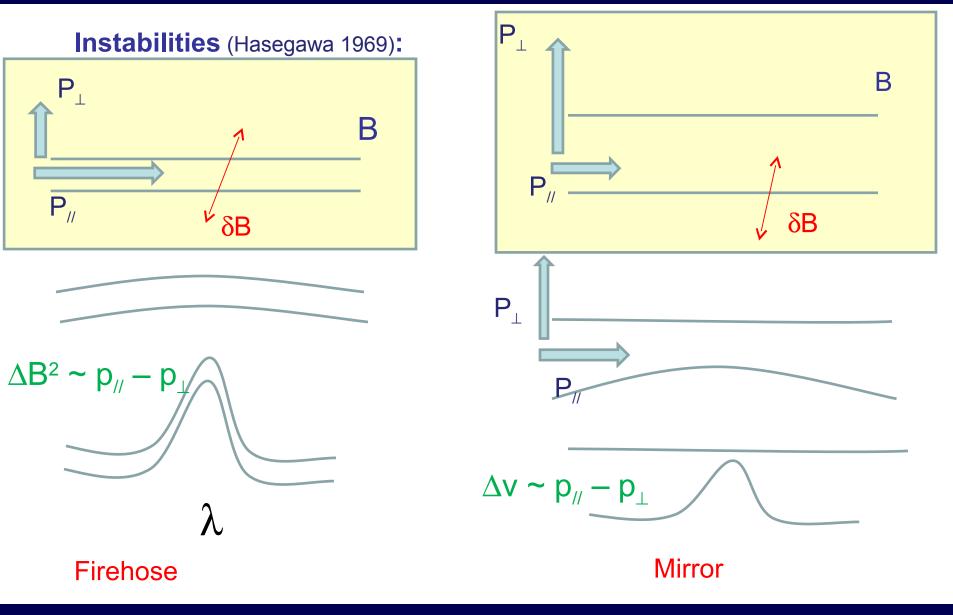
Closure
$$\frac{\partial S_{\perp}}{\partial t} + \nabla \cdot (S_{\perp} \boldsymbol{v}) = 0, \qquad S_{\perp} = p_{\perp} B^{1 - \gamma_{\perp}}$$
$$\frac{\partial S_{\parallel}}{\partial t} + \nabla \cdot (S_{\parallel} \boldsymbol{v}) = 0, \qquad S_{\parallel} = p_{\parallel} (B/\rho)^{\gamma_{\parallel} - 1}$$

(Chew, Goldberger, Low 1956)

Instabilities: $\beta_{\prime\prime} > 2\beta_{\perp}$ $\delta p_{\perp} \sim (1-T_{\perp}/T_{\prime\prime}) \delta B$

Double-isothermal case (the simplest case – not much): $p_{\perp} = a_{\perp}^2 \rho$ and $p_{\parallel} = a_{\parallel}^2 \rho$

Modeling



Kowal, Falceta-Goncalves & Lazarian 2011

Model	B_0	a_{\parallel}	a_{\perp}	$\xi \equiv a_{\parallel}^2/a_{\perp}^2$	$\mathcal{M}_s \equiv \langle v /a_\perp \rangle$	$\mathcal{M}_A \equiv \langle v /c_A \rangle$
1	1.0	1.0	2.0	0.25	0.7	0.7
2	1.0	1.0	0.5	4.00	0.7	0.7
3	0.1	0.1	0.2	0.25	7.0	2.0
4	0.1	0.1	0.05	4.00	7.0	2.0
5	0.1	1.0	0.5	4.00	0.7	2.0
6	1.0	0.1	0.2	0.25	7.0	0.7

Res = 256^3 512³

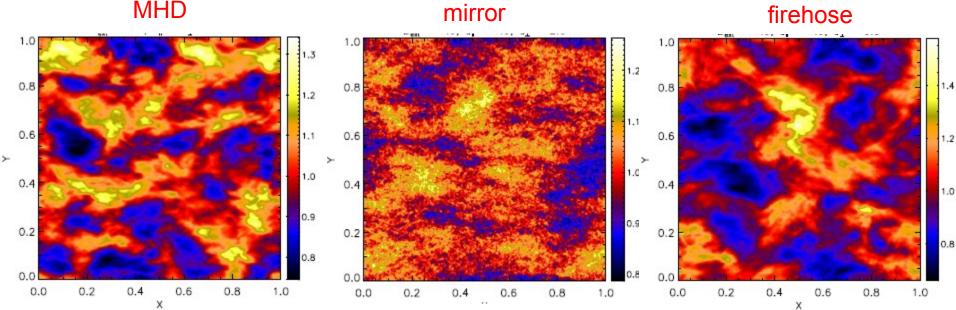
1024³

- Subsonic (supersonic) sub-alfvenic (super-alfvenic)
- Firehose / mirror-unstable

Modeling

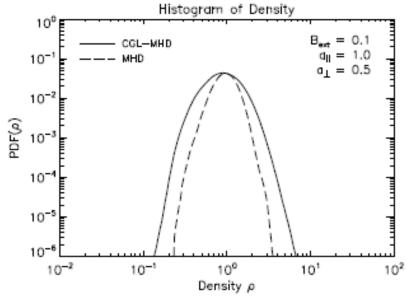
Sub-alfvenic (super-alfvenic are quite similar)

MHD

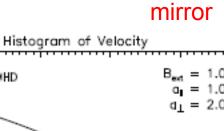


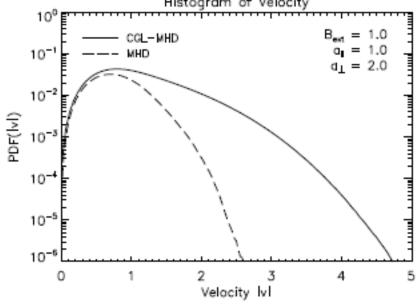
Modeling

Density and Velocity PDFs

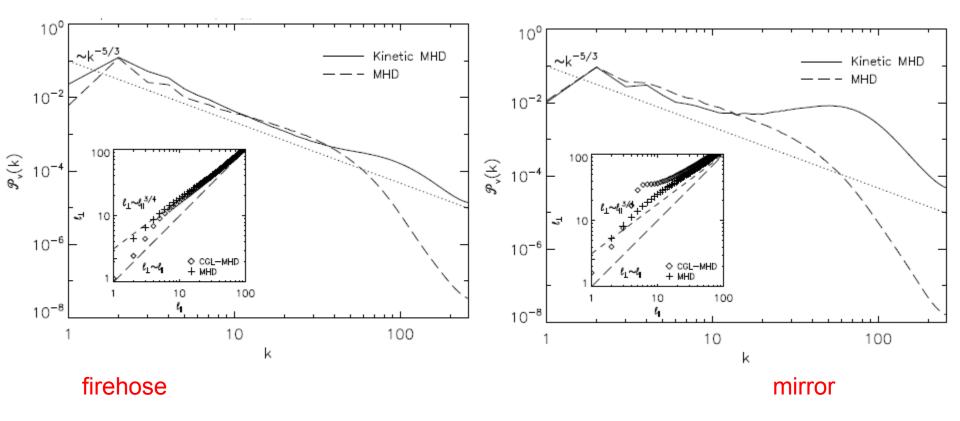




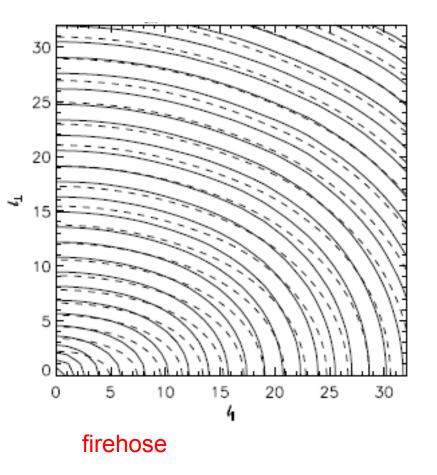


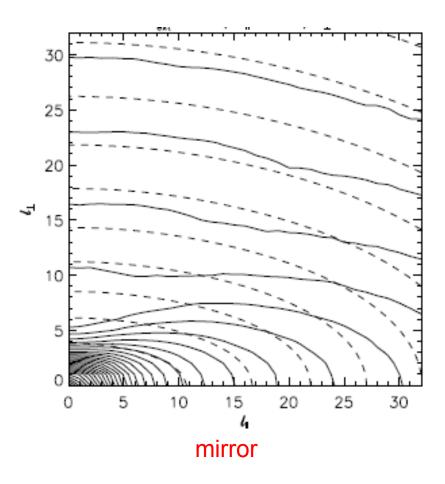


Spectra of velocity



2nd order structure function





Saturation of instability is fast!

$$\tau_{sat} \sim 0.1 \ \tau_{dyn} \ (<< \tau_{alf})$$

CGL-MHD turbulent dynamo (can be) is faster than MHD turbulent dynamo (do not miss Santos-Lima's talk)

Caveat: double-isothermal closure

External/internal processes could sustain pressure anisotropy (cosmic rays, anisotropic heat conduction, etc.)

The Intracluster medium

... best application

The ICM

The ICM and the cluster total mass

Usually, the ICM is believed to be in quasi-equilibrium :

$$\nabla P = -\rho \nabla \Phi$$

$$M(r) = -3.68 \times 10^{13} T(r) r \left(\frac{d \log \rho_{gas}}{d \log r} + \frac{d \log T}{d \log r} \right)$$

Sarazin (1988)

Density profile

For an isothermal plasma (isotropic), the emission is free-free:

Cavaliere & Fusco-Femiano (1978)

$$\Sigma(r) = \Sigma_0 \left(1 + \frac{r^2}{r_c^2}\right)^{-3\beta + 1/2} \longrightarrow n_e(r) = n_0 \left(1 + \frac{r^2}{r_c^2}\right)^{-3\beta / 2}$$

Vikhlinin et al. (2006)

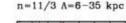
Relaxed clusters present cusps at central region, so the β model is not the best option

The ICM

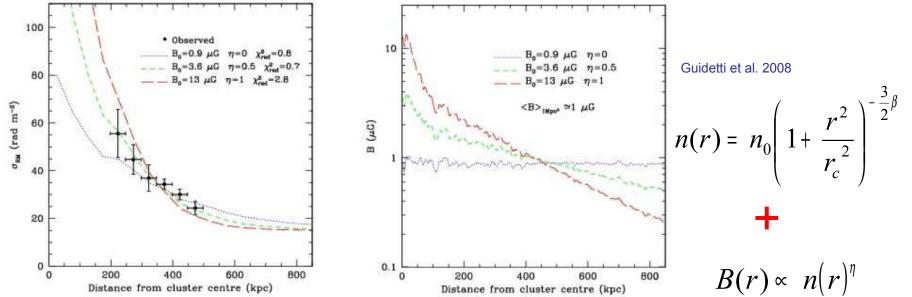
Faraday Rotation in the ICM

$$\psi_{obs} = \psi_{in} + \lambda^2 RM$$

$$RM \cong 812 \int_{0}^{L(kpc)} n_e(cm^{-3}) B_{//}(\mu G) dl \qquad (rad / m^2)$$



n=11/3 A=6-35 kpc



Turbulence in the ICM

$$K \approx \frac{U_g}{2}$$
 $T \approx 10^{7-8} K$
 $\sigma \approx 700 - 1000 \ km/s$ $C_s \approx 300 - 1000 \ km/s$

Subsonic (or mildly supersonic)

Injection scales ~ few hundreds of kiloparsecs

Total energy ~ 10⁶³ ergs

Energy density ~ 10⁻¹¹ erg/cm³

The ICM

KHMD models Probing magnetic fields

... or turbulent modes

0.4

0.2

0.0

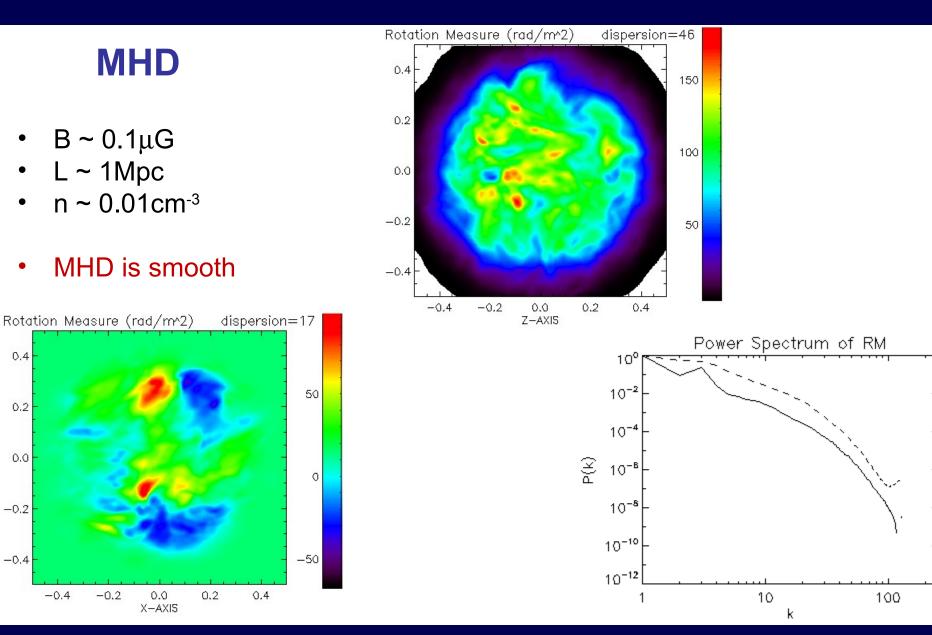
-0.2

-0.4

-0.4

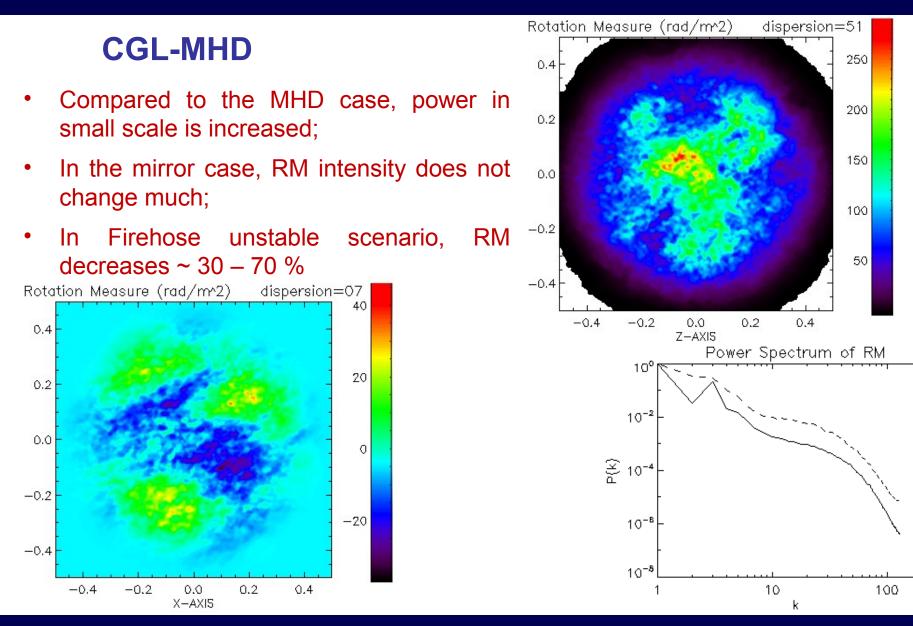
-0.2

Rotation Measure

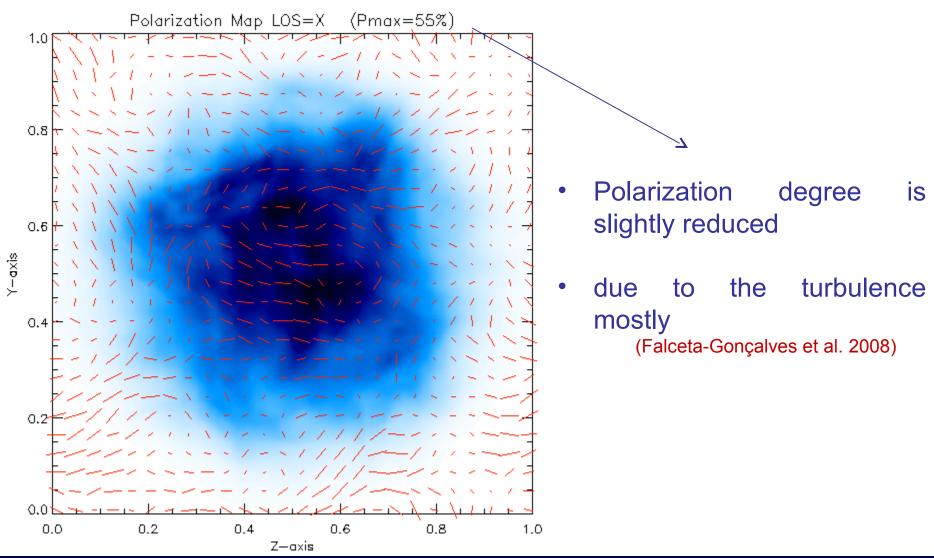


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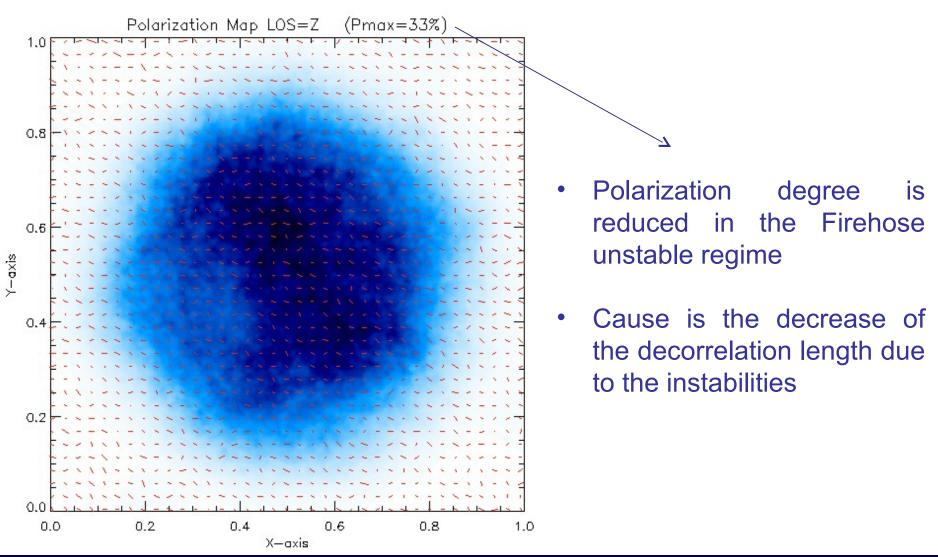


MHD



Polarization

CGL-MHD



Polarization

Conclusions

Some kinetic effects taken into account may change the dynamics of the plasma;

Depending on the regime, fast increase of magnetic field may take place;

As well as acceleration (fluid motions);

These occur at smaller scales;

Result in modifications of statistics of RM and synchrotron polarization

May increase reconnection rate and particle acceleration (seed for 2nd order Fermi)