Particle Acceleration by Relativistic Magnetic Reconnection in Blazar Jets

Lorenzo Sironi (ITC-Harvard) Krakow Jet Meeting, April 22nd 2015 LS & Spitkovsky, A. 2014, ApJL, 783, L21 LS, Petropoulou, M. & Giannios, D. 2015, arXiv:1502.0102

Powerful emission and hard TeV spectra

Blazar phenomenology:

(1) blazars are efficient emitters(radiated power ~ 10% of jet power)

(2) rough energy equipartition between emitting particles and magnetic field

(3) extended power-law distributions of the emitting particles, often with hard slope







Internal dissipation in blazar jets



Internal shocks in jets are likely to be poor non-thermal accelerators. (Giannios' and Spitkovsky's talks)

(LS+ 13, LS & Spitkovsky 09,11)

Shocks or Reconnection?

Internal shocks in blazars:

- electron-proton or pair plasma
- trans-relativistic (γ₀~a few)
- magnetized (σ >0.01)

$$\sigma = \frac{B_0^2}{4\pi\gamma_0 n_0 m_p c^2}$$

• toroidal field around the jet \rightarrow field \perp to the shock normal



Relativistic magnetic reconnection



Can relativistic reconnection self-consistently produce non-thermal particles?

The PIC method





No approximations, full plasma physics of ions and electrons



Tiny length and time scales (electron scales) need to be resolved
→ huge simulations, limited time coverage

• Relativistic 3D e.m. PIC code TRISTAN-MP (Buneman 93, Spitkovsky 05, LS+ 13, 14)

Dynamics and particle spectrum



- Reconnection is a hierarchical process of island formation and merging.
- The field energy is transferred to the particles at the X-points, in between the magnetic islands.

Hierarchical reconnection

 σ =10 electron-positron



- Reconnection is a hierarchical process of island formation and merging.
- The field energy is transferred to the particles at the X-points, in between the magnetic islands.
- Localized regions exist at the X-points where E>B.

Inflows and outflows

σ =10 electron-positron



- Inflow into the X-line is non-relativistic, with speed $v_{in} \sim 0.1$ c.
- Outflow from the X-points is ultra-relativistic, reaching the Alfven speed $v_A = c \sqrt{\frac{0}{1-1-c}}$

The particle energy spectrum

$\sigma=10$ electron-positron



• At late times, the particle spectrum in the current sheet approaches a broad power-law tail $dn/d\gamma \propto \gamma^{-p}$ of slope $p\sim 2$.

• The normalization increases, as more and more particles enter the current sheet.

 The mean particle energy in the current sheet reaches ~σ/2
 → energy equipartition

(LS & Spitkovsky 14)

The maximum particle energy





• The reconnection rate stays nearly constant in time, if the evolution is not artificially inhibited by the boundaries.

• The maximum energy grows at a rate proportional to the reconnection rate v_{in} , so that $\gamma_{max} \propto t$ (compare to $\gamma_{max} \propto t^{1/2}$ in relativistic shocks).





In 3D, the in-plane tearing mode and the out-of-plane drift-kink mode coexist.
The drift-kink mode is the fastest to grow, but the physics at late times is governed by the tearing mode, as in 2D.

3D: particle spectrum

$\sigma=10$ electron-positron



• At late times, the particle spectrum approaches a powerlaw tail of slope $p\sim2$, extending in time to higher and higher energies. The same as in 2D.

• The maximum energy grows as $\gamma_{max} \propto t$ (compare to $\gamma_{max} \propto t^{1/2}$ in shocks). The reconnection rate is $v_{in}/c \sim 0.02$ in 3D (vs $v_{in}/c \sim 0.1$ in 2D).

Particle acceleration mechanism



Two acceleration phases: (1) at the X-point; (2) in between merging islands

(2) Fermi process in between islands





 The particles are accelerated by a Fermi-like process in between merging islands.



 Island merging is essential to shift up the spectral cutoff energy.

 In the Fermi process, the rich get richer. But how do they get rich in the first place?

(1) Acceleration at X-points



• In cold plasmas, the particles are tied to field lines and they go through X-points.

- The particles are accelerated by the reconnection electric field at the X-points, and then advected into the nearest magnetic island.
- The energy gain can vary, depending on where the particles interact with the sheet.

Dependence on the flow conditions



As σ increases:

• the reconnection rate v_{in}/c increases with magnetization and it saturates at ~ 0.2 for σ >1 (LS & Spitkovsky 14; confirmed by Guo+ 14, Werner+ 14)

As σ increases:

- the power-law slope becomes harder
- \Rightarrow a probe of the flow magnetization?

Implications for blazar emission

Powerful emission and hard TeV spectra

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Implications for blazar emission

Relativistic reconnection in blazars:

(1) it transfers up to 50% of the flow energy (in pair plasmas) to the emitting particles.

(2) results in rough energy equipartition between particles and magnetic field.

(3) it produces power-law distributions, with slope harder than p=2 for $\sigma \ge 10$.



(LS & Spitkovsky 14, confirmed by Guo+ 14, Werner+ 14)



 B_{0}^{2}



• Relativistic magnetic reconnection in magnetically-dominated blazar jets is an efficient particle accelerator, in 2D and 3D.

• Relativistic reconnection ($\sigma \gtrsim 1$) can accelerate particles into a power-law tail with slope between -4 and -1 (harder for higher magnetizations), and max energy growing linearly with time, if slope is less than -2.

• The reconnection rate (and so, the rate of growth of the max energy) is around 0.1 - 0.2 *c* in 2D and ~ 0.02 *c* in 3D for the case of zero guide field. In 3D, the drift-kink mode is unimportant for the long-term evolution. The reconnection rate decreases with the strength of the guide field, in both 2D and 3D.

Magnetic reconnection satisfies all the basic requirements for blazar emission (Giannios' talk): it is fast and efficient, produces non-thermal particles with hard slopes, and results in rough energy equipartition between particles and fields.
With increasing guide field strength, the efficiency drops and the islands become more magnetically-dominated.