### Semi-Analytical GRMHD Jet Model -an complementary approach to GRMHD simulations-



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### Important Ingredients of "BH powered" Relativistic Jet

- Extraction of BH rotational energy by a hole-threading field line:
  - MHD version of BZ process (Takahashi et al. 1990)
  - Poynting flux dominated GRMHD flow
  - the development of outflow is constraint by the inflow
- Jet acceleration:
  - efficient conversion from magnetic to kinetic energy
  - fast magnetosonic surface exists, such that magnetic nozzle effect can take place (e.g. Camenzind 1989; Li et al. 1992; Begelman & Li 1994; Takahashi & Shibata 1998)
- Parabolic field line (e.g. observation of M87; Asada & Nakamura 2012)

### Outline

- GRMHD flow structure in the vicinity of BH
- Model setup and Results (based on ApJ 801:56 2015)
- Next Setup and Preliminary results
- Summary

### **GRMHD** Flow Structure

#### **GRMHD Simulation** (a/M=0.9375)

 $B_p$  field lines and characteristic surfaces



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## Model Setup

- stationary, axisymmetry magnetic dominated flow in Kerr spacetime
- cold flow; velocity=0 at separation surface (determined by a/M, field geometry, and field angular velocity)
- Four conserved quantities along field line: (E, L, mass loading, field angular velocity)
  = (separation surface, Alfven surface, magnetizaton, field angular velocity)
- remaining two condition can be determined by: passing fast surface + matching condition

# passing fast surface (I)

inflow: pass fast surface always exist (required by causality)



# passing fast surface (II)

- outflow: no fast surface exist when a force-free field line is considered
- MHD perturbed of force-free field is essential for the existence of fast surface (Beskin and Nokhrina 2006)
- prescribed field:  $\Psi = \Psi_0 + \epsilon f_1$

$$\Psi_0 = \frac{\pi C}{\Omega_F} \sinh^{-1} \left( \Omega_F r (1 - \cos \theta) \right)$$

$$\epsilon f = \epsilon \pi C \Omega_F r \sin \theta, \quad \epsilon \ll 1$$



The matching condition single out a <u>unique</u> outflow solution for a given perturbation!

Outward energy flux  $\mathcal{E}^{r} = \mathcal{E}_{FL}^{r} + \mathcal{E}_{EM}^{r}$   $= nE_{FL}u^{r} + nE_{EM}u^{r}$   $= -n\mu u_{t}u^{r} - \frac{\Omega_{F}}{4\pi} \frac{B_{\phi}}{\Sigma \sin \theta} A_{\phi,\theta}$ 

$$\left(\mathcal{E}_{\mathrm{EM}}^{r}\right)^{-} = \left(\mathcal{E}_{\mathrm{EM}}^{r}\right)^{+}$$

 $rac{\Omega_F}{4\pi} (B_\phi)^- = rac{\Omega_F}{4\pi} (B_\phi)^+$ 







## Next Step

- consider super-fast magnetosonic regime
- consider multi field lines
  - mass loading is different for different field line

### Trial and preliminary results

q=0 : force-free case; fast surface exist when q is not zero

fixed stream line geometry (BZ parabolic field)

$$\Psi(r,\theta) = \frac{\mathcal{C}}{2} \left\{ r(1-\cos\theta) + 2\left(1+\cos\theta\right) \left[1-\ln\left(1+\cos\theta\right)\right] \right\}$$
  
modified flux function  
(related to neighbour field lines)  
$$\Phi = \sqrt{-g} \tilde{B}_{p} = \sqrt{-g} \left[B_{p}^{2}\Pi(q)\right]^{1/2}$$
  
$$\Pi(q) = \left(\frac{r}{r_{c}}\right)^{-q}$$
  
r<sub>c</sub> is chosen at t



**r**c is chosen at the outer light cylinder







## Summary (I)

- Extraction of BH rotational energy by a hole-threading field line:
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# Summary (II)

- MHD perturbation is essential for the existence of the fast magnetosonic surface
- Matching condition: the outward energy flux is continuously propagate outward
- The matching condition single out a unique outflow solution
- Next step: exploring solutions in super-fast regime, and solutions of global field configuration