

Large jets from small-scale magnetic fields



Dimitrios Giannios
Purdue, Department of Physics

Relativistic Jets: Creation, Dynamics and Internal Physics
Krakow, April 20-24

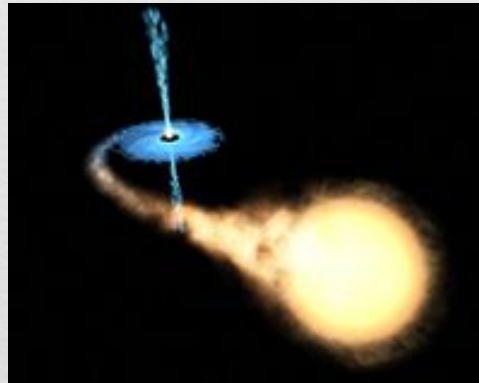
Relativistic jets from black holes and disks of all sizes/powers



jets in galactic centers

X-ray binaries

gamma-ray bursts



MBH $\sim 10^9 M_{\odot}$

$\sim 10 M_{\odot}$

$\sim 3 M_{\odot}$

Power $\sim 10^{44} \dots 10^{49} \text{ erg/s}$

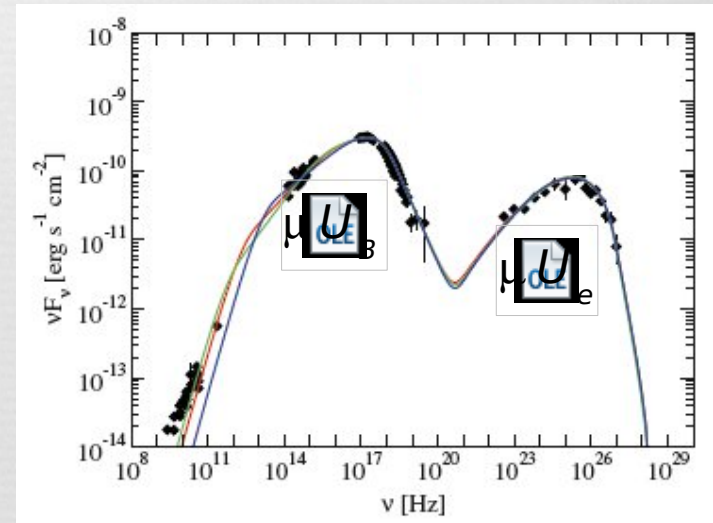
$\sim 10^{38} \text{ erg/s}$

$\sim 10^{52} \text{ erg/s}$

Basic Observed properties



SSC model fit of the spectral energy distribution of Mrk 421



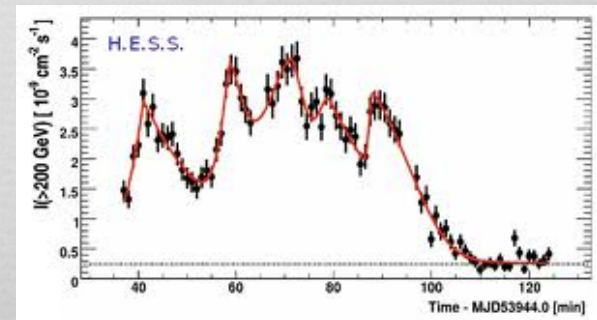
Efficient radiators ~10% radiative efficiency or more!

Effective particle accelerators up to multi TeV energies, *at least*

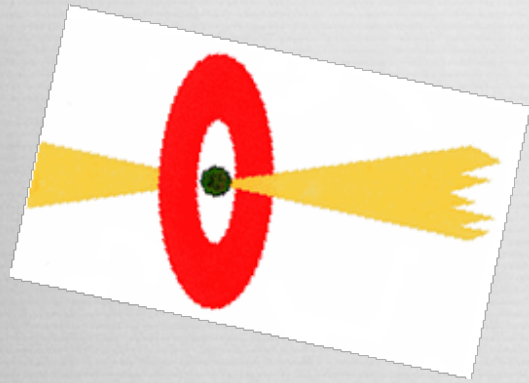
Radiating particles and magnetic field in very rough equipartition:

$$U_e \approx U_B$$

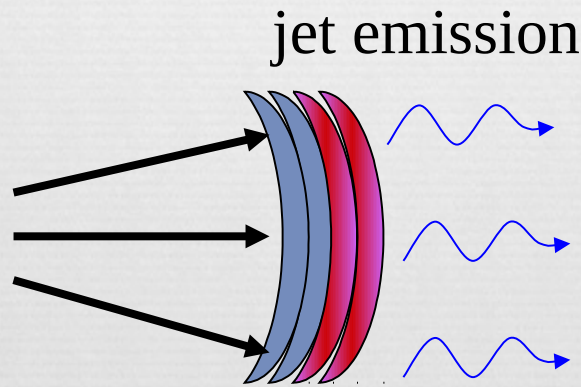
Other properties: highly variable ($t < R_g/c$), interesting polarization swings...



How is the jet flow launched?



Central engine



Acceleration

Internal dissipation

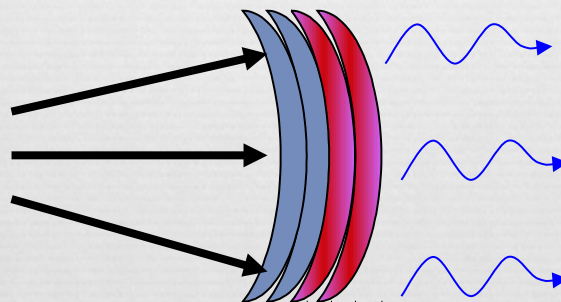
$$U_e \left[\frac{U}{c} \right] U_B$$

How is the jet flow launched?



jet emission

$$U_e \sim \frac{U}{c} ???$$



Acceleration

Internal
dissipation

$$\sigma_{initial} = \frac{B^2}{4\pi\rho c^2} \gg 1$$

Blandford & Znajek 1977

Begelman & Li 1992

Meier et al. 2001

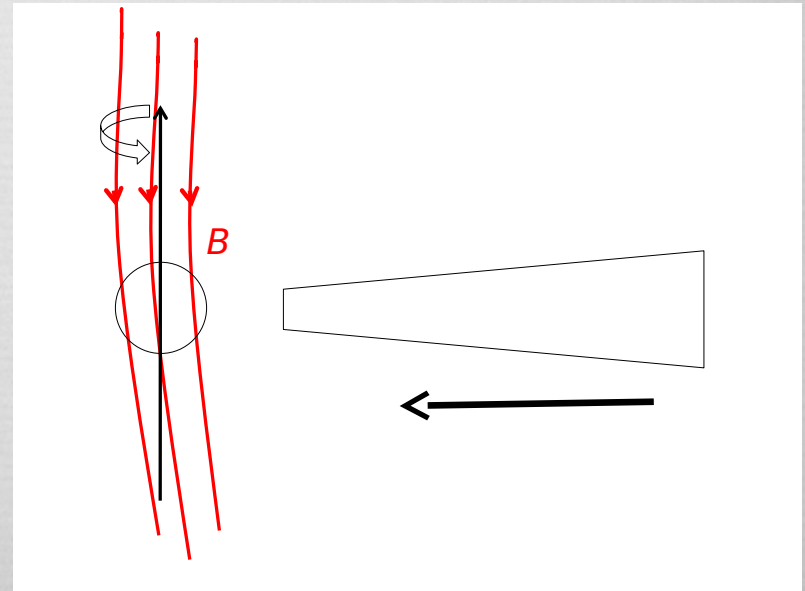
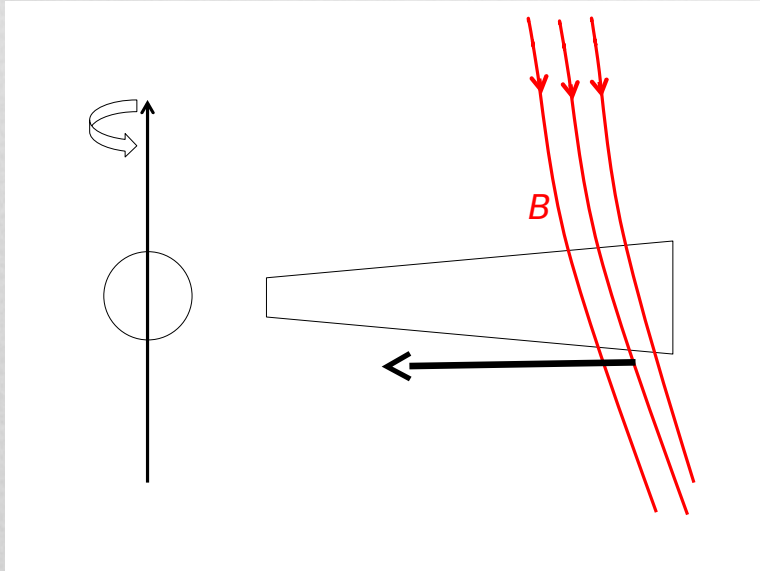
Koide et al. 2001

van Putten 2001

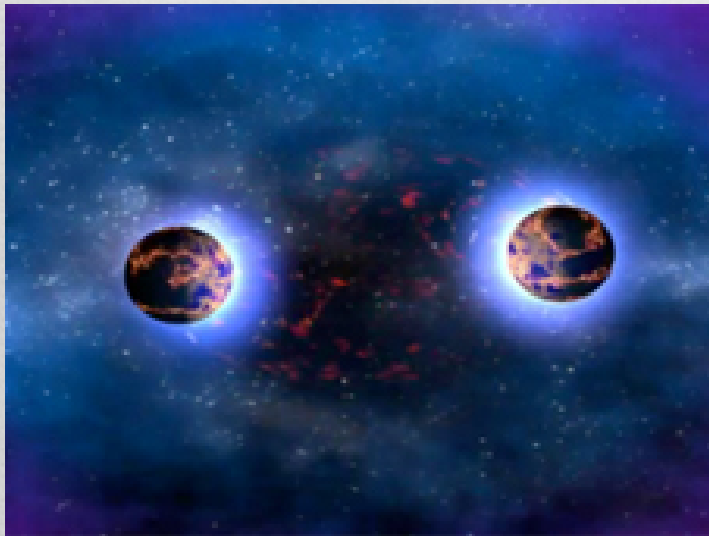
Komissarov, Lyubarky, McKinney, Tchekhovskoy

...

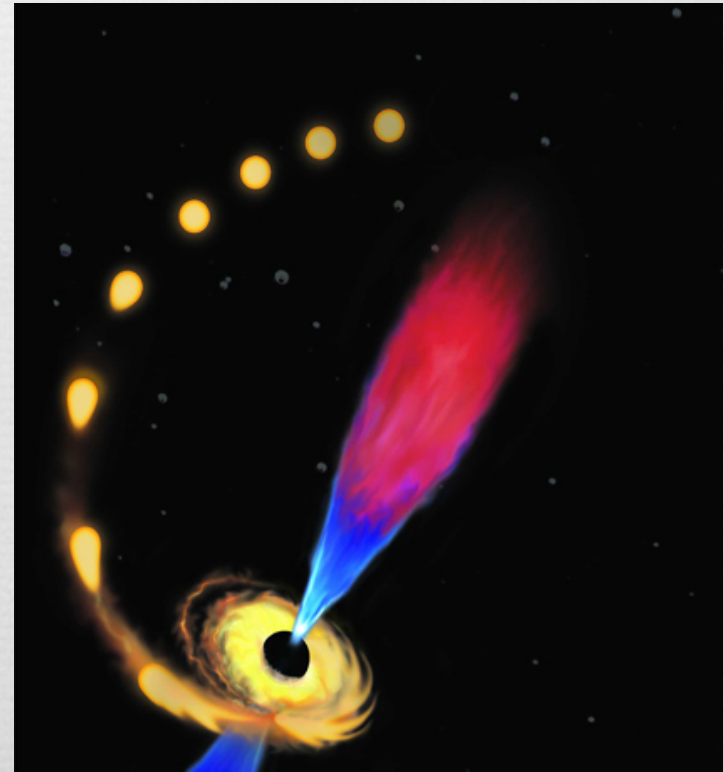
Origin of magnetic fields: carried in from large scales?



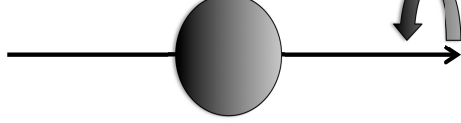
When the fields at the progenitor are *not* sufficient



binary neutron star mergers:
Short GRBs



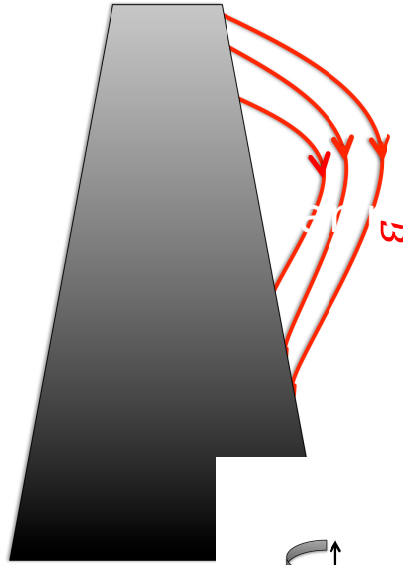
Stellar Tidal Disruption Jets (J1644)
Giannios & Metzger 2011; Bloom et al.
2011



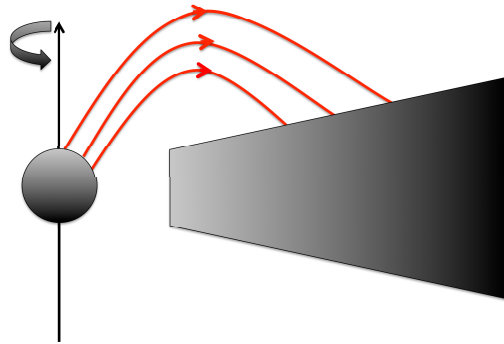
magnetic fields: locally by the disk?

annios, Beloborodov 2015

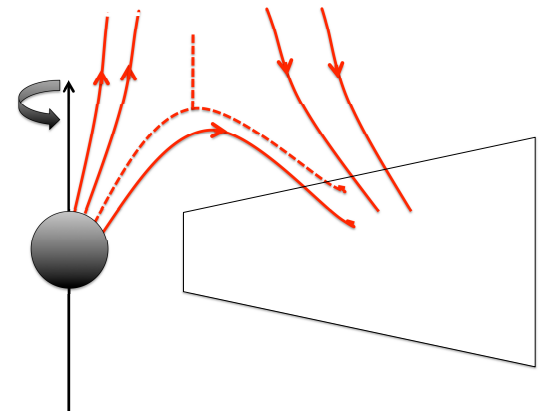
part of the
ected to



Step III
Differential
rotation inflates
loop
□ field lines open
jet forma

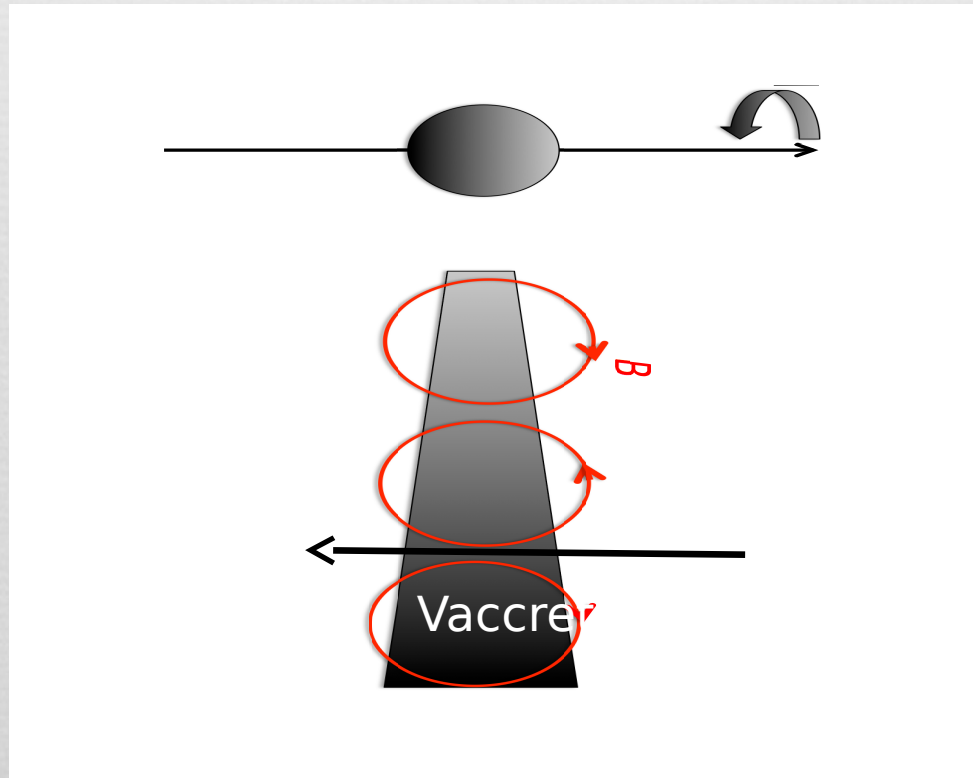


Step I
magnetic loop
emerging from the
disk
loop scale ~disk
scale high

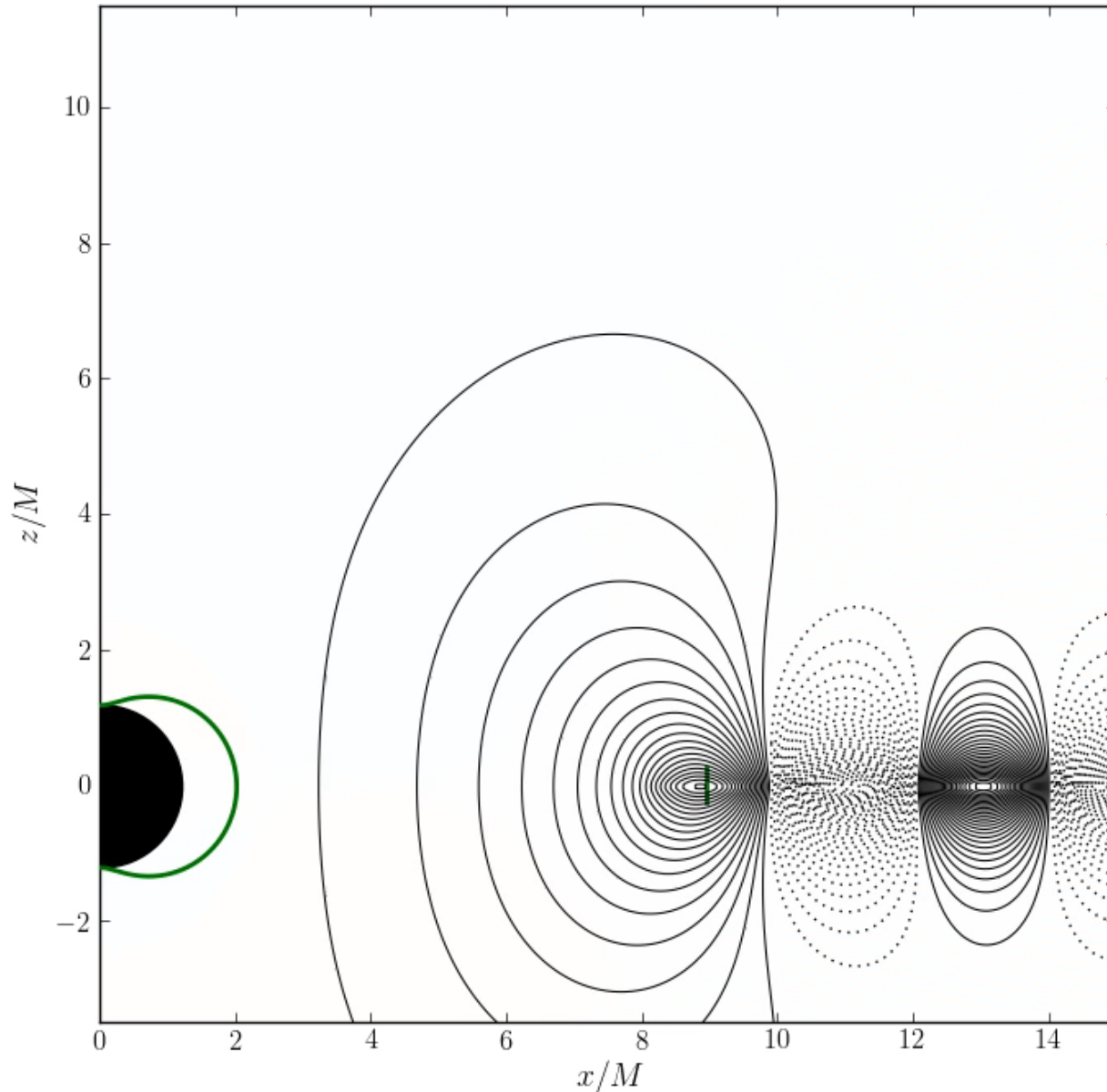


GR (force free) MHD simulations

Parfrey, Giannios, Beloborodov 2015



$t = 0 M$



Retrograde disc

loop width = $2 r_g$

loop direction

↙ clockwise

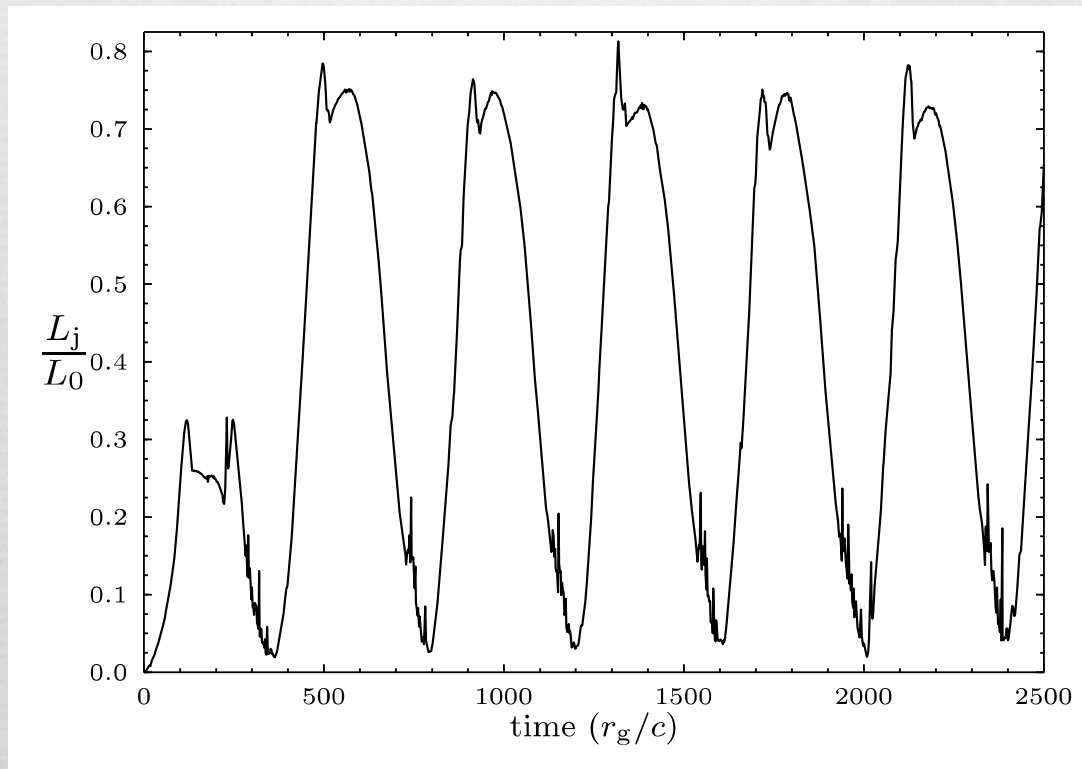
↘ anti-clockwise

Colour: H_ϕ positive
negative

$$v_{\text{accrete}} = c/100$$

$$a/M = 0.98$$

The jet power



The Jet Power (continued)



$$L_j \propto \frac{\Phi^2 a^2}{R_g^2} c, \quad \Phi = 2\pi r(l/2)B$$

$$\frac{B^2}{8\pi} = \beta P_g = \beta \rho c_s^2 = \beta \alpha \left(\frac{H}{r} \right)^2 v_K, \quad \dot{M} = 4\pi r H \rho v_r$$

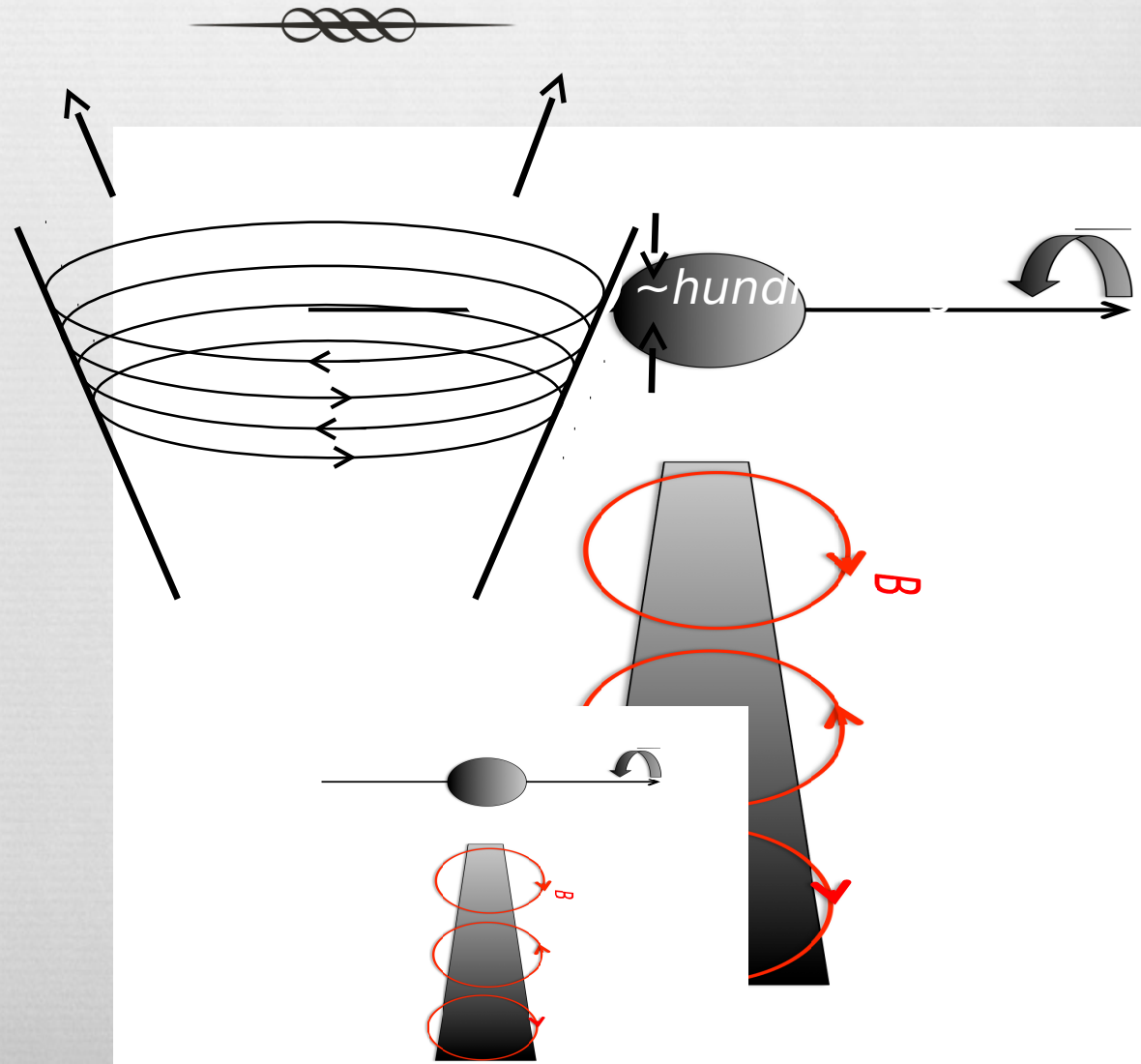
$$L_j \approx 0.16 \left(\frac{H}{r} \right) \left(\frac{r_{isco}}{r_g} \right)^{3/2} a^2 \dot{M} c^2$$

For *retrograde* orbits $r_{isco}/r_g \sim 10$ and *thick*

$H/r \sim 1$ disks the jet power reaches

$$L_j \sim \dot{M} c^2$$

Part II: large scale jet and emission



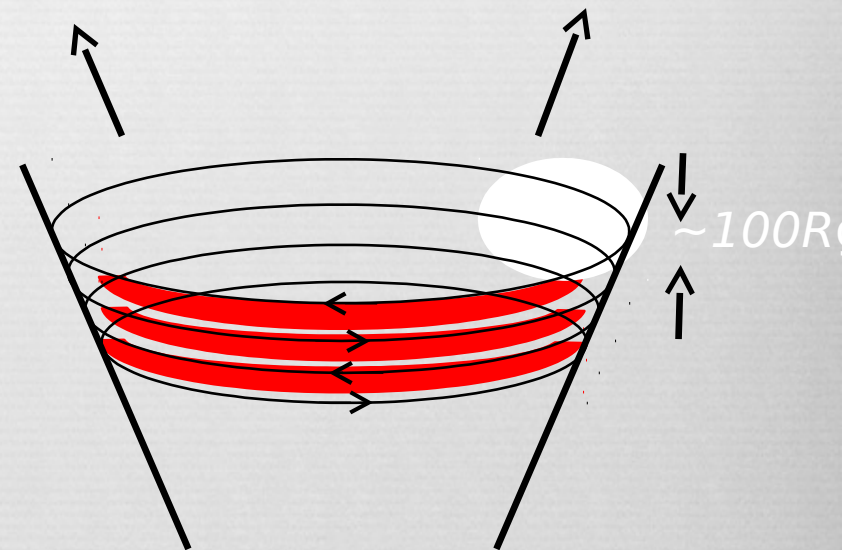
Part II: How to make you jet shine



Dissipation at large scale in the jet



- Jet may contain field reversals on small scale \sim hundreds r_g
- It remains magnetically dominated
- magnetic-reconnection becomes effective when $\frac{B^2}{4\pi\rho c^2} > 1$
 $t_{exp} \sim t_{rec}$
 where $v_{rec} = \epsilon c$



$$r_{diss} / \Gamma_j c \sim 100 \Gamma_j r_g / \epsilon c$$

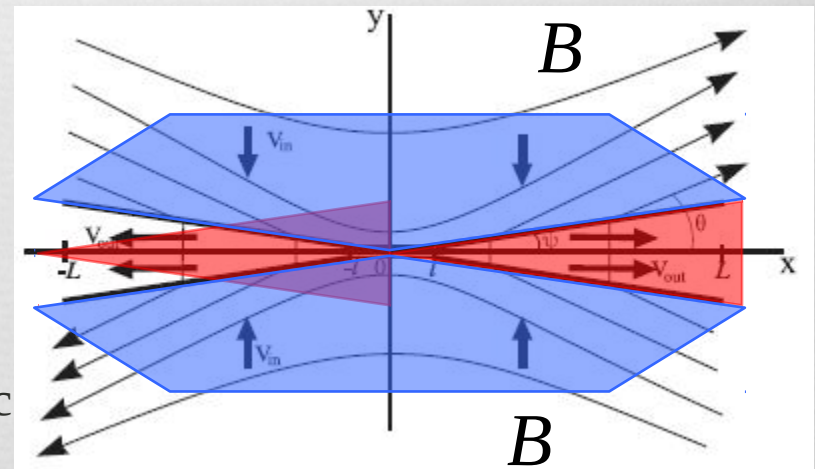
$$r_{diss} \sim \Gamma_j^2 100 r_g \frac{v}{c} \sim 1 M_8 \Gamma_{j,10}^2 \epsilon^{-1} \text{pc}$$

$$t_{flare} \sim r_{diss} / \Gamma_j^2 c \sim M_8 \epsilon^{-1} \text{days}$$

Relativistic Magnetic Reconnection



- An efficient convertor of magnetic energy into bulk motion, heat, energetic particles
 - cold, magnetized plasma enters the reconnection region
 - plasma leaves the reconnection region at the Alfvén speed $\Gamma_{out} \sim (1 + \sigma)^{1/2} > 1$
 - reconnected material contains energetic (nonthermal) particles



Relativistic Petschek Reconnection

Lyubarsky 2005

*Reconnection downstream:
emitting region*

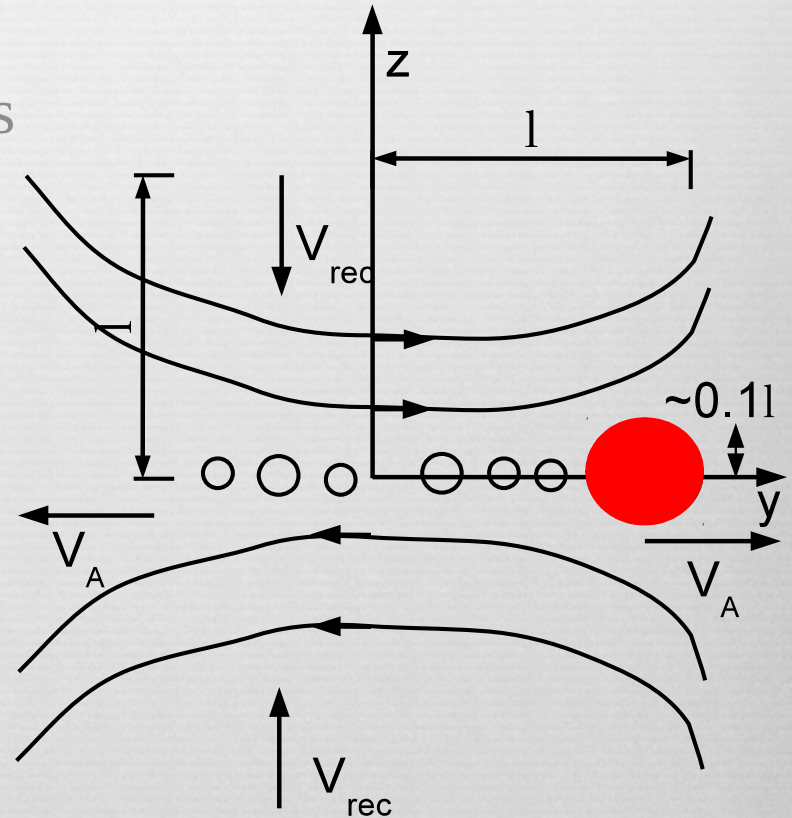
Plasmoid-dominated reconnection in blazars



- Current sheet fragments to plasmoids
Loureiro et al. 2007; Uzdensky et al. 2010; Loureiro et al. 2012+++

- Plasmoids merge/grow leaving the layer at $V_A \sim c$

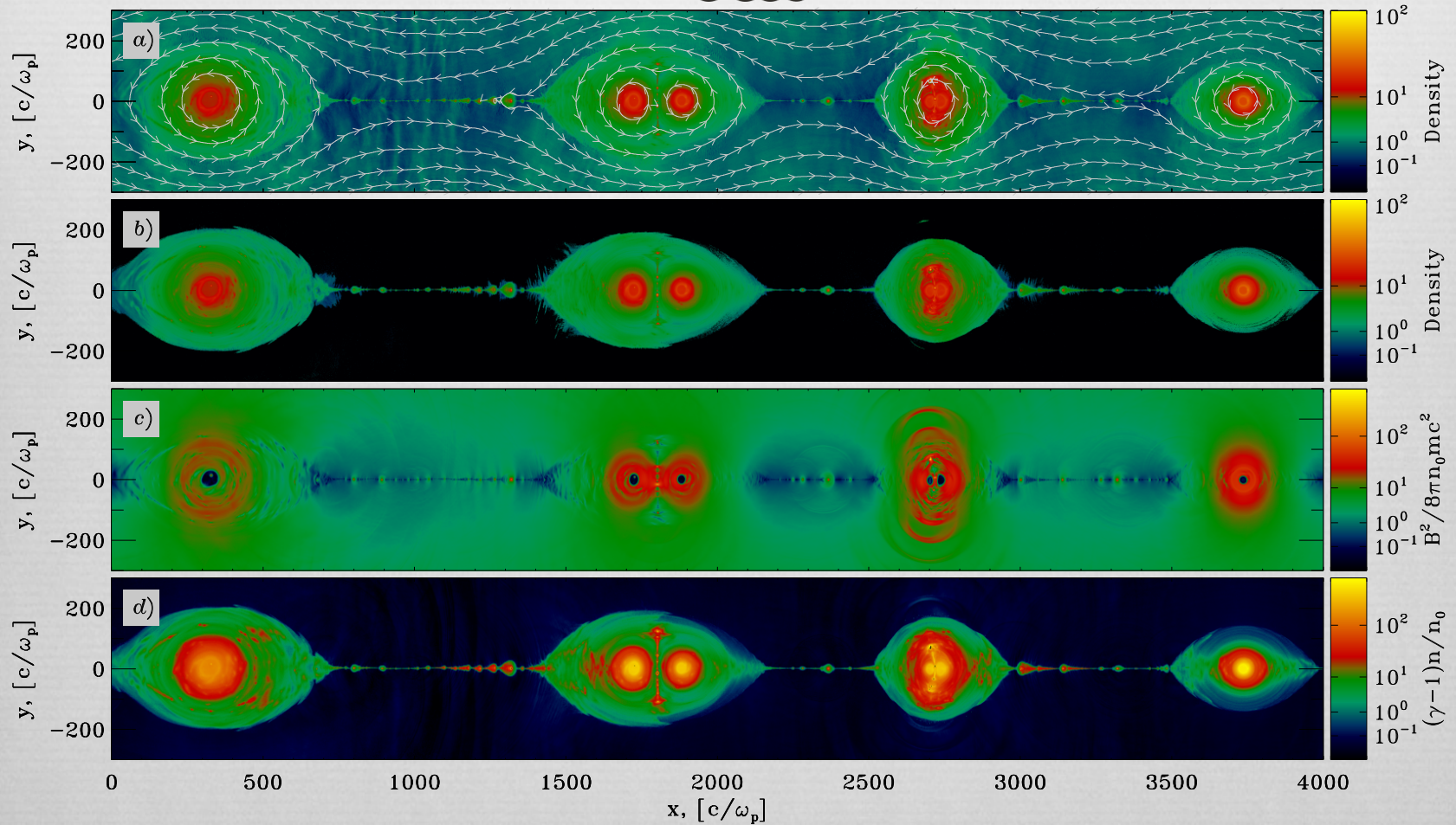
- The largest plasmoids can power bright/ultrafast blazar flares
Giannios 2013



Giannios et al. 2009; 2010; Giannios 2013

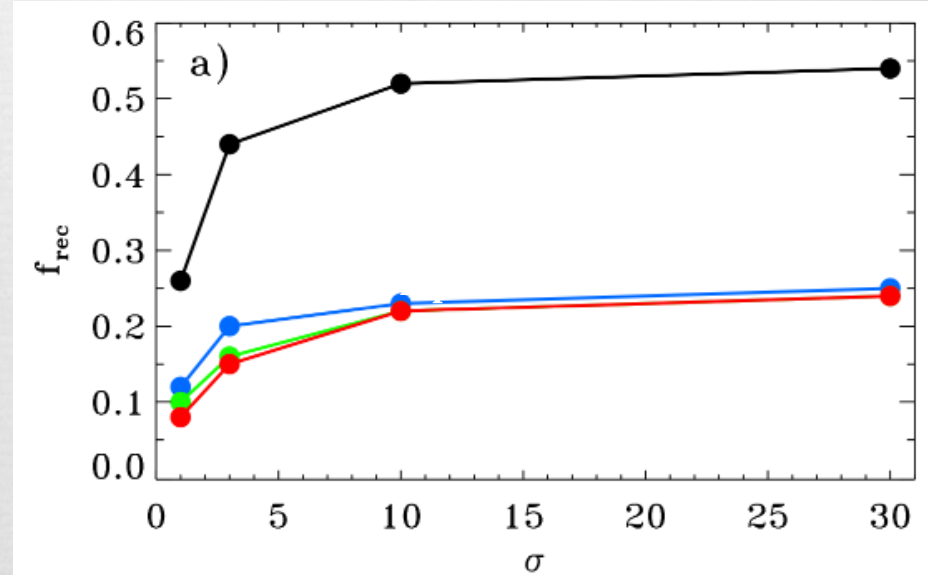
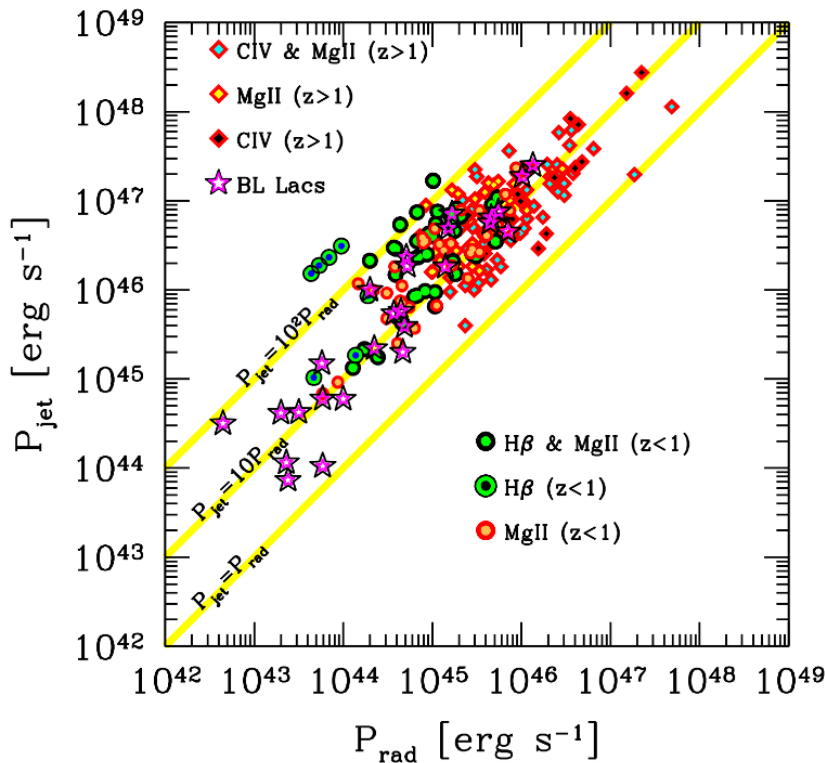
Magnetic reconnection

Sironi, Petropoulou & Giannios 2015



(1) Relativistic reconnection is efficient

$$f_{\text{rec}} \equiv \frac{\sum_i \int_{V_i} U_e dV_i}{\sum_i \int_{V_i} (e + \rho c^2 + U_B) dV_i}$$



(Sironi, Petropoulou, Giannios 15)

Blazar phenomenology:

- blazars are efficient emitters (radiated power \sim 10% of jet power)

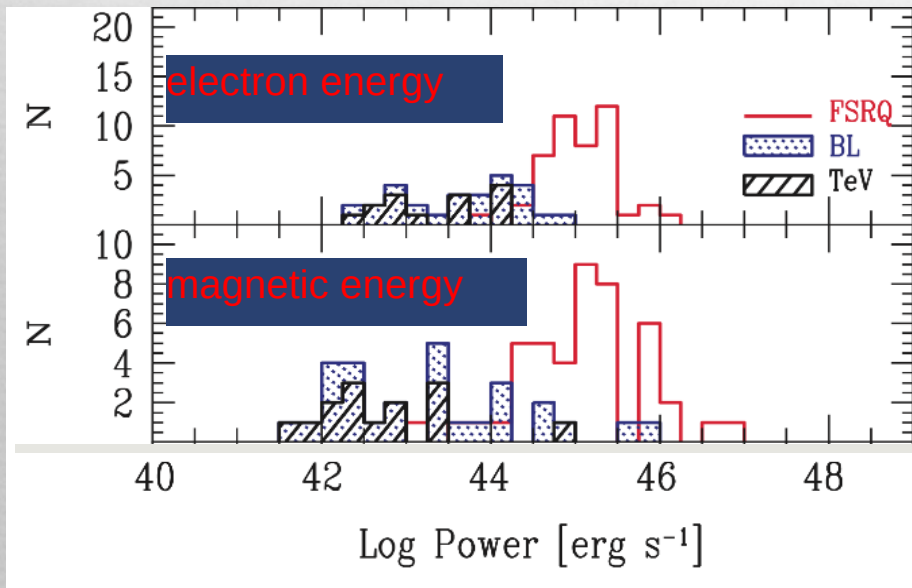
Relativistic reconnection:

- ✓ it transfers \sim 50% of the flow energy (electron-positron plasmas) or \sim 25% (electron-proton) to the emitting particles

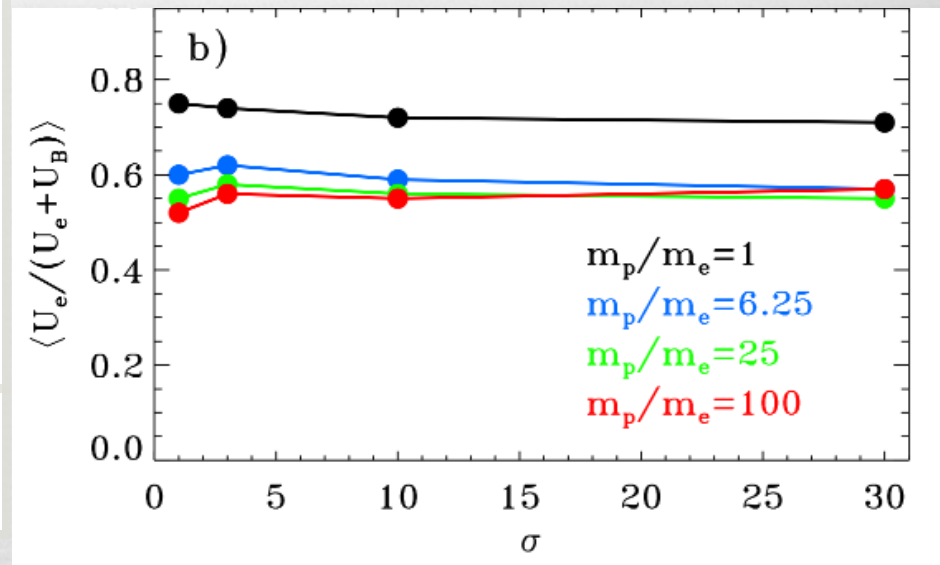
(2) Equipartition of particles and fields



$$\left\langle \frac{U_e}{U_e + U_B} \right\rangle \equiv \frac{\sum_i \int_{V_i} U_e \frac{U_e}{U_e + U_B} dV_i}{\sum_i \int_{V_i} U_e dV_i}$$



(Celotti+ 08)



(Sironi, Petropoulou & Giannios 2015)

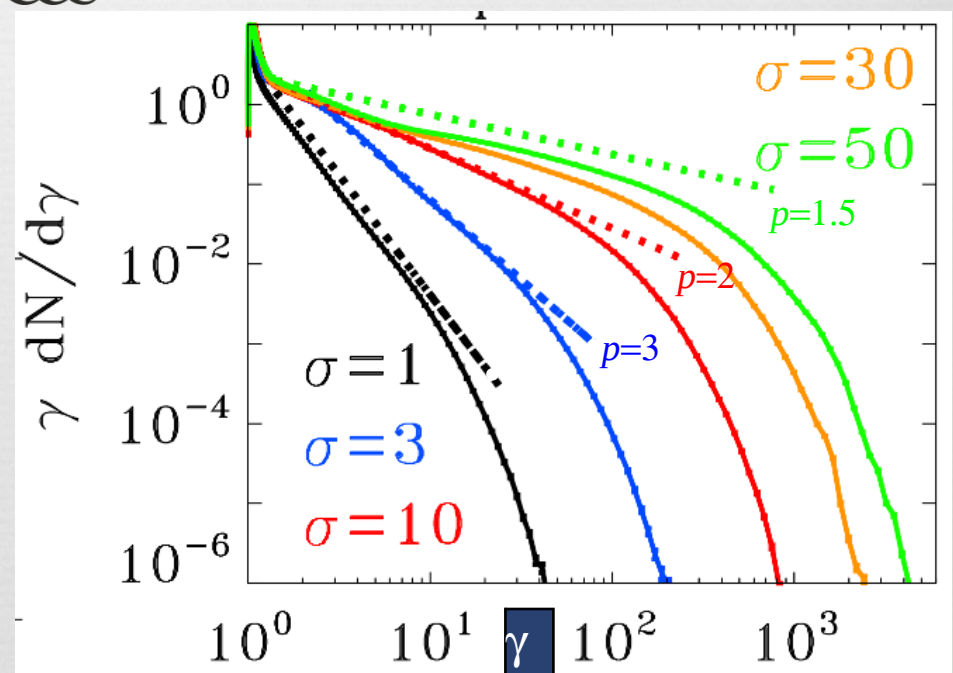
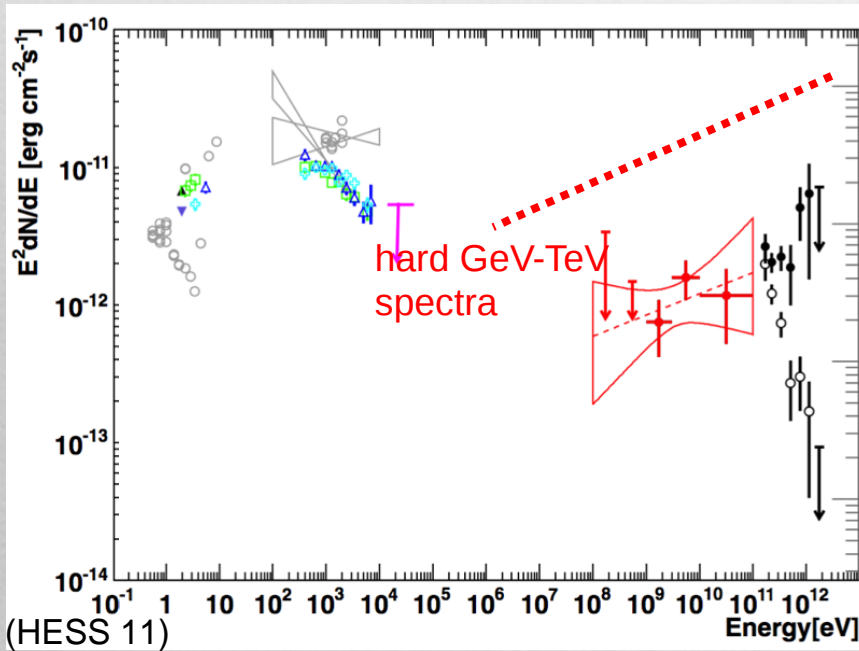
Blazar phenomenology:

- rough energy equipartition between emitting particles and magnetic field

Relativistic reconnection:

- ✓ in the magnetic islands, it naturally results in rough energy equipartition between particles and magnetic field

(3) Extended non-thermal distributions



(Sironi & Spitkovsky 14, Guo et al. 14, Werner et al. 14)

Blazar phenomenology:

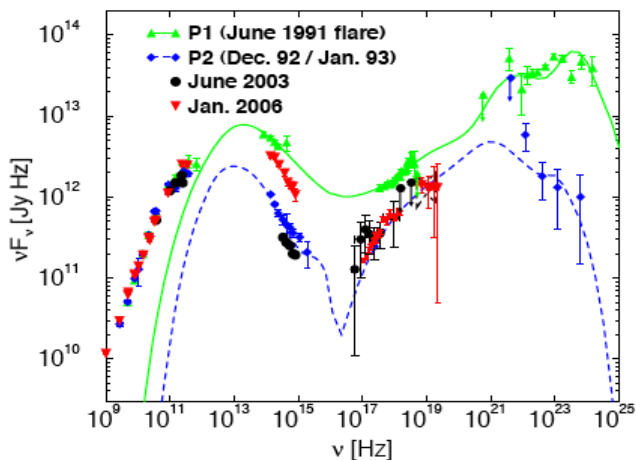
- extended power-law distributions of the emitting particles, with hard slope

$$\frac{dN}{d\gamma} \propto \gamma^{-p} \quad p \leq 2$$

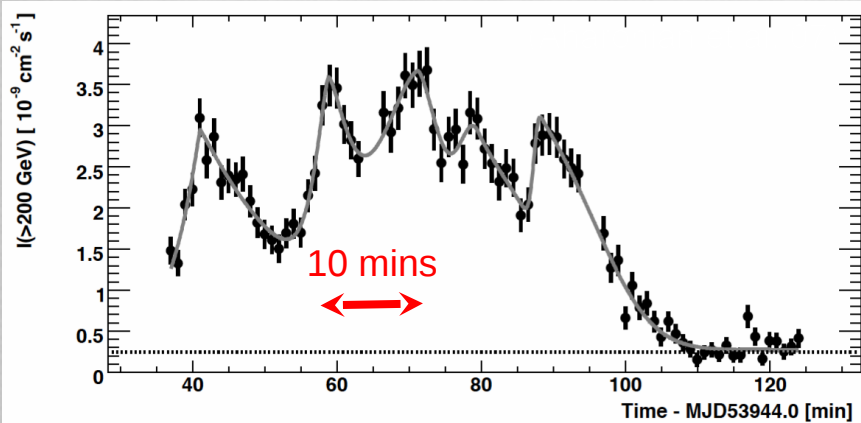
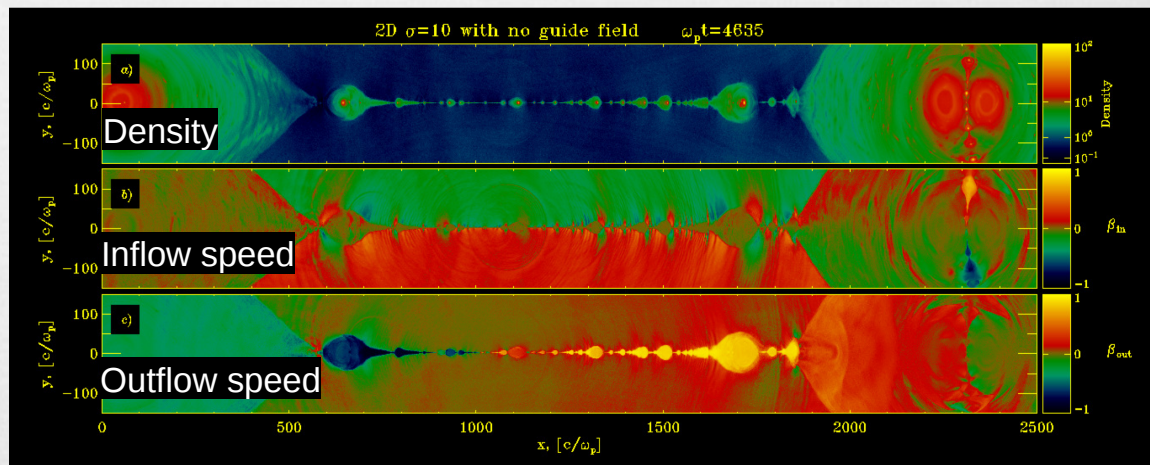
Relativistic reconnection:

- ✓ it produces extended non-thermal tails of accelerated particles, whose power-law slope can be harder than $p=2$

(4) Fast time variability at TeV energies

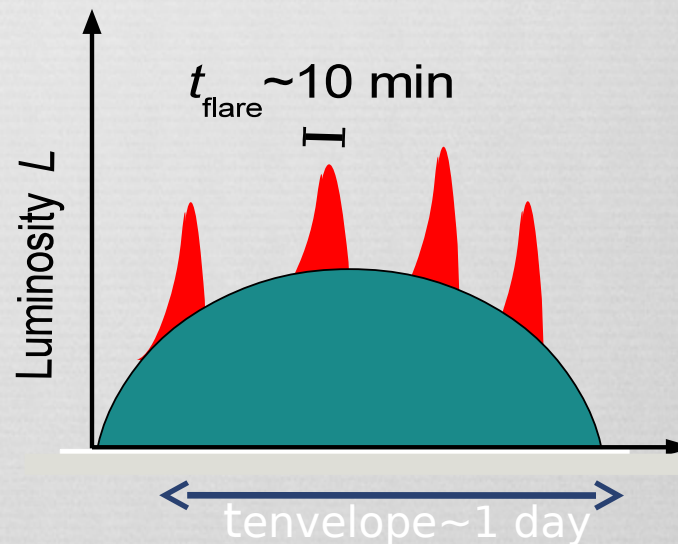


(Coppi & Aharonian 99, Boettcher 07)



Blazar phenomenology:

- at TeV energies, fast (~ 10 minutes) flares on a high-state envelope lasting for \sim days



Relativistic reconnection:

- ✓ large/fast islands might be a promising source of fast variability

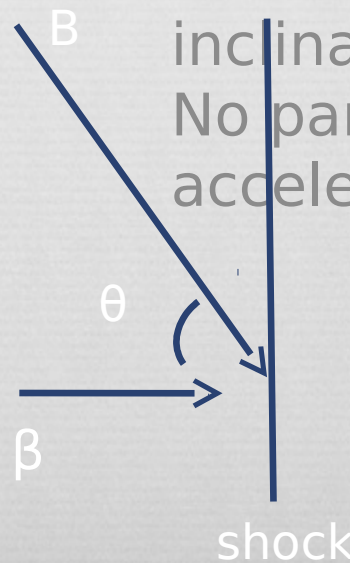
Shocks in jets



□ efficiency requires that the shocks are mildly *relativistic*, at least

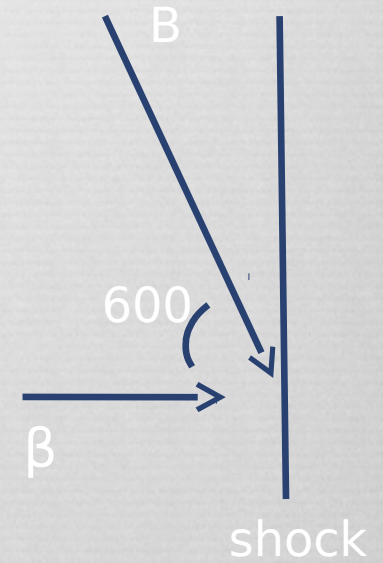
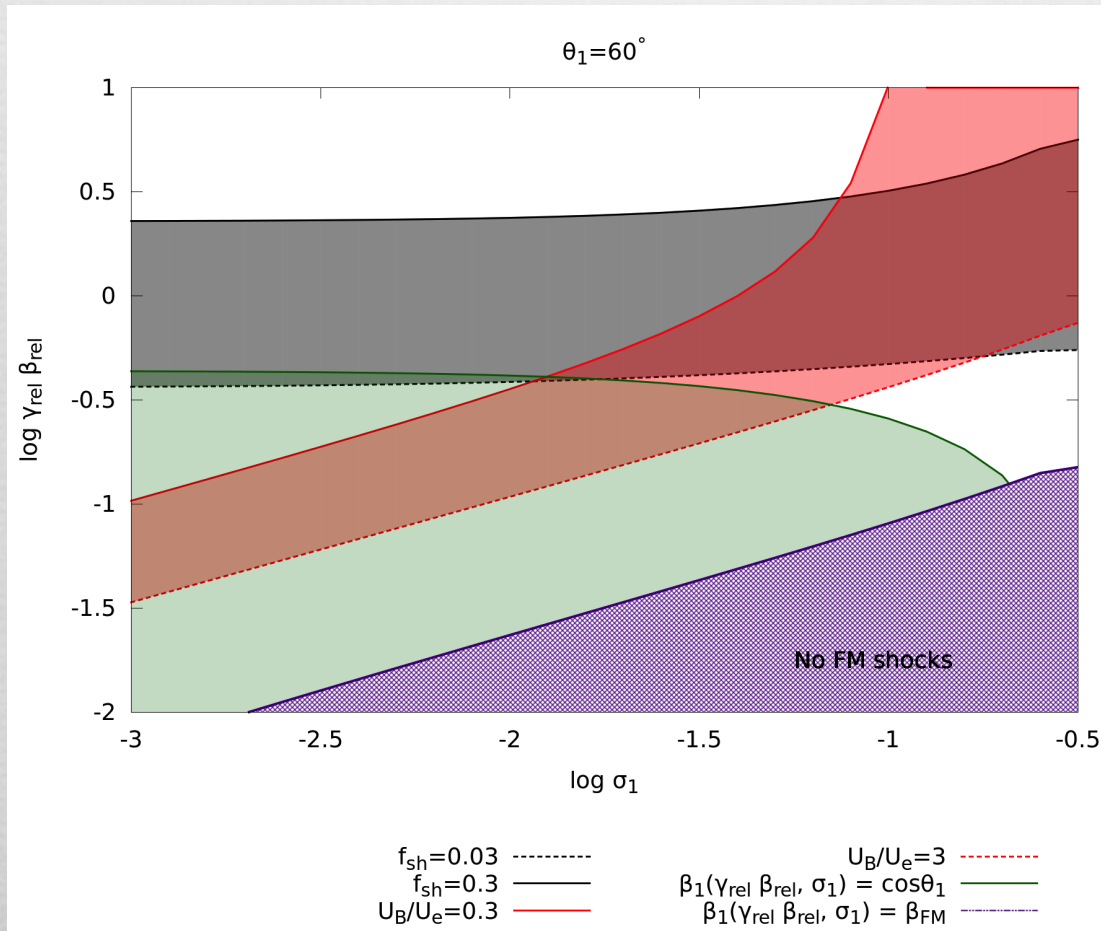
$$f_{sh} = \frac{U_{te}}{\rho c^2 + e} \approx 0.5 \frac{\gamma_{rel} - 1}{\gamma_{rel}}$$

□ (mildly) relativistic shocks are **superluminal** for most B-field inclinations □ No particle acceleration

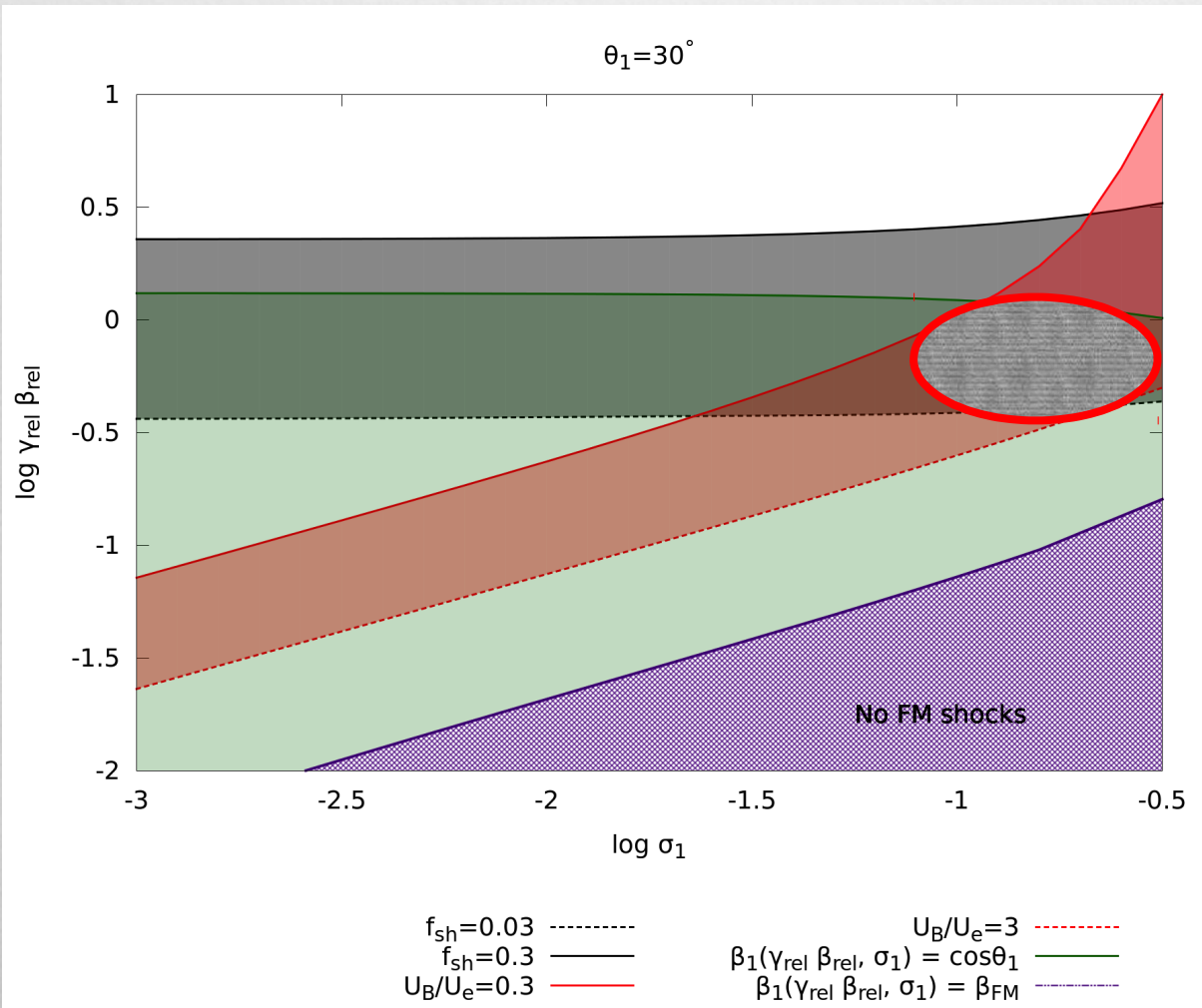


Shock is subluminal when $\beta < \cos \theta$

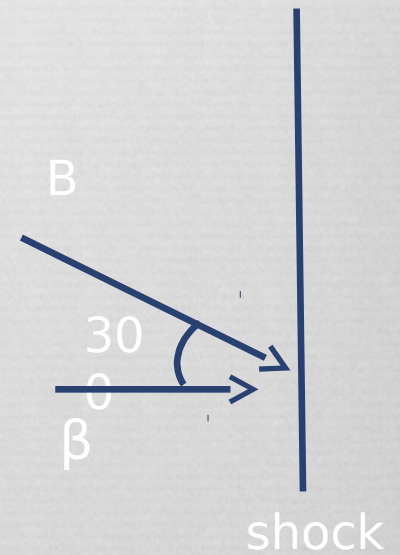
Quasi Perpendicular Shocks



Quasi-Parallel Shocks



Sironi, Petropoulou
& Giannios 2015



Concluding



□ Take Home messages

- 1) Are large-scale fields required to make Jets?
 - locally generated fields may launch jets as well
- 2) Jets are observed to be efficient particle accelerators, quasi equipartition objects
 - Shocks may not be both efficient *and* accelerate particles
 - Magnetic reconnection can be both and predicts equipartition conditions at the emitting region