

A cosmic-ray current driven instability in parallel shocks

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Cosmic-ray current driven instability

- ▶ We consider the region upstream of a quasi-parallel shock
- ▶ 3-component plasma
- ▶ relativistic beam of protons (Γ_b) along zeroth order field
- ▶ thermal electron/proton distribution $k_b T / mc^2 = \Theta \ll \Gamma_b$
- ▶ linear dispersion relation for circularly polarised waves
- ▶ plasma susceptibility:

$$\omega^2 \chi \approx \frac{\omega'_{pb}{}^2 \omega'}{\epsilon \omega_c} - \frac{\omega'_{pb}{}^2 \omega'}{\epsilon \omega_c + \omega'} + \frac{c^2 \omega^2}{v_A^2} + \frac{\omega_p^2 \omega}{\epsilon \omega_c^3} (c^2 k^2 - \omega^2) \langle u_{\perp}^2 \rangle_p$$

$$\omega \Theta_j, ck \Theta_j \ll |\omega_{cj}|$$

Cosmic-ray current driven instability

- ▶ Neglecting thermal effects, and provided $\Gamma_b \beta^2 (n_b/n_p) (\omega_{pp}/\omega_{cp})^2 \gg 1$
maximum growth rate of

$$\text{Im}(\omega) = \frac{1}{2} \frac{n_b}{n_p} \beta_b \omega_{pp}$$

Same result as Bell 2004

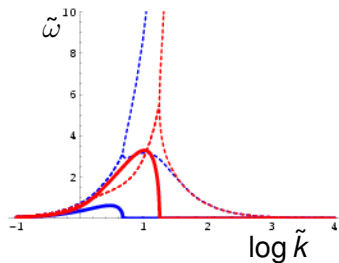
- ▶ Including thermal effects, if $\Theta_p \gg v_A/c$

$$\text{Im}(\omega) = \frac{\sqrt{3}}{2} \left(\frac{n_b}{n_p} \right)^{\frac{2}{3}} \left(\frac{v_A}{c} \right)^{\frac{2}{3}} \left(\frac{\omega_{pp}}{\omega_c} \right)^{\frac{2}{3}} \left(\frac{\beta_b^2}{\langle u_{\perp}^2 \rangle} \right)^{\frac{1}{3}} \omega_c$$

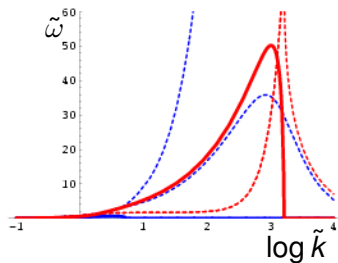
- ▶ May provide saturation mechanism in relativistic plasmas

Cosmic-ray current driven instability

e.g., $v_A = 2 \times 10^{-5}$, $\Gamma_b = 10$, $n_b/n_p = 1/10$, $\epsilon = -1$, $\epsilon = +1$



$$\Theta = 1/10$$



$$\Theta = 1/1000$$

cf. BR, Kirk & Duffy 2006, PPCF

Cosmic-ray current driven instability

- ▶ saturation when currents associated with waves:

$$|\mathbf{k} \times \mathbf{B}| \approx 4\pi n_{cr} e\beta$$

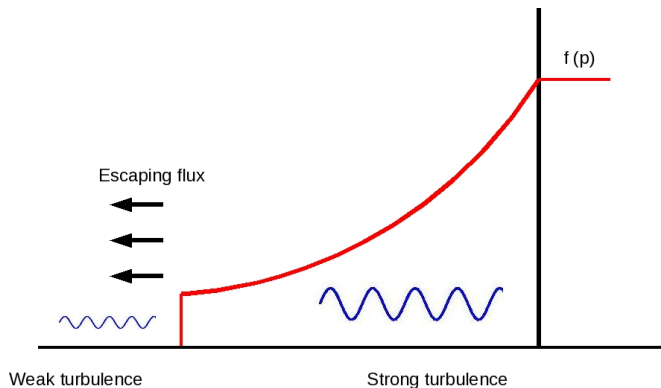
- ▶ $k \sim 1/r_g$ - saturated field energy

$$\frac{B_w^2}{8\pi} = \frac{1}{2} n_{cr} \Gamma_b m_p c^2$$

- ▶ Entire energy of the beam goes into magnetic field production
- ▶ How do we include this in acceleration models?

Cosmic-ray modified shocks

- ▶ Particles *escape* beyond some boundary



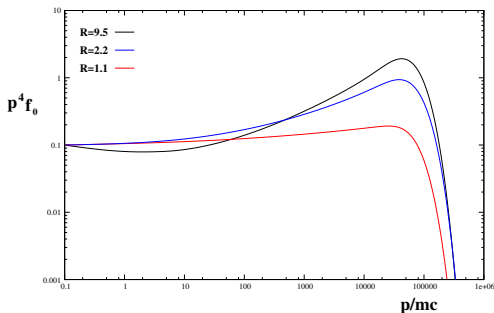
- ▶ escaping particles drive magnetic field growth

Steady-state solutions

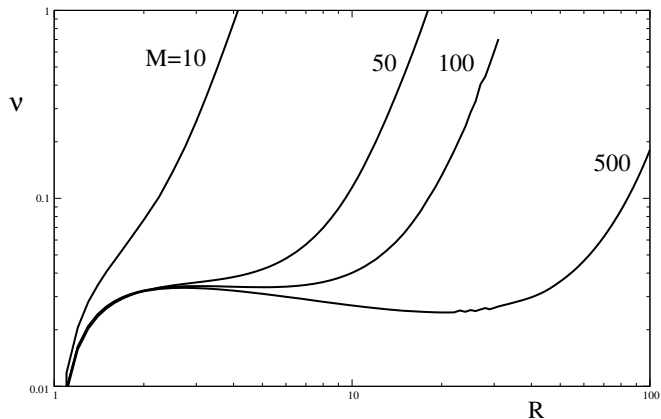
- ▶ coupled hydrodynamic - kinetic equations

$$\frac{\partial f}{\partial t} + \frac{\partial}{\partial x} \left(uf - \kappa(x, p) \frac{\partial f}{\partial x} \right) = \frac{\partial}{\partial p^3} \left(p^3 f \frac{\partial u}{\partial x} \right)$$

$$L_{\text{esc}} = \frac{\kappa(p^*)}{u_0}, \quad \text{e.g. } p^* = 10^5$$



Injection efficiency ν – R diagram



$$u_0 = 5000 \text{ km s}^{-1}, n_0 = 1 \text{ cm}^{-3}, p^* = 10^3$$

$$\nu = \frac{4\pi}{3} \frac{mc^2}{\rho_0 u_0^2} p_0^4 f_0(p_0),$$

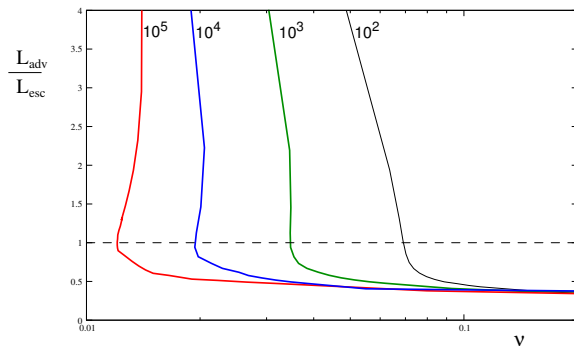
Maximum momentum

- ▶ what is the location of escape boundary - transition zone from weak to strong turbulence
- ▶ $L_{\text{esc}} = \kappa(p^*)/u_0$ determines maximum confined energy
- ▶ diffusive current at escape boundary drives nonresonant instability

$$j_{\text{cr}}(-L_{\text{esc}}) = -4\pi e \int_{p_0}^{\infty} \kappa \frac{\partial f}{\partial x} p^2 dp$$

- ▶ How do we determine L_{esc} in a self-consistent manner?
- ▶ Calculate CR flux (and p^*) from numerical SS solution
- ▶ Transition zone determined from condition
 $L_{\text{adv}} \equiv u_{\text{sh}}/\Gamma_{\text{max}} \sim L_{\text{esc}}$

Maximum momentum



BR, Kirk & Duffy, ApJ submitted

- ▶ Maximum energy can be calculated as a function of injection parameter

Summary

- ▶ Efficient acceleration of particles at collisionless shocks connected with magnetic field amplification
- ▶ nonresonant mode (Bell 2004) appears to be of greatest importance in efficient cosmic-ray accelerating shocks
- ▶ reduced diffusion coefficients leads to more rapid acceleration
- ▶ field amplification incorporated using free escape boundary L_{esc}
- ▶ maximum energy determined by boundary position - calculated self-consistently
- ▶ system is self-organising - maximum energy function of injection