Radiation-Dominated Relativistic Current Sheets

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Relativistic Current Sheet Instabilities

$V_A \sim c, T/mc^2 > 1$, Electron and Positron Plasmas



Open issue for relativistic Current Sheet

radiation effect such as synchrotron cooling

$$\frac{\tau_{loss}}{\tau_{dyn}} \approx \left(\frac{10^2}{\tau_{dyn}\Omega_c}\right) \left(\frac{10^{14}G}{B}\right) \left(\frac{1}{E/mc^2}\right)^2$$
e.g. Magnetor,
Pulsar magnetosphere...







Abraham-Lorentz Formula for Radiation Drag Force

$$mc \frac{du^i}{ds} = \frac{e}{c} F^{ik} u_k + g^i$$
 (Dirac Form)

$$g^{i} = \frac{2e^{2}}{3c} \left(\frac{d^{2}u^{i}}{ds^{2}} + u^{i} \frac{du^{k}}{ds} \frac{du_{k}}{ds} \right)$$

$$= \frac{2e^{3}}{3mc^{3}} \frac{\partial F^{ik}}{\partial x^{l}} u_{k} u^{l} - \frac{2e^{4}}{3m^{2}c^{5}} F^{ik} F_{lk} u^{l} + u^{i} \cdot \frac{2e^{4}}{3m^{2}c^{5}} (F^{kl}u_{l}) (F_{km}u^{m})$$

$$= T_{1} + T_{2} + T_{3}$$

Radiation Loss Effect in PIC Simulation Code

Abraham-Lorentz Formula for Radiation Drag Force

$$\mathbf{T}_{1} = \frac{2}{3}\gamma \cdot (\omega_{c0}\tau_{0}) \cdot (mc\omega_{c0}) \cdot ((\hat{\partial}_{t} + \hat{\mathbf{v}} \cdot \hat{\nabla})\hat{\mathbf{E}} + \hat{\beta} \times (\hat{\partial}_{t} + \hat{\mathbf{v}} \cdot \hat{\nabla})\hat{\mathbf{B}})$$

$$\mathbf{T}_{2} = \frac{2}{3} \cdot (\omega_{c0}\tau_{0}) \cdot (mc\omega_{c0}) \cdot (\hat{\mathbf{E}} \times \hat{\mathbf{B}} + \hat{\mathbf{B}} \times (\hat{\mathbf{B}} \times \beta) + \hat{\mathbf{E}}(\beta \cdot \hat{\mathbf{E}}))$$

$$\mathbf{T}_{3} = -\frac{2}{3}\gamma^{2} \cdot (\omega_{c0}\tau_{0}) \cdot (mc\omega_{c0}) \cdot \beta \cdot ((\hat{\mathbf{E}} + \beta \times \hat{\mathbf{B}})^{2} - (\hat{\mathbf{E}} \cdot \hat{\beta})^{2})$$

 τ_0 : Light crossing time over classical electron radius (e^2/mc^2)/c ~ 10^-23 s Main Radiation Effect is Synchrotron Radiation

(cf. Noguchi et al. 2005)

Time Evolution of MR & DKI



Growth Curves

 α : radiation loss coefficient ($\alpha = 0$: No radiation loss, $\alpha = 10^{-10}$: Strong radiation loss)

Drift-Kink Mode

Tearing Mode (Reconnection)



Comparison of Growth Rate



Reason for Fast Growth

- Decrease of gas pressure by radiation loss
- Shrink of plasma sheet & thin plasma sheet

Why reconnection has super-fast growth?

• Temperature Anisotropy $T_{//} \neq T_{\perp}$

Temperature Anisotropy (Early Stage)



Linear Growth Rate (Γ) under Temperature Anisotropy

• Tearing Mode

- strong dependence on $T_{\perp}/T_{//}$ (e.g. Chen et al., 1984)
- $-T_{\perp} > T_{//} \rightarrow \Gamma$ increase
- $-T_{\perp} < T_{//} \rightarrow \Gamma$ decrease



 $T_{\perp}/T_{//} \rightarrow isotropization$

• Drift-Kink Mode – weak dependence on $T_{\perp}/T_{//}$



Late Nonlinear Stage

 $\alpha = 0$

$\alpha = 10^{-12}$

$\alpha = 10^{-11}$



Tearing

Drift-Kink

Temperature Anisotropy in Late Nonlinear Stage





reconnection suppressed, shifted to small k mode

Relativistic Reconnection under Synchrotron Cooling



Summary

- Relativistic Reconnection with Radiation Cooling
 - thin current sheet due to radiation cooling
 - fast growth of reconnection/drift-kink instability
- Nonlinear Evolution of Reconnection/Drift-Kink
 - Early Stage: $T_{\perp} > T_{//}$
 - Super-Fast Reconnection
 - Fast Drift-Kink Instability (weak effect of Temp. Anisotropy)
 - Late Nonlinear Stage: $T_{\perp} < T_{//}$
 - Transition to Sweet-Parker Reconnection